

DOCUMENT RESUME

ED 211 340

SE 035 956

AUTHOR Onega, Ronald J.
TITLE Nuclear Engineering Computer Modules: Reactor Dynamics, RD-1 and RD-2.
INSTITUTION Virginia Polytechnic Inst. and State Univ., Blacksburg. Coll. of Engineering.
SPONS AGENCY National Science Foundation, Washington, D.C.
PUB DATE [73]
GRANT NSF-GZ-2888; NSF-SEB-73-06276.
NOTE 148p.; Not available in paper copy due to marginal legibility of original document.
EDRS PRICE MF01 Plus Postage. PC Not Available from EDRS.
DESCRIPTORS College Science; *Computer Programs; Engineering; *Engineering Education; Higher Education; *Kinetics; *Mathematical Models; *Nuclear Physics; Physics; Science Education
IDENTIFIERS *Nuclear Reactors

ABSTRACT

The objective of the Reactor Dynamics Module, RD-1, is to obtain the kinetics equation without feedback and solve the kinetics equations numerically for one to six delayed neutron groups for time varying reactivity insertions. The computer code FDMCKI (Fundamental Mode Kinetics) will calculate the power as a function of time for either uranium or plutonium. Either fuel can be used with one to six delayed neutron groups and one of three types of reactivity insertions: a constant reactivity, sinusoidal, or a ramp. The code does not compute any parameters so the neutron generation time must be provided. The user has the option of studying the effects of various time steps in solving the system. The objective of Module RD-2 is to examine the temperature feedback mechanism of a pressurized water reactor (PWR) and solve the one delayed neutron model with temperature feedback for a step insertion and a ramp insertion of reactivity. A PWR core with a two-path feedback is considered. The core region is the only one of interest in this module. The program name is FUMOTEM (Fundamental Mode Kinetics with Temperature Feedback). There are four types of reactivity inputs that the program can accommodate. (Author)

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REACTOR DYNAMICS MODULE, RD-1
THE REACTOR KINETICS EQUATIONS

by

Ronald J. Onega

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The University gratefully acknowledges the support of the
Division of Higher Education of the National Science Foundation
for support of this work performed under Grant GZ-2888 and the
support of Duke Power Company, North Carolina Power and Light Company,
and Virginia Electric and Power Company.

Project Director: Milton C. Edlund

ACKNOWLEDGMENT

The author would like to extend his appreciation and thanks to Mr. Tjeri Surjanto who did the programming for this module.

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KINETICS MODULE 1

THE REACTOR KINETICS EQUATIONS

1.1 Object of Module

The object of this module is to:

- 1) Obtain the Kinetics Equation without feedback
- and 2). Solve the kinetics equations numerically for one to six delayed neutron groups for time varying reactivity insertions.

The time dependence of a modern reactor is really very complicated. The control rod motion is a local perturbation so the time dependence of the flux cannot completely be divorced from the space dependence. The fundamental mode kinetics equations do provide a rough idea of what the time behavior of a reactor will be. In this module we will develop the kinetics equations and indicate how they can be solved numerically. Feedback effects will be introduced in later modules.

The computer code FUMOKI (Fundamental Mode Kinetics) will calculate the power as a function of time for either uranium or plutonium. Either fuel can be used with one to six delayed neutron groups and one of three types of reactivity insertions:

- 1) a constant reactivity ρ_0
- 2) sinusoidal $\rho(t) = \rho_0 \sin b_2 t$
- or 3) a ramp $\rho(t) = \rho_0 (1 + b_3 t)$.

The ramp is to simulate the rod withdrawal or insertion reactivity input.

The code does not compute any parameters so the neutron generation time must be provided. Also the user has the option of studying the effects of various time steps in solving the system. The time step can be surprisingly large (0.05 sec) and still yield good results in most cases.

1.2 The Kinetics Equations

The time behavior of a reactor is a very important consideration in the operation of a nuclear power plant. Also, the safety analysis of a plant depends upon a thorough knowledge of the kinetics equations. The many types of reactor designs necessitates consideration of various reactivity coefficients and dynamic response characteristics.

The neutronic considerations of a reactor cannot be divorced from associated feedback mechanisms such as heat transport, fluid flow, mechanical changes etc. There are many ways to delineate the dynamics problems of a power plant but a natural one seems to be to classify problems according to the time constants involved. There are four areas that we can study:

- 1) Very slow transients - fuel depletion with time constants of the order of a year or so. We will not consider this as a dynamics problem at all but rather in the statics sections.
- 2) Slow transients - Xenon and Samarium effects. The time constants here are of the order of hours.
- 3) Normal transients - Changes in fuel and moderator temperature, void changes, delayed neutron considerations etc. The time constants here are of the order of a second or so and small reactivity changes are involved.
- 4) Fast transients - Control rod dropped in or withdrawn at its maximum rate etc. Reactivity inputs are of 50¢ and up. The time constants are of the order of 10^{-4} sec. and serious safety questions are raised.

In addition to these rather general time reference frames, another important aspect of the reactor dynamics problem is whether or not the core is so large that spatially dependent analysis is necessary. The solution of one or more dimensional kinetics problems necessitates the use of a relatively large computer. Xenon oscillations are of importance here. Also reactivity insertions are usually localized so that hot spots may develop.

Coupled core kinetics considerations are of interest in some types of reactors with large cores or core regions which are loosely coupled. When the neutron flight times between different regions of the reactor are not negligible, then coupled core kinetics may be a useful tool. Coupled core kinetics equations are generally differential-integral equations with time-lag kernels. Coupled core kinetics involves writing the kinetics equations for each region of the reactor and then coupling the regions by neutron leakage from one region to the other.

When a control rod in a reactor is removed, the neutron flux distribution is disturbed so that not only the fundamental mode is present but the higher modes are also present. However, these higher harmonics die out very rapidly so that while the rod is in motion (in the order of seconds) the fundamental mode is the only important one. Therefore we will deal only with the fundamental mode for the space dependence and the kinetics equations we solve will simply indicate how the amplitude of this fundamental mode changes with time. For example, in a spherical reactor the flux is

$$\phi(r,t) = \phi_0(t) \frac{\sin Br}{Br}$$

and this spatial dependence persists through most transients, and the kinetics equations simply yield the $\phi_0(t)$. In a word, the spatial and time dependence are separable for most transients.

We will derive the kinetics equations using one group diffusion theory but the same equations are obtained from multi-group or transport theory. The thermal neutron diffusion equation is

$$D \nabla^2 \phi - \Sigma_a \phi + S = \frac{1}{v} \frac{\partial \phi}{\partial t} \quad (1.2.1)$$

The source is composed of three terms, i.e.

$$S = S_p + S_D + S_{ext} \quad (1.2.2)$$

where S_p = prompt neutron source,

S_D = delayed neutron source,

and S_{ext} = a source of neutrons entered externally, as from a Plutonium-Beryllium source.

Not all neutrons are emitted immediately from the fission process.

The prompt neutrons are emitted within 10^{-10} sec. of the fission process

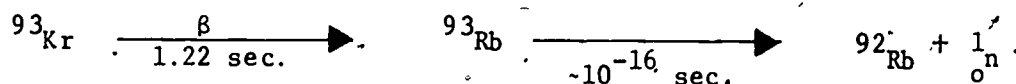
itself but there are other neutrons called delayed neutrons. They arise

from the beta decay of some of the fission fragments. As an example,

Kr-93 is a fission fragment which has a half life of 1.22 sec. This beta-

decays to Rb-93, but it is formed in such a highly excited state that a

neutron is emitted following the beta decay and the reaction is thus,



There are some 30 of these neutron precursors (${}^{93}_{\text{Kr}}$) which produce neutrons

through beta decay.

Since the half-lives of many of these precursors are very close to each other, it is necessary only to consider six delayed neutron precursor groups which are averages of the 30 or so precursors, appropriately weighted. The delayed neutron precursors we talk about are thus fictitious in that they do not actually exist, but are averages of the actual precursors. Table 1.5.1 (section 5) indicates the groups of delayed neutrons obtained from these six precursors for neutron induced fission in uranium and plutonium.

The fraction of neutrons which are delayed is called β . For ^{235}U , $\beta = 0.0064$ and the total delayed neutron fraction is the sum of the individual ones, i.e.,

$$\beta = \sum_{i=1}^6 \beta_i \quad (1.2.3)$$

Returning now to the diffusion equation, we can set

$$S_p = (1-\beta) \cdot v \Sigma_f \phi(\vec{r}, t) \quad (1.2.4)$$

and

$$S_D = \sum_{i=1}^6 \lambda_i C_i(\vec{r}, t) \quad (1.2.5)$$

where $C_i(\vec{r}, t)$ is the precursor concentration of group "i", i.e., the number of precursors/cm³ existing at point \vec{r} at time t . Each time a precursor C_i decays, it is assumed a delayed neutron of group "i" results.

So our diffusion equation becomes (Equation 1.2.1)

$$D \nabla^2 \phi - \Sigma_a \phi + (1-\beta) v \Sigma_f \phi + \sum_{i=1}^6 \lambda_i C_i(\vec{r}, t) = \frac{1}{v} \frac{\partial \phi}{\partial t} \quad (1.2.6)$$

Equation(1.2.6) has four independent variables, x, y, z and t , and seven dependent variables $\phi, C_1, C_2, \dots, C_6$. In order to solve such a system of equations, we will have to write six more equations, one for each of the precursor concentrations. The rate at which the precursor concentration changes is

$$\frac{\partial C_i(\vec{r}, t)}{\partial t} = \underbrace{\beta_i \nu \Sigma_f \phi(\vec{r}, t)}_{\text{rate of formation}} - \underbrace{\lambda_i C_i(\vec{r}, t)}_{\text{rate of decay}}. \quad (1.2.7)$$

$$i = 1, 2, \dots, 6.$$

This system of equations is relatively difficult to solve so we shall work analytically with Equation (1.2.6) and (1.2.7) for a while.

Assume that we want to expand the spatial dependence of the flux $\phi(\vec{r}, t)$ and the precursor concentrations $C_i(\vec{r}, t)$ in terms of a set of eigenfunctions of the Helmholtz equation

$$\nabla^2 Y_n(\vec{r}) + B_n^2 Y_n(\vec{r}) = 0. \quad (1.2.8)$$

The reason for doing this is that the solution for the steady state satisfies this same equation with coefficients which are time independent. In the time dependent problems the coefficients are functions of time. The eigenfunctions $Y_n(\vec{r})$ are $\cos B_n x$ for a slab reactor, Bessel's functions for a cylindrical reactor, etc., In any event, we set

$$\phi(\vec{r}, t) = \sum_{n=0}^{\infty} \phi_n(t) Y_n(\vec{r}) \quad (1.2.9)$$

and

$$C_i(\vec{r}, t) = \sum_{n=0}^{\infty} \tilde{C}_{in}(t) \cdot Y_n(\vec{r}). \quad (1.2.10)$$

We note that $\tilde{\phi}_n(t)$ will yield the amplitude of the n^{th} harmonic and $Y_n(\vec{r})$ its spatial distribution. Substituting these expressions into Equations (1.2.6) and (1.2.7) and using Equation (1.2.8) we obtain

$$\sum_{n=0}^{\infty} \left[-DB_n^2 \tilde{\phi}_n(t) - \Sigma_a \tilde{\phi}_n(t) + (1-\beta) v \Sigma_f \tilde{\phi}_n(t) \right] Y_n(\vec{r}) + \sum_{i=1}^6 \sum_{n=0}^{\infty} \lambda_i \tilde{C}_{in}(t) Y_n(\vec{r}) = \frac{1}{v} \sum_{n=0}^{\infty} Y_n(\vec{r}) \frac{d\tilde{\phi}_n(t)}{dt}$$

and

$$\sum_{n=0}^{\infty} Y_n(\vec{r}) \frac{d\tilde{C}_{in}(t)}{dt} = \sum_{n=0}^{\infty} \left[\beta_i v \Sigma_f \tilde{\phi}_n(t) - \lambda_i \tilde{C}_{in}(t) \right] Y_n(\vec{r}).$$

Now using the fact that the $Y_n(\vec{r})$ form an orthogonal set, the preceding equation is simplified by multiplying it by $Y_m(\vec{r})$ and integrating over the reactor volume (taking the inner product) to yield

$$\frac{1}{v} \frac{d\tilde{\phi}_n(t)}{dt} = -(DB_n^2 + \Sigma_a) \tilde{\phi}_n(t) + (1-\beta) v \Sigma_f \tilde{\phi}_n(t) + \sum_{i=1}^6 \lambda_i \tilde{C}_{in}(t), \quad (1.2.11)$$

and

$$\frac{d\tilde{C}_{in}(t)}{dt} = \beta_i v \Sigma_f \tilde{\phi}_n(t) - \lambda_i \tilde{C}_{in}(t). \quad (1.2.12)$$

Now it is necessary to solve these two equations for the expansion coefficients $\tilde{\phi}_n(t)$ and $\tilde{C}_{in}(t)$. This is a very difficult task so we

make some definitions and some approximations.

First we define the multiplication k_n as

$$k_n = \frac{v \Sigma_f / \Sigma_a}{1 + B_n^2 L^2} \quad (1.2.13)$$

and the thermal neutron lifetime

$$l_n = \frac{1}{v \Sigma_a (1 + B_n^2 L^2)} \quad (1.2.14)$$

The next thing we do is assume that there are no delayed neutrons (there are, but assume for the moment we can neglect them). Then Equation (1.2.11) becomes (if $\beta = 0$),

$$\frac{d\tilde{\phi}_n}{dt}(t) = -v \Sigma_a (1 + L^2 B_n^2) \tilde{\phi}_n(t) + v v \Sigma_f \tilde{\phi}_n(t).$$

Now using Equation (1.2.14) we have

$$\begin{aligned} \frac{d\tilde{\phi}_n}{dt}(t) &= -\frac{\tilde{\phi}_n}{l_n}(t) + v v \frac{\Sigma_f}{\Sigma_{au}} \frac{\Sigma_{au}}{\Sigma_a} \Sigma_a \tilde{\phi}_n(t), \\ &= -\frac{\tilde{\phi}_n}{l_n}(t) + \frac{1 + B_n^2 L^2}{1 + B_n^2 L^2} \tilde{\phi}_n(t), \\ &= \frac{k_n - 1}{l_n} \tilde{\phi}_n(t). \end{aligned} \quad (1.2.15)$$

The solution of Equation (1.2.15) is obviously

$$\tilde{\phi}_n(t) = \tilde{\phi}_n(0) e^{\frac{k_n - 1}{l_n} t} \quad (1.2.16)$$

For a typical light water reactor, $\ell_o = \ell = 10^{-4}$ sec. and $k_o = k = 1.0020$.

From Equation (1.2.14) we note that

$$\frac{\ell_n}{\ell} = \frac{1 + B_o^2 L^2}{1 + B_n^2 L^2} \sim \frac{1}{n^2}$$

so that the higher harmonics die out rapidly. We will keep only the fundamental mode, not only in the solution of the no delayed neutron case but also for the solution of Equations (1.2.11) and (1.2.12).

Returning now to Equations (1.2.11) and (1.2.12) and omitting the subscript n since we are concerned only with the fundamental mode, we set

$$\phi(t) = v n(t)$$

where $n(t)$ is the neutron density (cm^{-3}), and after a bit of algebra come up with

$$\frac{dn(t)}{dt} = (1-\beta) \frac{\Sigma_f}{\Sigma_a} v \Sigma_a n(t) - v \Sigma_a (1 + B^2 L^2) n(t) + \sum_{i=1}^6 \lambda_i \tilde{C}_i(t)$$

and

$$\frac{d\tilde{C}_i(t)}{dt} = \beta_i v \Sigma_a \frac{\Sigma_f}{\Sigma_a} n(t) - \lambda_i \tilde{C}_i(t).$$

Using the definitions given in Equations (1.2.13) and (1.2.14) the above equations become

$$\frac{dn(t)}{dt} = (1-\beta) \frac{k}{\ell} n - \frac{n}{\ell} + \sum_{i=1}^6 \lambda_i \tilde{C}_i(t)$$

and

$$\frac{d\tilde{C}_i(t)}{dt} = \beta_i \frac{k}{\ell} n(t) - \lambda_i \tilde{C}_i(t).$$

k is the effective multiplication, and the effective neutron lifetime is denoted as ℓ . If we drop the tilde on the C_i 's, then we have

$$\frac{dn(t)}{dt} = \frac{(1 - \beta) k - 1}{\ell} n(t) + \sum_{i=1}^6 \lambda_i C_i(t) \quad (1.2.18)$$

and

$$\frac{dC_i(t)}{dt} = \beta_i \frac{k}{\ell} n(t) - \lambda_i C_i(t). \quad (1.2.19)$$

These are called the fundamental mode reactor kinetics equations, in that spatial and time dependence are assumed separable.

Sometimes it's convenient to cast these equations into a slightly different form. The thermal neutron lifetime is

$$\ell = \frac{\lambda_a}{v} \frac{1}{1 + B^2 L^2} \quad (1.2.20)$$

i.e., the mean time a neutron spends in the system from its birth as a thermal neutron until it's absorbed or leaks out of the reactor. The neutron generation time is defined as

$$\Lambda = \frac{\ell}{k} = \frac{1}{v \Sigma_a (1 + B^2 L^2)} \cdot \frac{\Sigma_a (1 + B^2 L^2)}{v \Sigma_f} = \frac{1}{v \Sigma_f} \quad (1.2.21)$$

The generation time is the mean time that it takes one neutron to generate one more prompt neutron or one precursor. The neutron lifetime is thus the reciprocal of the destruction rate of neutrons and the generation time is the reciprocal of the production rate of neutrons. With this definition and the definition of reactivity ρ ,

$$\rho \equiv \frac{k-1}{k}, \quad (1.2.22)$$

the kinetics equations (1.2.18 and 1.2.19) become

$$\frac{dn}{dt}(t) = \frac{\rho(t) - \beta}{\Lambda} n(t) + \sum_{i=1}^6 \lambda_i C_i(t) \quad (1.2.23)$$

$$\frac{dC_i}{dt}(t) = \frac{\beta_i}{\Lambda} n(t) - \lambda_i C_i(t) \quad (1.2.24)$$

$$i = 1, 2, \dots, 6$$

Equations (1.2.23) and (1.2.24) are used to describe the time behavior of a reactor. This form of the equations results if multigroup diffusion theory or transport theory is used to derive the equations but the definitions of ℓ , Λ , ρ , k are modified. We will assume they are input parameters to the system of equations.

Problem 1.2.1

Estimate the neutron lifetime and neutron generation time in an infinite stack of graphite and U-235 if $\sigma_a = 5 \times 10^{-3}$ b for the graphite and there are 500 atoms of carbon to each atom of U-235. There is no U-238 present. $\sigma_a = 680$ b for the U-235 and take v to be 2200 m/sec.

Problem 1.2.2

Estimate the period of a reactor if there are no delayed neutrons and if the $k = 1.0010$ and $\ell = 10^{-4}$ sec. The reactor period is the time it takes for the flux or neutron density to increase by a factor of e .

Problem 1.2.3

Let $Y_i = C_i(t) \frac{\Lambda \lambda_i}{\beta_i}$ and $\alpha_R = \beta/\Lambda$. Show that the kinetics equations are then

$$\frac{dn}{dt} = \alpha_R [n(t) (\rho' - 1) + \sum_{i=1}^6 a_i Y_i]$$

and

$$\frac{dY_i}{dt} = \lambda_i [n(t) - Y_i(t)]$$

where $\rho' = \frac{\rho(t)}{\beta}$ and $a_i = \beta_i/\beta$

1.3 Analytical Solutions of the Reactor Kinetics Equations

The kinetics equations are relatively difficult to solve both analytically and numerically due to the large difference in the time constants in the equations as well as due to the fact that there are seven coupled ordinary differential equations if there are six delayed neutron groups.

Equations (1.2.23) and (1.2.24) are repeated here as

$$\frac{dn}{dt} = \frac{\rho(t) - \beta}{\Lambda} n(t) + \sum_{i=1}^6 \lambda_i C_i(t) \quad (1.3.1)$$

and

$$\frac{dC_i(t)}{dt} = \frac{\beta_i n(t)}{\Lambda} - \lambda_i C_i(t) \quad i = 1, 2, \dots, 6. \quad (1.3.2)$$

If we assume that we can take an appropriate average for the delayed neutron decay constant, then it is possible to collapse these seven equations into two equations. For one effective group of delayed neutrons, Equations (1.3.1) and (1.3.2) reduce to

$$\frac{dn(t)}{dt} = \frac{\rho(t) - \beta}{\Lambda} n(t) + \lambda C(t) \quad (1.3.3)$$

and

$$\frac{dC(t)}{dt} = \frac{\beta}{\Lambda} n(t) - \lambda C(t). \quad (1.3.4)$$

The decay constant, λ for this average delayed neutron precursor $C(t)$ is

$$\frac{1}{\lambda} = \frac{1}{\beta} \sum_{i=1}^6 \frac{\beta_i}{\lambda_i} \quad (1.3.5)$$

This average is not unique and at times it may be advantageous to use another expression for the average. If very small reactivities are added to the reactor, then instead of averaging over all six groups, perhaps only the longest three would be used.

If only one group is desired Equation (1.3.5) is used to determine the effective λ and β is simply the sum of the β_i . For two groups, we split the β 's into two groups of three each and

$$\beta_1 = \sum_{i=1}^3 \beta_i, \quad \frac{1}{\lambda_1} = \frac{1}{\beta_1} \sum_{i=1}^3 \frac{\beta_i}{\lambda_i}, \quad (1.3.5a)$$

and

$$\beta_2 = \sum_{i=4}^6 \beta_i, \quad \frac{1}{\lambda_2} = \frac{1}{\beta_2} \sum_{i=4}^6 \frac{\beta_i}{\lambda_i}. \quad (1.3.5b)$$

For three groups, the β 's are split into three groups of two components each and for four groups, the shortest two groups are averaged together as well as the next shortest two and the 22 and 55 sec. groups are treated separately.

We will solve Equations (1.3.3) and (1.3.4) subject to a constant reactivity insertion ρ_0 at time zero. In order to do this we recast our equations into a matrix form because of the similarity to the numerical technique (Hansen's method) used in the solution of the equations. In matrix form, Equations (1.3.3) and (1.3.4) become

$$\frac{d\psi(t)}{dt} = \underline{A} \psi(t) \quad (1.3.6)$$

with

$$\underline{\psi}(t) = \begin{bmatrix} n(t) \\ C(t) \end{bmatrix} \quad \text{and} \quad \underline{A} = \begin{bmatrix} \frac{\rho_0 - \beta}{\Lambda} & \lambda \\ \frac{\beta}{\Lambda} & -\lambda \end{bmatrix}$$

Notice that the \underline{A} matrix is independent of time if the reactivity is a constant. Equation (1.3.6) can formally be integrated to yield

$$\underline{\psi}(t) = e^{\underline{A}t} \underline{\psi}(0). \quad (1.3.7)$$

The initial condition vector $\underline{\psi}(0)$ is

$$\underline{\psi}(0) = \begin{bmatrix} n(0) \\ C(0) \end{bmatrix} = n(0) \begin{bmatrix} 1 \\ \frac{\beta}{\lambda\Lambda} \end{bmatrix} \quad (1.3.8)$$

where we have used Equation (1.3.4) and set the derivative $\frac{dC}{dt} = 0$ for $t \leq 0$.

When reactivity is added at time zero, the delayed neutron precursor concentration $C(t)$ does not change until after some time. It is assumed that the reactor has been operating for a long time at a reactor power consistent with a neutron density $n(0)$.

The formal solution given by Equation (1.3.7) is not much good if the explicit functional dependence of $n(t)$ and $C(t)$ cannot be obtained. The reason no explicit functional dependence is achieved from Equation (1.3.7) is that matrix \underline{A} is not diagonal. This means we really haven't separated the equations from each other. In order to demonstrate how the solution can be obtained, we imitate the technique used for the solution of an

n^{th} order differential equation, i.e. set

$$\underline{\psi}(t) = e^{\omega t} \underline{v}, \quad (1.3.9)$$

where \underline{v} is a constant vector and ω is a scalar independent of time.

Substituting this into Equation (1.3.6) we have

$$\omega e^{\omega t} \underline{v} = \underline{A} e^{\omega t} \underline{v}$$

or

$$\underline{A} \underline{v} = \omega \underline{v}. \quad (1.3.10)$$

Now it is apparent that Equation (1.3.10) is an eigenvalue equation so

$$|\underline{A} - \omega \underline{I}| = 0 \quad (1.3.11)$$

must be satisfied to determine the eigenvalue ω . This equation is called the characteristic equation and in reactor dynamics it is called the "in-hour equation" because it relates the reactivity inserted into \underline{A} with the ω which is the reciprocal of the reactor period.

The inhour equation (Equation 1.3.11) for the one group delayed neutron model is

$$\det \begin{bmatrix} \frac{\rho_0 - \beta}{\Lambda} - \omega & \lambda \\ \frac{\beta}{\Lambda} & -\lambda - \omega \end{bmatrix} = 0 \quad (1.3.12)$$

or

$$\Lambda \omega^2 + (\beta - \rho_0 + \Lambda \omega) - \lambda \rho_0 = 0. \quad (1.3.13)$$

Equation (1.3.13) can also be written as

$$\rho_0 = \Lambda \omega + \frac{\beta \omega}{\omega + \lambda} \quad (1.3.14)$$

If we had kept all six groups of delayed neutrons, Equation (1.3.11) would be a 7 x 7 determinant and Equation (1.3.14) would be enlarged to

$$\rho_0 = \Lambda \omega + \sum_{i=1}^6 \frac{\beta_i \omega}{\omega + \lambda_i} \quad (1.3.15)$$

The \underline{v} of Equation (1.3.10) is the eigenvector associated with the eigenvalue ω . Now if we perform a similarity transformation \underline{B} on \underline{A} , then both \underline{A} and $\underline{B}^{-1} \underline{A} \underline{B}$ have the same characteristic equation. Also the fact that the trace of \underline{A} is the sum of the eigenvalues of the matrix \underline{A} is a useful check on the actual calculation of the eigenvalues of \underline{A} . From this last fact it is apparent that the sum of the two roots ω_1 and ω_2 (the two eigenvalues) of Equation (1.3.13) is given by the relation

$$\omega_1 + \omega_2 = \frac{\rho_0 - \beta}{\Lambda} - \lambda \quad (1.3.16)$$

The eigenvectors \underline{v}_1 and \underline{v}_2 associated with ω_1 and ω_2 respectively are obtained by taking the cofactors of the element in any row of Equation (1.3.12). To see how this works, we set

$$\underline{v}_1 = \begin{bmatrix} v_{11} \\ v_{21} \end{bmatrix} \quad \text{and} \quad \underline{v}_2 = \begin{bmatrix} v_{12} \\ v_{22} \end{bmatrix}$$

and have v_{11} , the cofactor of the element $a_{11} - \omega_1$, as

$$v_{11} = -(\lambda + \omega_1),$$

and the cofactor of the element λ of the first row as

$$v_{21} = -\beta/\Lambda.$$

Similarly, for the vector \underline{v}_2 , we have for the second row of $|\underline{A} - \underline{I}\omega|$,

$$v_{12} = -\lambda, \quad v_{22} = \frac{\rho_0 - \beta}{\Lambda} - \omega_2.$$

Notice that to get the \underline{v}_1 , we use the eigenvalue ω_1 and for \underline{v}_2 , we use the eigenvalue ω_2 in Equation (1.3.12). Therefore, the eigenvectors of \underline{A} are (to within a constant)

$$\underline{v}_1 = \begin{bmatrix} \lambda + \omega_1 \\ \beta/\Lambda \end{bmatrix} \quad \text{and} \quad \underline{v}_2 = \begin{bmatrix} \lambda \\ \frac{\beta - \rho_0}{\Lambda} + \omega_2 \end{bmatrix}$$

These vectors \underline{v}_1 and \underline{v}_2 are linearly independent since the eigenvalues are distinct.

Now since we have

$$\underline{A} \underline{v}_i = \omega_i \underline{v}_i \quad \text{for } i = 1, 2,$$

it is apparent that the similarity transformation

$$\underline{A} \underline{B} = \underline{B} \underline{D} \tag{1.3.17}$$

holds, where \underline{D} is a diagonal matrix having the eigenvalues as its elements,

i.e.

$$\underline{\underline{D}} = \begin{bmatrix} \omega_1 & 0 \\ 0 & \omega_2 \end{bmatrix}$$

and

$$\underline{\underline{B}} = \begin{bmatrix} \underline{v}_1 & \underline{v}_2 \end{bmatrix}$$

These two matrices are now completely determined for the one delayed neutron model.

If we make a transformation

$$\underline{\psi}(t) = \underline{\underline{B}} \underline{\underline{Z}}(t), \quad (1.3.18)$$

then

$$\begin{aligned} \frac{d\underline{\underline{Z}}(t)}{dt} &= \underline{\underline{B}}^{-1} \underline{\underline{A}} \underline{\underline{B}} \underline{\underline{Z}}(t) \\ &= \underline{\underline{D}} \underline{\underline{Z}}(t) \end{aligned}$$

and its solution is obviously

$$\underline{\underline{Z}}(t) = \begin{bmatrix} e^{\omega_1 t} & 0 \\ 0 & e^{\omega_2 t} \end{bmatrix} \underline{\underline{Z}}(0) = e^{\underline{\underline{D}} t} \underline{\underline{Z}}(0) \quad (1.3.19)$$

Now using Equations (1.3.18) and (1.3.19) we have

$$\underline{\psi}(t) = \underline{\underline{B}} e^{\underline{\underline{D}} t} \underline{\underline{B}}^{-1} \underline{\psi}(0) \quad (1.3.20)$$

This is the solution of the kinetics equations. Notice now that the exponential

$e^{\frac{D}{dt}}$ is a diagonal matrix so the separation of the equations is effected.

We now write out the details of this procedure. The transformation matrix \underline{B} is

$$\underline{B} = \begin{bmatrix} v_1 & v_2 \end{bmatrix} = \begin{bmatrix} \lambda + \omega_1 & \lambda \\ \beta/\Lambda & \frac{\beta - \rho_0}{\Lambda} + \omega_2 \end{bmatrix}$$

and the determinant of \underline{B} is

$$\det \underline{B} = \omega_1 \omega_2 + \frac{\beta - \rho_0}{\Lambda} \omega_1 + \lambda \omega_2 - \frac{\lambda \rho_0}{\Lambda} \quad (1.3.21)$$

The inverse of \underline{B} is

$$\underline{B}^{-1} = \frac{1}{\det \underline{B}} \begin{bmatrix} \frac{\beta - \rho_0}{\Lambda} + \omega_2 & -\lambda \\ -\frac{\beta}{\Lambda} & \lambda + \omega_1 \end{bmatrix}$$

The solution of the one delayed neutron group equations is obtained from Equations (1.3.20) and (1.3.8)

$$\underline{\psi}(t) = \frac{n(0)}{\det \underline{B}} \begin{bmatrix} \lambda + \omega_1 & \lambda \\ \frac{\beta}{\Lambda} & \frac{\beta - \rho_0}{\Lambda} + \omega_2 \end{bmatrix}$$

$$x \begin{bmatrix} e^{\omega_1 t} & 0 \\ 0 & e^{\omega_2 t} \end{bmatrix}$$

$$x \begin{bmatrix} \frac{\beta - \rho_o}{\Lambda} + \omega_2 & -\lambda \\ -\frac{\beta}{\Lambda} & \lambda + \omega_1 \end{bmatrix}$$

$$x \begin{bmatrix} 1 \\ \frac{\beta}{\Lambda \lambda} \end{bmatrix}$$

$$= \frac{n(t)}{\text{Det } B} \begin{bmatrix} (\omega_2 - \frac{\rho_o}{\Lambda})(\lambda + \omega_1) e^{\omega_1 t} + \frac{\omega_1 \beta}{\Lambda} e^{\omega_2 t} \\ \frac{\beta}{\Lambda} (\omega_2 - \frac{\rho_o}{\Lambda}) e^{\omega_1 t} + (\frac{\beta - \rho_o}{\Lambda} + \omega_2) \frac{\omega_1 \beta}{\Lambda \lambda} e^{\omega_2 t} \end{bmatrix} \quad (1.3.23)$$

Even though Equations (1.3.23) represent the exact solution of the problem, we are generally not concerned with $c(t)$ so we look only at the $n(t)$ equation. Also from Equation (1.3.13) we have

$$\omega_{1,2} = \frac{-(\beta - \rho_o + \Lambda \lambda) \pm \sqrt{(\beta - \rho_o + \Lambda \lambda)^2 + 4 \Lambda \lambda \rho_o}}{2 \Lambda} \quad (1.3.24)$$

Note too that

$$\omega_1 \omega_2 = - \frac{\lambda \rho_o}{\Lambda} \quad (1.3.25)$$

and in agreement with Equation (1.3.16),

$$\omega_1 + \omega_2 = - \frac{\beta - \rho_o + \lambda \Lambda}{\Lambda} \quad (1.3.26)$$

Now if we approximate ω_1 (the positive sign) and ω_2 (the negative sign) in Equation (1.3.24) by assuming that

$$(\beta - \rho_o + \lambda \Lambda)^2 \gg 4 \lambda \rho_o \Lambda,$$

which is true for most reactors, we have the radical of Equation (1.3.24) equal to

$$(\beta - \rho_o + \lambda \Lambda) \sqrt{1 + \frac{4 \lambda \rho_o \Lambda}{(\beta - \rho_o + \lambda \Lambda)^2}}$$

or

$$\left[1 + \frac{1}{2} \frac{4 \lambda \rho_o \Lambda}{(\beta - \rho_o + \lambda \Lambda)^2} + \dots \right] (\beta - \rho_o + \lambda \Lambda)$$

When this is used we obtain

$$\omega_1 \approx \frac{\lambda \rho_o}{\beta - \rho_o} \quad \text{and} \quad \omega_2 \approx - \frac{\beta - \rho_o}{\Lambda} \quad (1.3.27)$$

Note that if ρ_o is positive, then ω_1 is also positive.

The determinant of \underline{B} is involved in the solution of $n(t)$ so with the above approximations as well as by using Equation (1.3.25) we have

$$\det \underline{\underline{B}} \doteq -\frac{\lambda \rho_o}{\Lambda} + \frac{\beta - \rho_o}{\Lambda} \cdot \frac{\lambda \rho_o}{\beta - \rho_o} - \frac{\lambda}{\Lambda} (\beta - \rho_o) - \frac{\lambda \rho_o}{\Lambda} = -\frac{\lambda \beta}{\Lambda} \quad (1.3.28)$$

So Equation (1.3.23) becomes, using Equation (1.3.28),

$$n(t) = n(o) \left[\frac{\beta}{\beta - \rho_o} e^{\frac{\lambda \rho_o}{\beta - \rho_o} t} - \frac{\rho_o}{\beta - \rho_o} e^{-\frac{\beta - \rho_o}{\Lambda} t} \right] \quad (1.3.29)$$

This approximate solution is very useful in obtaining checks on numerical solutions. The character of the solution is also exponential. This will play a role in our numerical technique.

Problem 1.3.1

Assume that $\rho(t) = 0$ in Equations (1.3.3) and (1.3.4) and obtain a solution of the kinetics equations.

Problem 1.3.2

Prove the theorems

a) $|\underline{\underline{B}}^{-1} \underline{\underline{A}} \underline{\underline{B}} - \omega \underline{\underline{I}}| = |\underline{\underline{A}} - \omega \underline{\underline{I}}|$

and

b) $\text{trace } \underline{\underline{A}} \doteq \sum_{i=1}^2 \omega_i$

Problem 1.3.3

Why does Equation (1.3.17) follow from the eigenvector equation

$$\underline{\underline{A}} \underline{\underline{v}} = \omega \underline{\underline{v}}?$$

Problem 1.3.4

Using the same approximations as used in developing Equation (1.3.29) establish a relation for $c(t)$.

1.4 Numerical Solution of the Kinetics Equations

The reactor kinetics equations are difficult to solve numerically by "standard" Runge-Kutta or predictor-corrector methods. The basic reason becomes apparent by looking at the one-delayed neutron group equations which we repeat as

$$\frac{dn(t)}{dt} = \frac{\rho(t) - \beta}{\Lambda} n(t) + \lambda c(t) \quad (1.4.1)$$

and

$$\frac{dc(t)}{dt} = \frac{\beta}{\Lambda} n(t) - \lambda c(t) \quad (1.4.2)$$

The very short time response Λ of the prompt neutrons is of the order of 10^{-4} sec. whereas the delayed neutron time response is $\frac{1}{\lambda}$ or about 10 sec, a factor of 10^5 greater. The implication of these facts is that in order to obtain the prompt response, very small time steps, of the order of 10^{-4} sec, are required. But then before the delayed neutron term can come into play, many time steps are required. Also, to examine the response out to even one second, 10,000 steps of calculation would be required.

There are several methods that are used to ameliorate this difficult problem:

- 1) Using a Laplace transform technique.
- 2) Transferring the differential equations to integral equations.
- 3) Using the eigenvalue method.

We choose the last technique and refer to it as Hansen's method (4) after its originator. The method works for varying reactivity and can be extended

to systems with feedback.

The basic idea of Hansen's method is relatively simple. We again write Equations (1.4.1) and (1.4.2) as a matrix and set

$$\frac{d\psi(t)}{dt} = \underline{A} \psi \quad (1.4.3)$$

where \underline{A} and ψ are defined as in Section 3 of this module. We will also only perform the operations for our one delayed neutron group model.

Now set

$$\underline{A} = \underline{L} + \underline{D} + \underline{U} \quad (1.4.4)$$

where

$$\underline{L} = \begin{bmatrix} 0 & 0 \\ \frac{\beta}{\Lambda} & 0 \end{bmatrix}, \quad \underline{D} = \begin{bmatrix} \frac{\rho - \beta}{\Lambda} & 0 \\ 0 & -\lambda \end{bmatrix}$$

and

$$\underline{U} = \begin{bmatrix} 0 & \lambda \\ 0 & 0 \end{bmatrix}$$

Of course, for all six delayed neutron groups these matrices would still have the same meaning. Equation (1.4.4) can be written using these definitions as

$$\frac{d\psi(t)}{dt} = \underline{D} \psi(t) = (\underline{L} + \underline{U}) \psi(t) \quad (1.4.5)$$

Equation (1.4.5) does not appear any simpler to solve than Equation (1.4.3) and in fact it isn't. The reason for splitting it up in this fashion is to develop an iteration procedure. We assume we begin this calculation from a time t_0 and advance to a time t_1 . We set

$$h = t_1 - t_0. \quad (1.4.6)$$

Since \underline{D} is a diagonal matrix, an integrating factor for Equation (1.4.5) is $e^{-\underline{D}t}$ if the reactivity doesn't change much during the time interval h . Therefore, Equation (1.4.5) becomes

$$e^{-\underline{D}t} \frac{d\underline{\psi}(t)}{dt} - e^{-\underline{D}t} \underline{D} \underline{\psi}(t) = e^{-\underline{D}t} (\underline{L} + \underline{U}) \underline{\psi}(t)$$

or

$$\frac{d}{dt} (e^{-\underline{D}t} \underline{\psi}(t)) = e^{-\underline{D}t} (\underline{L} + \underline{U}) \underline{\psi}(t) \quad (1.4.6)$$

Integrating now from 0 to h we have

$$e^{-\underline{D}t} \underline{\psi}(t) \Big|_0^h = \int_0^h e^{-\underline{D}t} (\underline{L} + \underline{U}) \underline{\psi}(t) dt$$

or

$$\underline{\psi}(t_0 + h) = e^{\underline{D}h} \underline{\psi}(t_0) + \int_0^h e^{\underline{D}(h-\theta)} (\underline{L} + \underline{U}) \underline{\psi}(t_0 + \theta) d\theta, \quad (1.4.7)$$

where

$$t_0 \leq \theta \leq t_1 = t_0 + h,$$

and obviously in this interval

$$d\theta = dt.$$

This is an integral equation since the function we're looking for is part of the integrand. In order to provide a reasonable approximation to $\psi(t_0 + \theta)$, we recall that the analytical solution is exponential so we assume

$$\psi(t_0 + \theta) = e^{\omega_0 \theta} \psi(t_0), \quad (1.4.8)$$

and ω_0 is the largest eigenvalue of the matrix \underline{A} . This means that we will have to solve the equation (notice this is simply the inhour equation)

$$|\underline{A} - \omega \underline{I}| = 0 \quad (1.4.9)$$

in each time interval for which the reactivity has changed since the reactivity will generally be a time dependent quantity.

Inserting Equation (1.4.8) into (1.4.7) we have

$$\begin{aligned} \psi(t_0 + h) &= e^{\underline{D}h} \psi(t_0) + \int_0^h e^{\underline{D}(h-\theta)} (\underline{L} + \underline{U}) e^{\omega_0 \theta} \psi(t_0) d\theta \\ &= e^{\underline{D}h} \psi(t_0) + \left[\int_0^h e^{\underline{D}(h-\theta)} e^{\omega_0 \underline{I}\theta} d\theta \right] (\underline{L} + \underline{U}) \psi(t_0) \\ &= e^{\underline{D}h} \psi(t_0) + \left[e^{\underline{D}h} \cdot \int_0^h e^{(\omega_0 \underline{I} - \underline{D})\theta} d\theta \right] (\underline{L} + \underline{U}) \psi(t_0) \\ &= e^{\underline{D}h} \psi(t_0) + (\omega_0 \underline{I} - \underline{D})^{-1} \left[e^{\omega_0 \underline{I}h} - e^{\underline{D}h} \right] (\underline{L} + \underline{U}) \psi(t_0). \end{aligned} \quad (1.4.10)$$

If we write

$$\psi(t_0) = \psi_j \quad (1.4.11)$$

and

$$\psi(t_0 + h) = \psi_{j+1}$$

Then Equation (1.4.10) becomes

$$\psi_{j+1} = e^{\underline{D}h} \psi_j + (\omega_{0,\underline{D}})^{-1} \left[e^{\omega_0 h \underline{I}} - e^{\underline{D}h} \right] (\underline{L} + \underline{U}) \psi_j \quad (1.4.11)$$

$$\equiv \underline{G} \psi_j \quad (1.4.12)$$

This \underline{G} matrix obviously represents

$$\underline{G} = e^{\underline{D}h} + (\omega_{0,\underline{D}})^{-1} \left[e^{\omega_0 h \underline{I}} - e^{\underline{D}h} \right] (\underline{L} + \underline{U}), \quad (1.4.13)$$

and it can be written as

$$\underline{G} = \begin{bmatrix} e^{\frac{\rho-\beta}{\Lambda} h} & \frac{e^{\omega_0 h} - e^{\frac{\rho-\beta}{\Lambda} h}}{\omega_0 - \frac{\rho-\beta}{\Lambda}} \lambda \\ \frac{e^{\omega_0 h} - e^{-\lambda h}}{\omega_0 + \lambda_1} \frac{\beta}{\Lambda} & e^{-\lambda h} \end{bmatrix} \quad (1.4.14)$$

For completeness, if there are N delayed neutron groups, we include the following relation for \underline{G} ,

$$\underline{G} = \begin{bmatrix} \frac{\rho - \beta}{\Lambda} h & \frac{e^{\omega_0 h} - e^{\frac{\rho - \beta}{\Lambda} h}}{\omega_0 - (\frac{\rho - \beta}{\Lambda})} & \lambda_1 & \dots & \frac{e^{\omega_0 h} - e^{\frac{\rho - \beta}{\Lambda} h}}{\omega_0 - (\frac{\rho - \beta}{\Lambda})} & \lambda_N \\ \frac{e^{\omega_0 h} - e^{-\lambda_1 h}}{\omega_0 + \lambda_1} \frac{\beta_1}{\Lambda} & e^{-\lambda_1 h} & \dots & \dots & 0 & \dots \\ \vdots & \vdots & \ddots & \ddots & \vdots & \vdots \\ \frac{e^{\omega_0 h} - e^{-\lambda_N h}}{\omega_0 + \lambda_N} \frac{\beta_N}{\Lambda} & \dots & \dots & \dots & e^{-\lambda_N h} & \dots \end{bmatrix} \quad (1.4.15)$$

If there are six delayed neutron groups, it is apparent that \underline{G} will be a 7 x 7 matrix. Also this iteration technique as given by Equation (1.4.12) is unconditionally stable and yields the asymptotically correct eigen-solution.

The numerical procedure for the solution of the kinetics equations is thus:

1. Determine the number of delayed neutron groups desired and read in the pertinent parameters such as Λ_1 , β_1 etc. Also choose a time step "h".
2. Construct the vector $\underline{\psi}(0)$. This will usually be

$$\underline{\psi}(0) = n(0) \begin{bmatrix} 1 \\ \beta_1 / \lambda_1 \Lambda \\ \vdots \\ \beta_N / \lambda_N \Lambda \end{bmatrix}$$

3. Determine the largest eigenvalue of the equation

$$|\underline{A} - \omega \underline{I}| = 0.$$

This is a rather difficult step since it means solving an algebraic equation of perhaps degree 7 to determine its largest root ω_0 .

4. Construct the \underline{G} matrix using Equation (1.4.15).
5. Determine the vector $\underline{\psi}_1$ where

$$\underline{\psi}_1(h) = \underline{G} \underline{\psi}(0).$$

6. Repeat the above steps starting with step 3.

The technique is not involved and can yield very accurate results.

The determination of the root of an algebraic equation needs some discussion. We choose the Newton-Raphson technique to solve the equation

$$f(\omega) = \sum_{n=0}^N a_n \omega^n = 0 \quad \text{for } 1 \leq N \leq 7. \quad (1.4.16)$$

Let the N roots of Equation (1.4.17) be labelled $\omega_0, \omega_1, \dots, \omega_N$ where the roots are ordered such that $\omega_0 > \omega_1 > \omega_2 > \dots > \omega_N$. For our problem, all the roots are real and there will only be one positive root depending on whether ρ is positive at the time step of interest. We are interested only in ω_0 . Also we assume that the interval of interest in the roots is limited by

$$\omega_0 \leq \left| \frac{\rho_0}{A} \right|$$

and

$$|\omega_N| \leq |\lambda_N|,$$

with λ_N the decay constant corresponding to the shortest lived delayed neutron group. The reactivity can range from negative β to positive β in

most practical situations. Note that if

$$\rho = 0 \text{ then } \omega_0 = 0$$

and if $\rho = \beta$, then $\omega_0 \rightarrow \infty$ while if $\rho = -\beta$, then $\omega_0 \rightarrow -\lambda_1$. For the last situation, regardless of the amount of negative reactivity introduced into the reactor, the reactor cannot shut down faster than a period of

$$T = \omega_0^{-1} = \frac{1}{\lambda_1} \approx 80 \text{ sec.}$$

The Newton-Raphson method is relatively simple to use. Assume that we can expand $f(\omega)$ in a Taylor series about ω_0 , where ω_0 is the root of interest. Then

$$f(\omega_0) = f(\omega_{01}) + h \frac{df(\omega_{01})}{d\omega} + \frac{h^2}{2} \frac{d^2 f(\omega_{01})}{d\omega^2} + \dots \quad (1.4.17)$$

The ω_{01} is a first "guess" at the solution which we assume to be $\rho(t)/\Lambda$. If we truncate Equation (1.4.17) after the first two terms on the right, we have

$$f(\omega_0) \approx f(\omega_{01}) + h \frac{df}{d\omega}(\omega_{01}) = 0$$

or

$$h = - \frac{f(\omega_{01})}{\frac{df(\omega_{01})}{d\omega}} \quad (1.4.18)$$

The next approximation to ω_0 is then

$$\omega_{02} = \omega_{01} + h \quad (1.4.19)$$

and then Equation (1.4.18) is used again. If this iteration procedure is used a sufficient number of times, then the roots of Equation (1.4.16) can be obtained.

Problem 1.4.1

- Use the Newton-Raphson method to solve the equation

$$x^2 + 3x - 8 = 0.$$

Check by the quadratic formula.

Problem 1.4.2

Extend the Newton-Raphson method to systems of equations. In particular if

$$f_1(x_1, x_2) = 0$$

and

$$f_2(x_1, x_2) = 0$$

then show that if h_1 and h_2 are the increments to the assumed roots,

$$h_1 = \frac{f_2 \frac{\partial f_1}{\partial x_2} - f_1 \frac{\partial f_2}{\partial x_2}}{\det \underline{J}}$$

and

$$h_2 = \frac{f_1 \frac{\partial f_2}{\partial x_2} - f_2 \frac{\partial f_1}{\partial x_1}}{\det \underline{J}}$$

where \underline{J} is the Jacobian matrix, i.e.

$$J = \begin{bmatrix} \frac{\partial f_1}{\partial x_1} & \frac{\partial f_1}{\partial x_2} \\ \frac{\partial f_2}{\partial x_1} & \frac{\partial f_2}{\partial x_2} \end{bmatrix}$$

and $\det J$ is the determinant of the Jacobian matrix.

Using this formulation find the solution to the equations

$$f_1(x_1, x_2) = \sin x_1 x_2 - x_1 + x_2^2 = 0$$

$$f_2(x_1, x_2) = 2x_1^2 - x_2^2 + x_1 x_2 = 0.$$

1.5 The Computer Program and the Kinetics Equations

As has already been indicated, the kinetics equations offer a challenge for their successful computer solution. The purpose of this section is to indicate what models the program will solve. The program will solve the kinetics equations by means of Hansen's method for the following reactivity inputs:

- 1) A constant reactivity,

$$\rho(t) = \rho_0. \quad (1.5.1)$$

- 2) Sinsusoidal variation of reactivity,

$$\rho(t) = \rho_0 \sin b_2 t, \quad (1.5.2)$$

where ρ_0 and b_2 must be input.

- 3) Linear reactivity insertion to approximate the insertion or withdrawal of a control rod:

$$\rho(t) = \rho_0(1 + b_3 t). \quad (1.5.3)$$

These inputs are the most common and can be used to simulate many situations. In each of the above situations ρ_0 must be read into the computer as well as the "b" or rate of insertion in sec^{-1} . The reactivity units are in dollars. A dollar of reactivity is the amount inserted or withdrawn which equals the delayed neutron fraction. For example, a reactivity of 50¢ for U-235 where $\beta = 0.0065$ is 0.00325 whereas for Pu-239, it is 0.00135 since $\beta = 0.0027$.

The program also has an option of either doing the calculations for

U-235 or Pu-239. The delayed neutron decay constants are taken from Keepins' data and are given in Table 1.5.1.

Table 1.5.1 Delayed-Neutron Half-Lives, Decay Constants and Yields from Thermal Fission of U-235 and Pu-239.

URANIUM-235				PLUTONIUM-239		
Group Index	Half-Life (sec)	Decay Constant λ (sec ⁻¹)	Relative Abundance β_i/β	Half-Life (sec)	Decay Constant λ (sec ⁻¹)	Relative Abundance β_i/β
1	55.72	0.0124	0.033	54.28	0.0128	0.035
2	22.72	0.0305	0.219	23.04	0.0301	0.298
3	6.22	0.111	0.196	5.60	0.124	0.211
4	2.30	0.301	0.395	2.13	0.325	0.326
5	0.610	1.14	0.115	0.618	1.12	0.086
6	0.230	3.01	0.042	0.257	2.69	0.044
		$\beta=0.0065$			$\beta=0.0027$	

The computer program for this module is good only for low power reactors since feedback effects are not taken into account. It does give some idea of the behavior of the power of the reactor for various reactivities as well as for the two different fuels.

I.6 Input-Output Data for Code FUMOKI

The input data required for the program are presented below:

Card 1. IFUEL indicates the type of fuel the reactor has. It can either be ^{235}U or ^{239}Pu . If IFUEL = 5 - its Uranium -235

IFUEL = 9 - the fuel is Plutonium 239.

The format is I2.

Card 2. NN - The number of groups of delayed neutrons desired. NN can be any integer from 1 to 6.

The format is again I2.

Card 3. NRO - Type of reactivity inserted into the reactor.

The format is I2.

If NRO = 1; $\rho(t) = \rho_0$,

NRO = 2; $\rho(t) = \rho_0 \sin b_2 t$,

NRO = 3; $\rho(t) = \rho_0 (1 + b_3 t)$.

Card 3 must have the reactivity ρ_0 , in dollars, read in with an F 12.9 format. The b_2 and b_3 are read in with units of sec^{-1} .

The b's are read with an F 12.9 format on this card.

Card 4. XL - Neutron generation time.

XL must be in seconds and is in an F 12.9 format.

Card 5. TIME: - the total time period for which the solution is desired (seconds)

H - The time step desired, again in seconds. This cannot be taken arbitrarily large. Generally one should choose an H of between 0.005 and 0.05 sec. Both of these are F 10.5 formats.

✓

XXN can be anything you choose but it is in an F 10.5 format.

An example of an input data set is given below. This is an example where we use Uranium fuel, with one group of delayed neutrons for a constant reactivity input of about 34¢.

VIRGINIA TECH

15117435

The core requirements for FUMOKI are about 40 K bytes and the compilation time increases approximately with the square of the number of delayed neutron groups, all other parameters being the same.

The output of the program is simple. The time, neutron density and G matrix can be read out. Also a plot of the density versus time is printed out.

Problem 1.6.1

Run the sample problem, i.e. the data shown.

Problem 1.6.2

Find the neutron density if you start with 100 neutrons/cm³ for a ²³⁹Pu reactor with all six delayed neutron groups. Obtain the solution for a 5 second time span and do the same problem with H = 0.01 and 0.1 sec. Compare the results. Do not write the G matrix.

Problem 1.6.3

Run FUMOKI for the case for three delayed neutron groups for ²³⁹Pu, with a neutron generation time of 10⁻⁴ sec and the reactivity

$$\rho(t) = \$0.25 \sin 5t,$$

Run it for a total time of 1 second, time intervals of 0.01 sec and an initial relative flux of 1.000. Notice that the reactivity at 0.5 sec is \$0.17 and the power has risen to 1.275.

If you do the exact same problem as above for 2 delayed neutron groups notice that the power has become 1.268 at 0.5 sec.

Problem 1.6.4

Given a reactor fueled with ^{239}Pu and assume that you wish to use a 5-delayed-neutron-group model. For a relative power of 1.000 initially, show that after one second the power is 2.157 if a rod is being removed such that the reactivity inserted is

$$\rho(t) = \rho_0(1 + 5t).$$

All other parameters are as in problem 1.6.3.

Problem 1.6.5

Assume a ^{235}U fueled system with a reactivity variation of

$$\rho(t) = 0.25 \sin 5t.$$

If the reactor is to be simulated for three seconds after the reactivity is inserted and, if all parameters are the same as in example 1.6.3 except for the number of delayed neutron groups then:

a) use the two delayed neutron model to obtain the power as a function of time

and

b) use the five delayed neutron model to obtain the power variation with time.

How much longer did the computer take to do case b) than case a)?

Answer: About 10 times as long. On the IBM 370, the times were 11.03 sec and .97 sec respectively. Notice also that the power at the end of 3.0 sec is 0.8154 and 0.8161 for the 2 and 5 group cases respectively.

REFERENCES

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List of Symbols for FUMOKI

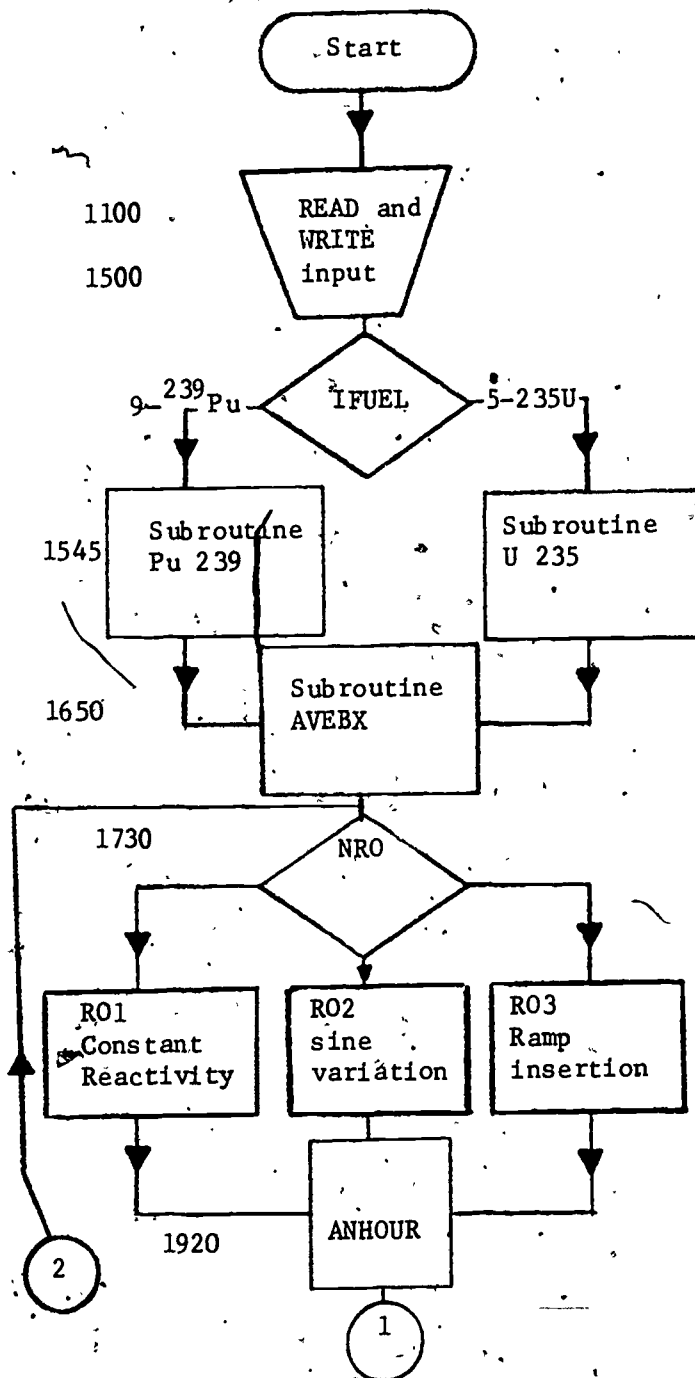
The following symbols are listed in the order of their appearance in program FUMOKI.

IFUEL		Type of fuel (^{235}U or ^{239}Pu) (5 or 9)
NN		Number of delayed neutron groups
NRO	$\rho(t) = \rho_0 \sin b_2 t$ $\rho_0(1 + b_3 t)$	Type of reactivity to be inserted. 1 = constant reactivity 2 = sine insertion 3 = ramp insertion
RO	ρ_0	Reactivity inserted
B1,B2,B3	b_1, b_2, b_3	Rate or period at which reactivity is inserted.
X(I)	λ	precursor decay constants
XL	Λ	neutron generation time
TIME		total time the reactor transient (s) is to be simulated
H	Δt	the time step
IG		output option to print the G matrix after 1st iteration (to print any other iteration change statement MKI 2520)
XN	n_0	initial relative power
B(I)		fraction of neutrons for each group
BB		total delayed neutron fraction
TH		cumulative time used in TPLOT
LL		total number of time steps H fits into TIME
R	$\rho(t)$	Reactivity at any time
MM		number of coefficients in the inhour equation. (number of delayed neutron groups + 2)

AW(I)		coefficients of the inhour polynomial
M		degree of the inhour polynomial (number of delayed neutron groups + 1)
WO	ω_0	largest eigenvalue of $ \underline{A} - \omega \underline{I} = 0$
XNCO(I), GNXO(I)	$\underline{\phi}$	column vector $\underline{\phi}$
XXNC2		logarithm of power
GG,MAD		power
MOD(I)	H*(step number)	dimensioned TH used in TPLOT
Z1(I)		precursor density of delayed neutron group 1
Z2(I)		reactivity
AVEBX		subroutine to obtain the average β and λ depending on how many delayed neutron groups are desired
ANHOUR		subroutine to calculate the coefficients of the inhour equation
GMTRX		a subroutine which forms the \underline{G} matrix
GXM		subroutine which multiplies \underline{G} by column vector $\underline{\phi}$, i.e. $\underline{G}\underline{\phi}_0$
POLRT		A Newton-Raphson iteration technique is used to obtain the roots of the inhour equation to obtain ω_0
TPLOT		A subroutine to plot the vector $\underline{\phi}$ up to and including five dependent variables. The reactivity, logarithm of the power, power and some of the delayed neutron precursors are plotted.

Flow Chart for FUMOKI

Statement Numbers



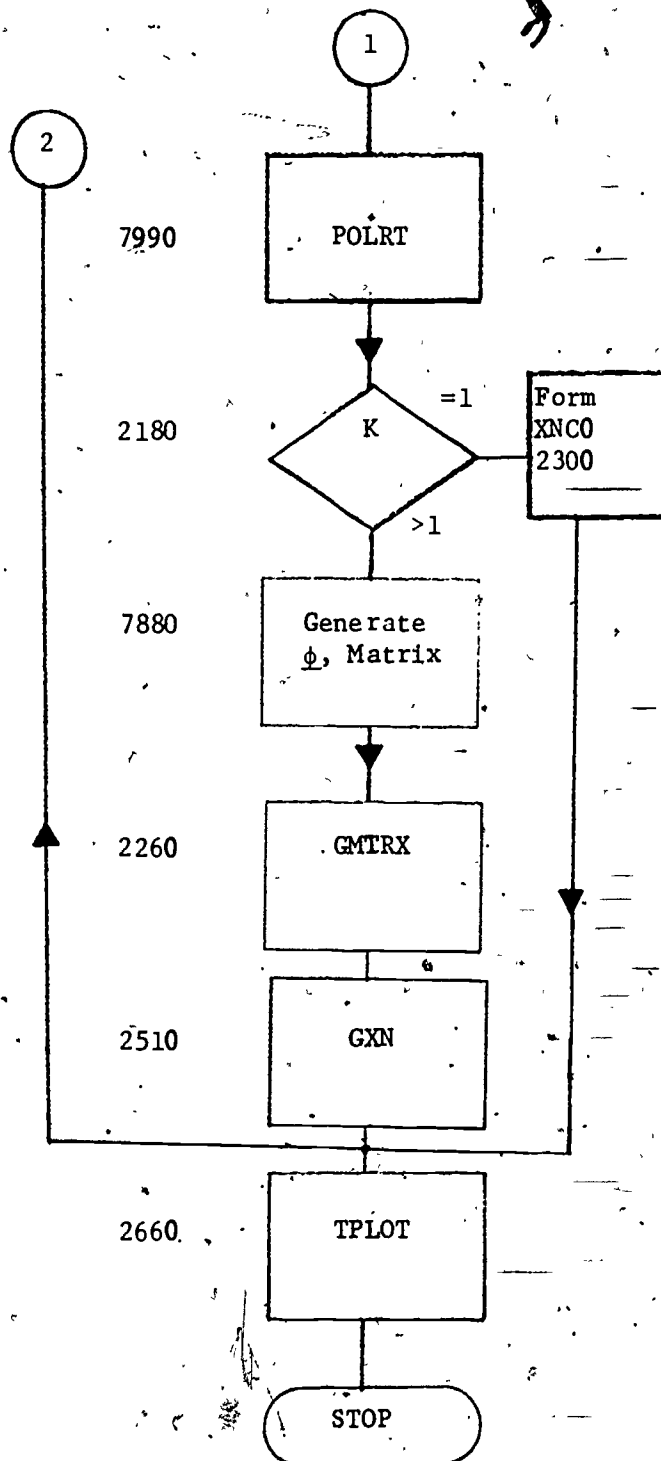
Obtains six group β 's and λ 's for isotope.

Calculates average β 's and λ 's depending on the number of delayed groups desired.

The type of reactivity variation desired determines which path to take.

Routines which determine the reactivity as functions of time.

Subroutine which calculates the polynomial form of the inhour equation.



Subroutine to determine the largest root of the inhour equation using Newton-Raphson method.

Subroutine which forms initial ϕ matrix; ϕ_0

Generates the ϕ matrix for all cases except the first

Forms the \underline{G} matrix

Forms the product $\underline{G} \phi$

Plots the power, etc.

//WATFIV SMART, GNEGA, PAGES=80, TIME=150
MAIN PROGRAM

KINETICS MODULE

CODE NAME F U M K I
OBJECTIVE TO OBTAIN THE KINETICS EQUATION WITHOUT FEEDBACK AND
SOLVE THE KINETICS EQUATIONS NUMERICALLY FOR ONE, TO
SIX DELAYED NEUTRON GROUPS FOR TIME VARYING REACTIVITY
INSERTIONS.
THE KINETIC EQUATION DERIVED FROM ONE GROUP DIFFUSION
EQUATIONS.
FEEDBACK EFFECTS AND SPATIAL DEPENDENT EFFECTS ARE NOT
INCLUDED.

ASSUMPTION THE BASIC ASSUMPTION OF THE METHOD IS THAT NEUTRON AND
PRECURSOR DENSITIES BEHAVE EXPONENTIALLY WITH A FREQUENCY
CHARACTERISTICS OF THE ASYMPTOTIC FREQUENCY CORRESPOND-
ING TO THE REACTIVITY.

PROGRAM WRITTEN IN SINGLE PRECISION

GENERAL DESCRIPTION OF PARAMETERS

SYMBOL	IN/OUT/V	DESCRIPTION	REAL/INT.	UNIT
IFUEL	IN	TYPE OF FUEL 5 = U-235 9 = PU-239	I	-
NN	IN	GROUP DELAYED NEUTRON FROM 1 TO 6	I	-
NRO	IN	TYPE OF REACTIVITY 1= CONSTANT 2= FSIN(R2*T) 3= R(1+P1*T)	I	-
B1	IN	CONSTANT B2 OR B3 IN NRO	R	1/SEC
B	OUT	FRACTION OF DELAYED NEUTRON	R	-
X	OUT	PRECURSOR DECAY CONSTANT	R	1/SEC
XL	IN	NEUTRON GENERATION TIME	R	SEC
R	V	REACTIVITY	R	-
A	V	COEFFICIENT OF INHOUR FORMULA IN POLYNOMIAL FORM	R	-
WO	V	RECIPROCAL OF PERIOD	R	1/SEC
T	V	STABLE PERIOD	R	SEC
G	V	G - MATRIX	R	-
XN	V	NEUTRON DENSITY	R	N/CM**3
C	V	PRECURSOR DENSITY	R	N/CM**3
TIME	N	TOTAL TIME	R	SEC
H	N	TIME INCREMENT	R	SEC

IN - INPUT
OUT - OUTPUT

MK1 20
MK1 30
MK1 40
MK1 50
MK1 60
MK1 70
MK1 80
MK1 90
MK1 100
MK1 110
MK1 120
MK1 130
MK1 140
MK1 150
MK1 160
MK1 170
MK1 180
MK1 190
MK1 200
MK1 210
MK1 220
MK1 230
MK1 240
MK1 250
MK1 260
MK1 270
MK1 280
MK1 290
MK1 300
MK1 310
MK1 320
MK1 330
MK1 340
MK1 350
MK1 360
MK1 370
MK1 380
MK1 390
MK1 400
MK1 410
MK1 420
MK1 430
MK1 440
MK1 450
MK1 460
MK1 470
MK1 480
MK1 490
MK1 500
MK1 510
MK1 520
MK1 530
MK1 540
MK1 550
MK1 560
MK1 570
MK1 580
MK1 590
MK1 600


```

15      WRITE (6,40) XL                      MK1 1210
16      WRITE (6,50) NKL,PF,NPD,B1          MK1 1220
17      WRITE (6,60) TIME,H                 MK1 1230
18      WRITE (6,70) IG,XN                  MK1 1240
19      WRITE (6,100)                        MK1 1250
20  20  FORMAT (1H1,6H*****KINETICS MODULE 1**MK1 1260
1*****//,20H*****//,1X,16H INPUT PARAMK1 1270
2METER,/,20H*****//,1X,16H INPUT PARAMK1 1280
21  30  FORMAT (2X,18H TYPE OF FUEL U=23,11,/,2X,29H NUMBER OF DELAYE MK1 1290
1UTE IN = ,12,/)                          MK1 1300
22  40  FORMAT (2X,27H NEUTRON GENERATION TIME = ,F11.5,2X,/(SEC),/) MK1 1310
23  50  FORMAT (2X,22H TYPE OF REACTIVITY = ,12,/,2X,10H RE = ,F15.3,MK1 1320
15X,/,11,/, = ,F11.5,/)                  MK1 1330
24  60  FORMAT (2X,19H TOTAL TIME USED = ,F10.3,2X,/(SEC),5X,/, TIME IMK1 1340
CINTERVAL = ,F7.3,/(SEC),/)              MK1 1350
25  70  FORMAT (2X,17H OUTPUT OPTION = ,12,/,2X,25H INITIAL RELATIVE FLUXMK1 1360
1 = ,F3.3,/)                              MK1 1370
26  80  FORMAT (5X,22H SIX GROUP LAMBDA ARE ,2X,6(F7.4,2X),/) MK1 1380
27  90  FORMAT (5X,20H SIX GROUP BETA ARE ,2X,6(F7.5,2X),/) MK1 1390
28  100 FORMAT (17H//,39H*****END OF INPUT DATA *****//) MK1 1400
29  110 FORMAT (12)                        MK1 1410
30  120 FORMAT (12)                        MK1 1420
31  130 FORMAT (110,2(F10.5))              MK1 1430
32  140 FORMAT (F10.5)                     MK1 1440
33  150 FORMAT (2(F10.5))                   MK1 1450
34  160 FORMAT (12)                         MK1 1460
35  170 FORMAT (F10.5)                     MK1 1470
C      -----COMPUTE # OF ITERATION AND TYPE OF FUEL USED MK1 1480
C      MK1 1490
C      MK1 1500
36      NI=1                                MK1 1510
37      JB=0.0                              MK1 1520
38      TH=0.0                              MK1 1530
39      H1=H                                MK1 1540
40      LL=(TIME+H1/2.)/H1                  MK1 1550
41      IF (IFUEL=5) 190,140,190            MK1 1560
42  180  CALL U235 (8,X)                     MK1 1570
43      GO TO 230                             MK1 1580
44  190  CALL PU239 (8,X)                     MK1 1590
45  200  WRITE (6,80) (X(I),I=1,6)           MK1 1600
46      WRITE (6,90) (U(I),I=1,6)           MK1 1610
C      MK1 1620
C      -----COMPUTE THE AVERAGE BETA AND LAMDA FOR THE PARTICULAR MK1 1630
C      FUEL . MK1 1640
C      MK1 1650
47      CALL AVEIX (8,X,6,NKP)               MK1 1660
48      WRITE (6,210)                        MK1 1670
49  210  FORMAT (7,5X,27H THE AVERAGE BETA AND LAMDA,/,5X,29H-----MK1 1680
1-----//)                                MK1 1690
50      DO 220 L=1,NKP                       MK1 1700
51  220  WRITE (6,230) L,B(L),L,X(L)          MK1 1710
52  230  FORMAT (6X,2H(,12,3H) = ,F11.5,5X,7HLAMBDA(,12,4H) = ,F11.5,/) MK1 1720
53      DO 460 K=1,LL                         MK1 1730
54      GO TO (240,250,260), NKL             MK1 1740
C      MK1 1750
C      -----SOL EQUATION 1.5.1 MK1 1760
C      MK1 1770
55  240  E=KJ                                MK1 1780
56      GO TO 270                             MK1 1790
C      MK1 1800

```

```

C -----SEE EQUATION 1.5.2 MK1 1810
C MK1 1820
57 250 R=RQ2(RQ,TH,B1) MK1 1830
58 GO TO 273 MK1 1840
C MK1 1850
C -----SEE EQUATION 1.5.3 MK1 1860
C MK1 1870
59 260 R=RQ3(RQ,TH,B1) MK1 1880
C MK1 1890
C -----COMPUTE THE COEFFICIENT OF INHOUR FORMULA, SEE EQUATION 1.5.4 MK1 1900
C ON 1.5.13 WHERE NN=1 MK1 1910
C MK1 1920
60 270 CALL ANHOUR (XL,X,B,R,A,NN) MK1 1930
61 MM=NN+2 MK1 1940
62 IF (K.EQ.1) WRITE (6,280) MK1 1950
63 280 FORMAT (/,4X,32H (COEFFICIENTS OF INHOUR FORMULA /,4X,35H----- MK1 1960
I-----/,/) MK1 1970
64 DO 300 I=1,MM MK1 1980
65 MW=MM-I+1 MK1 1990
66 AW(MW)=A(I) MK1 2000
67 IF (K.GT.1) GO TO 300 MK1 2010
68 WRITE (6,290) I,AW(MW) MK1 2020
69 290 FORMAT (/,5X,24H(,12,5H) = ,F11.4) MK1 2030
70 300 CONTINUE MK1 2040
71 M=NN+1 MK1 2050
C MK1 2060
C -----COMPUTE THE ROOT OF THE POLYNOMIAL WITH MM COEFFICIENTS MK1 2070
C AND DETERMINE THE LARGEST EIGEN VALUE WO IN STEP 3 MK1 2080
C MK1 2090
72 CALL POLRT (A,COF,M,W,ROOT1,IER,MM) MK1 2100
73 WO=W(1) MK1 2110
74 DO 310 J=1,M MK1 2120
75 VO=AMAX1(W(J),WO) MK1 2130
76 310 WO=VO MK1 2140
77 IF (K.NE.1) GO TO 370 MK1 2150
78 WRITE (6,350) WO MK1 2160
79 350 FORMAT (//,6X,' THE LARGEST EIGEN-VALUE = ',E11.4,/) MK1 2170
C MK1 2180
C -----FORM THE INITIAL VECTOR COLUMN PHI. MK1 2190
C INITIALIZE NEUTRON DENSITY, I=0 AND PRECURSOR DENSITY. MK1 2200
C MK1 2210
80 XNCO(1)=XN MK1 2220
81 DO 360 I=2,M MK1 2230
82 J1=I-1 MK1 2240
83 BB=BB+B(J1) MK1 2250
84 C(J1)=B(J1)/(X(J1)-X1) MK1 2260
85 360 XNCO(I)=C(J1)*XN MK1 2270
86 370 IF (K.GT.1) GG(K)=GNXCO(1) MK1 2280
87 GG(K)=XNCO(1) MK1 2290
88 KR=R*BM MK1 2300
89 IF (K-1) 380,400,380 MK1 2310
90 380 H=H1 MK1 2320
91 DO 390 J1=1,M MK1 2330
92 390 XNCO(J1)=GNXCO(J1) MK1 2340
93 GO TO 41) MK1 2350
94 400 H=0.0 MK1 2360
C MK1 2370
C -----CONSTRUCT THE G MATRIX CORRESPONDING TO EQUATION MK1 2380
C 1.4.15 AND STEP 4, AND 5 PAGE 26 AND MULTIPLY THE G-MATRIX BY INITIAL VECTOR. MK1 2390
C MK1 2400

```

MK1 2780
 MK1 2790
 MK1 2800
 MK1 2810
 MK1 2820
 MK1 2830
 MK1 2840
 MK1 2850
 MK1 2860
 MK1 2870
 MK1 2880
 MK1 2890
 MK1 2900
 MK1 2910
 MK1 2920
 MK1 2930
 MK1 2940
 MK1 2950
 MK1 2960
 MK1 2970
 MK1 2980

```

140      DO 80 J1=1,NGROUP
141      YB(J1)=ABETA(B,N3)
142      YX(J1)=ALAMDA(B,X,N3)
143      IF (J1.EQ.NGROUP) GO TO 150
144      DO 70 J2=1,N3
145      J3=J3+1
146      B(J2)=BX(J3)
147      X(J2)=XX(J3)
148      80 CONTINUE
149      GO TO 150
150      90 N4=NGROUP-2
151      J3=2
152      DO 110 J1=1,N4
153      YB(J1)=ABETA(B,N4)
154      YX(J1)=ALAMDA(B,X,N4)
155      DO 100 J2=1,N4
156      J3=J3+1
157      B(J2)=BX(J3)
158      100 X(J2)=XX(J3)
159      110 CONTINUE
160      DO 120 N3=3,NGROUP
161      N1=N3+2
162      YB(N3)=BX(N1)
163      120 YX(N3)=XX(N1)
164      GO TO 150
165      130 N5=NGROUP-1
166      YB(1)=ABETA(B,2)
167      YX(1)=ALAMDA(B,X,2)
168      DO 140 J1=2,NGROUP
169      J2=J1+1
170      YB(J1)=BX(J2)
171      140 YX(J1)=XX(J2)
172      GO TO 150
173      150 DO 160 J2=1,NGROUP
174      B(J2)=YB(J2)
175      160 X(J2)=YX(J2)
176      RETURN
177      END

```

MK1 2990
 MK1 3000
 MK1 3010
 MK1 3020
 MK1 3030
 MK1 3040
 MK1 3050
 MK1 3060
 MK1 3070
 MK1 3080
 MK1 3090
 MK1 3100
 MK1 3110
 MK1 3120
 MK1 3130
 MK1 3140
 MK1 3150
 MK1 3160
 MK1 3170
 MK1 3180
 MK1 3190
 MK1 3200
 MK1 3210
 MK1 3220
 MK1 3230
 MK1 3240
 MK1 3250
 MK1 3260
 MK1 3270
 MK1 3280
 MK1 3290
 MK1 3300
 MK1 3310
 MK1 3320
 MK1 3330
 MK1 3340
 MK1 3350
 MK1 3360
 MK1 3370
 MK1 3380
 MK1 3390
 MK1 3400
 MK1 3410
 MK1 3420

```

178      FUNCTION ALAMDA (P,X,N)
179      DIMENSION B(N), X(N)
180      B1=0.0
181      BL=0.0
182      DO 10 I=1,N
183      B1=B1+B(I)
184      BL=BL+B(I)/X(I)
185      ALAMDA=B1/BL
186      RETURN
187      END

```

MK1 3430
 MK1 3440
 MK1 3450
 MK1 3460
 MK1 3470
 MK1 3480
 MK1 3490
 MK1 3500
 MK1 3510
 MK1 3520
 MK1 3530^o
 MK1 3540
 MK1 3550
 MK1 3560

-----FUNCTION ABETA IS AVERAGING BETA .
SEE EQUATION 1.2.3

MK1 3570
MK1 3580

188 FUNCTION ABETA (S,I)
189 DIMENSION B(I)
190 ABETA=0.0
191 DO 10 I=1,N
192 10 ABETA=ABETA+B(I)
193 RETURN
194 END

MK1 3590
MK1 3600
MK1 3610
MK1 3620
MK1 3630
MK1 3640
MK1 3650
MK1 3660
MK1 3670
MK1 3680
MK1 3690
MK1 3700
MK1 3710

-----SUBROUTINE U-235 IS PROVIDING DATA FOR BETA AND LAMB
FOR SIX GROUP DELAYED NEUTRON FROM THERMAL FISSION.

195 SUBROUTINE U235 (R,X)
196 DIMENSION B(6), X(6)
197 B(1)=.00021
198 B(2)=.00141
199 B(3)=.00127
200 B(4)=.00255
201 B(5)=.00074
202 B(6)=.00027
203 X(1)=.0124
204 X(2)=.0305
205 X(3)=.111
206 X(4)=.301
207 X(5)=1.14
208 X(6)=3.01
209 RETURN
210 END

MK1 3720
MK1 3730
MK1 3740
MK1 3750
MK1 3760
MK1 3770
MK1 3780
MK1 3790
MK1 3800
MK1 3810
MK1 3820
MK1 3830
MK1 3840
MK1 3850
MK1 3860
MK1 3870
MK1 3880
MK1 3890
MK1 3900

-----SUBROUTINE PU-239 IS PROVIDING DATA FOR BETA AND LAMB
DA SIX GROUP DELAYED NEUTRON FROM THERMAL FISSION.

211 SUBROUTINE PU239 (R,X)
212 DIMENSION B(6), X(6)
213 B(1)=.00009450
214 B(2)=.00080460
215 B(3)=.00056970
216 B(4)=.00088020
217 B(5)=.00023220
218 B(6)=.00011880
219 X(1)=.0128
220 X(2)=.0301
221 X(3)=.124
222 X(4)=.325
223 X(5)=1.12
224 X(6)=2.69
225 RETURN
226 END

MK1 3940
MK1 3950
MK1 3960
MK1 3970
MK1 3980
MK1 3990
MK1 4000
MK1 4010
MK1 4020
MK1 4030
MK1 4040
MK1 4050
MK1 4060
MK1 4070
MK1 4080
MK1 4090
MK1 4100

C
C
C
C
C

-----THE MAIN PROGRAM IS TRYING TO CALCULATE THE COEFFICIENTS OF INHOUR FORMULA FOR ANY GROUP BETWEEN 1 TO 6 .

MK1 4110
MK1 4120
MK1 4130
MK1 4140
MK1 4150

```

227 SUBROUTINE ANHOUR (XL,X,B,R,A,NN)
228 DIMENSION X(NN), P(NN), A(8)
229 BB=0.0
230 DO 10 I=1,NN
231 10 BB=BB+B(I)
232 RO=R*BB
233 JJ=NN+2
234 DO 20 J=1,JJ
235 20 A(J)=0.0
C SEE 1.13.13 FOR NN=1
236 GO TO (30,40,50,60,70,80), NN
237 30 CALL G1 (XL,X,B,RO,A)
238 GO TO 90
239 40 CALL G2 (XL,X,B,RO,A)
240 GO TO 90
241 50 CALL G3 (XL,X,B,RO,A)
242 GO TO 90
243 60 CALL G4 (XL,X,B,RO,A)
244 GO TO 90
245 70 CALL G5 (XL,X,B,RO,A)
246 GO TO 90
247 80 CALL G6 (XL,X,B,RO,A)
248 GO TO 90
249 90 RETURN
250 END

```

MK1 4160
MK1 4170
MK1 4180
MK1 4190
MK1 4200
MK1 4210
MK1 4220
MK1 4230
MK1 4240
MK1 4250
MK1 4260
MK1 4270
MK1 4280
MK1 4290
MK1 4300
MK1 4310
MK1 4320
MK1 4330
MK1 4340
MK1 4350
MK1 4360
MK1 4370
MK1 4380
MK1 4390
MK1 4400
MK1 4410
MK1 4420

C
C
C
C
C
C
C
C
C
C
C
C

-----SUBROUTINE G1,G2,G3,G4,G5,G6, IS DIRECTLY CALCULATING COEFFICIENTS OF INHOUR FORMULA IN POLYNOMIAL FORM IN THE ORDER OF THE SMALLEST DEGREE TO THE LARGEST. THE FORM OUTPUT IS $A1*W**N + A2*W**(N-1) + A3*W**(N-2) + \dots + AN$ SEE EQUATION 1.3.15 FOR THE GENERAL INHOUR FORMULA

MK1 4430
MK1 4440
MK1 4450
MK1 4460
MK1 4470
MK1 4480
MK1 4490
MK1 4500
MK1 4510
MK1 4520
MK1 4530

SUPPORTING ROUTINE A62,A52,A42,A32,A22 AND XX,XY,XZ,XU,XV,AA

```

251 SUBROUTINE G1 (XL,X,B,RO,W)
252 DIMENSION X(1), B(1), W(3)
253 W(3)=XL
254 W(2)=B(1)+XL*X(1)-RO
255 W(1)=-RO*X(1)
256 RETURN
257 END

```

MK1 4540
MK1 4550
MK1 4560
MK1 4570
MK1 4580
MK1 4590
MK1 4600

```

258 SUBROUTINE G2 (XL,X,B,RO,W)
259 DIMENSION X(2), B(2), W(4)
260 W(4)=XL
261 W(3)=XX(B,2)+XL*XX(X,2)-RO

```

MK1 4610
MK1 4620
MK1 4630
MK1 4640

262	W(2)=XL*AA(X,2)+A62(B,X,2)-RO*XX(X,2)	MK1 4650
263	W(1)=-RO*AA(X,2)	MK1 4660
264	RETURN	MK1 4670
265	END	MK1 4680

266	SUBROUTINE G3 (XL,X,B,RC,W)	MK1 4690
267	DIMENSION X(3), B(3), W(5)	MK1 4700
268	W(5)=XL	MK1 4710
269	W(4)=XX(B,3)+XL*XX(X,3)-RO	MK1 4720
270	W(3)=XL*XY(X,3)+A62(B,X,3)-RO*XX(X,3)	MK1 4730
271	W(2)=XL*AA(X,3)+A52(B,X,3)-RO*XY(X,3)	MK1 4740
272	W(1)=-RO*AA(X,3)	MK1 4750
273	RETURN	MK1 4760
274	END	MK1 4770

275	* SUBROUTINE G4 (XL,X,B,PC,W)	MK1 4780
276	DIMENSION X(4), B(4), W(6)	MK1 4790
277	W(6)=XL	MK1 4800
278	W(5)=XX(B,4)+XL*XX(X,4)-RO	MK1 4810
279	W(4)=XL*XY(X,4)+A62(B,X,4)-RO*XX(X,4)	MK1 4820
280	W(3)=XL*XZ(X,4)+A52(B,X,4)-RO*XY(X,4)	MK1 4830
281	W(2)=XL*AA(X,4)+A42(B,X,4)-RO*XZ(X,4)	MK1 4840
282	W(1)=-RO*AA(X,4)	MK1 4850
283	RETURN	MK1 4860
284	END	MK1 4870

285	SUBROUTINE G5 (XL,X,B,PC,W)	MK1 4880
286	DIMENSION X(5), B(5), W(7)	MK1 4890
287	W(7)=XL	MK1 4900
288	W(6)=XX(B,5)+XL*XX(X,5)-RO	MK1 4910
289	W(5)=XL*XY(X,5)+A62(B,X,5)-RO*XX(X,5)	MK1 4920
290	W(4)=XL*XZ(X,5)+A52(B,X,5)-RO*XY(X,5)	MK1 4930
291	W(3)=XL*XU(X,5)+A42(B,X,5)-RO*XZ(X,5)	MK1 4940
292	W(2)=XL*AA(X,5)+A32(B,X,5)-RO*XU(X,5)	MK1 4950
293	W(1)=-RO*AA(X,5)	MK1 4960
294	RETURN	MK1 4970
295	END	MK1 4980

296	SUBROUTINE G6 (XL,X,B,RC,W)	MK1 4990
297	DIMENSION X(6), B(6), W(8)	MK1 5000
298	W(8)=XL	MK1 5010
299	W(7)=XX(B,6)+XL*XX(X,6)-RO	MK1 5020
300	W(6)=XL*XY(X,6)+A62(B,X,6)-RO*XX(X,6)	MK1 5030
301	W(5)=XL*XZ(X,6)+A52(B,X,6)-RO*XY(X,6)	MK1 5040
302	W(4)=XL*XU(X,6)+A42(B,X,6)-RO*XZ(X,6)	MK1 5050
303	W(3)=XL*XV(X,6)+A32(B,X,6)-RO*XU(X,6)	MK1 5060
304	W(2)=XL*AA(X,6)+A22(B,X,6)-RO*XV(X,6)	MK1 5070
305	W(1)=-RO*AA(X,6)	MK1 5080
306	RETURN	MK1 5090
307	END	MK1 5100

C
C
C
C
C
C

-----FUNCTION A62 FUNCTION AS :

A62 = P1*(X2+X3+X4+....) + B2*(B1+B3+B4+....) +
B3*(B1+B2+B3+B4+....) + IN THAT COMBINATIONS.

MK1 5130
MK1 5140
MK1 5150
MK1 5160

G			MK1 5170
308		FUNCTION A62 (B,X,M)	MK1 5180
309		DIMENSION X(M), AX1(6), B(M), W(6)	MK1 5190
310		A62=0.0	MK1 5200
311		DO 20 I=1,M	MK1 5210
312		DO 10 J=1,M	MK1 5220
313	10	AX1(J)=X(J)	MK1 5230
314		AX1(I)=0.0	MK1 5240
315		W(I)=B(I)*XX(AX1,M)	MK1 5250
316	20	A62=A62+W(I)	MK1 5260
317		RETURN	MK1 5270
318		END	MK1 5280
	C		MK1 5290
	C		MK1 5300
	C	-----A52 IS ALMOST THE SAME THING AS A62 IN IT'S FUNCTION	MK1 5310
	C	EXCEPT THAT X1 = Y1*Y2 WHICH THE POSTSCRIPT NEVER BE	MK1 5320
	C	EQUAL AND THE FIRST CHARACTER POSTSCRIPT IS ALWAYS	MK1 5330
	C	SMALLER THAN THE NEXT	MK1 5340
	C		MK1 5350
	C		MK1 5360

319		FUNCTION A52 (B,X,M)	MK1 5370
320		DIMENSION X(M), AX2(6), B(M), W(6)	MK1 5380
321		A52=0.0	MK1 5390
322		DO 20 I=1,M	MK1 5400
323		DO 10 J=1,M	MK1 5410
324	10	AX2(J)=X(J)	MK1 5420
325		AX2(I)=0.0	MK1 5430
326		W(I)=B(I)*XY(AX2,M)	MK1 5440
327	20	A52=A52+W(I)	MK1 5450
328		RETURN	MK1 5460
329		END	MK1 5470
	C		MK1 5480
	C		MK1 5490
	C	-----OIO AS A62 EXCEPT ITS ADDITION OF THREE CHARACTER	MK1 5500
	C		MK1 5510
	C		MK1 5520

330		FUNCTION A42 (B,X,M)	MK1 5530
331		DIMENSION X(M), AX3(6), B(M), W(6)	MK1 5540
332		A42=0.0	MK1 5550
333		DO 20 I=1,M	MK1 5560
334		DO 10 J=1,M	MK1 5570
335	10	AX3(J)=X(J)	MK1 5580
336		AX3(I)=0.0	MK1 5590
337		W(I)=B(I)*XZ(AX3,M)	MK1 5600
338	20	A42=A42+W(I)	MK1 5610
339		RETURN	MK1 5620
340		END	MK1 5630
	C		MK1 5640
	C		MK1 5650
	C	-----OEO AS A62 EXCEPT CHARACTER ADDITION IS FOUR	MK1 5660
	C		MK1 5670
	C		MK1 5680

```

341 FUNCTION A32 (B,X,M)
342 DIMENSION X(M), AX4(6), U(M), W(6)
343 A32=0.0
344 DO 20 I=1,M
345 DO 10 J=1,M
346 10 AX4(J)=X(J)
347 AX4(I)=0.0
348 W(I)=B(I)*XU(AX4,M)
349 20 A32=A32+W(I)
350 RETURN
351 END

```

C
C
C
C
C

-----USED AS A62 EXCEPT ADDITION CHARACTER IS FIVE .

MK1 5690
MK1 5700
MK1 5710
MK1 5720
MK1 5730
MK1 5740
MK1 5750
MK1 5760
MK1 5770
MK1 5780
MK1 5790
MK1 5800
MK1 5810
MK1 5820
MK1 5830
MK1 5840

```

352 FUNCTION A22 (B,X,M)
353 DIMENSION X(M), AX5(6), U(M), W(6)
354 A22=0.0
355 DO 20 I=1,M
356 DO 10 J=1,M
357 10 AX5(J)=X(J)
358 AX5(I)=0.0
359 W(I)=B(I)*XV(AX5,M)
360 20 A22=A22+W(I)
361 RETURN
362 END

```

C
C
C
C
C
C

-----XX IS COMPUTING THE TOTAL OF X
XX = (X1+X2+X3+X4+.....+XN)

MK1 5850
MK1 5860
MK1 5870
MK1 5880
MK1 5890
MK1 5900
MK1 5910
MK1 5920
MK1 5930
MK1 5940
MK1 5950
MK1 5960
MK1 5970
MK1 5980
MK1 5990
MK1 6000
MK1 6010

```

363 FUNCTION XX (X,M)
364 DIMENSION X(M)
365 XX=0.0
366 DO 10 I=1,M
367 10 XX=XX+X(I)
368 RETURN
369 END

```

C
C
C
C
C
C

-----XY IS COMPUTING THE ADDITION OF TWO PRODUCTS .
XY = (X1*X2+X1*X3+X1*X4+.....+X2*X3+X2*X4+.....)
THE FIRST POSTSCRIPT ALWAYS SMALLER THAN THE NEXT

MK1 6020
MK1 6030
MK1 6040
MK1 6050
MK1 6060
MK1 6070
MK1 6080
MK1 6090
MK1 6100
MK1 6110
MK1 6120
MK1 6130
MK1 6140
MK1 6150

```

370 FUNCTION XY (X,M)
371 DIMENSION X(M)
372 XY=0.0
373 DO 20 I=1,M
374 DO 20 J=1,M
375 A=X(I)
376 B=X(J)

```

MK1 6160
MK1 6170
MK1 6180
MK1 6190
MK1 6200
MK1 6210
MK1 6220

```

377 IF (I-J) 10,20,20
378 10 Y=A*B
379 XY=XY+Y
380 20 CONTINUE
381 RETURN
382 END

```

C
C
C
C
C

-----DEF AS XY EXCEPT IT IS A PRODUCT OF THREE

MK1 6230
MK1 6240
MK1 6250
MK1 6260
MK1 6270
MK1 6280
MK1 6290
MK1 6300
MK1 6310
MK1 6320
MK1 6330

```

383 FUNCTION XZ (X,M)
384 DIMENSION X(M)
385 XZ=0.0
386 DO 30 I=1,M
387 DO 30 J=1,M
388 DO 30 K=1,M
389 A=X(I)
390 B=X(J)
391 C=X(K)
392 IF (I-J) 10,30,30
393 10 IF (J-K) 20,30,30
394 20 Y=A*B*C
395 XZ=XZ+Y
396 30 CONTINUE
397 RETURN
398 END

```

C
C
C
C
C

-----DEF AS XY EXCEPT FOR PRODUCT OF FOUR

MK1 6340
MK1 6350
MK1 6360
MK1 6370
MK1 6380
MK1 6390
MK1 6400
MK1 6410
MK1 6420
MK1 6430
MK1 6440
MK1 6450
MK1 6460
MK1 6470
MK1 6480
MK1 6490
MK1 6500
MK1 6510
MK1 6520
MK1 6530
MK1 6540

```

399 FUNCTION XU (X,M)
400 DIMENSION X(M)
401 XU=0.0
402 DO 40 I=1,M
403 DO 40 J=1,M
404 DO 40 K=1,M
405 DO 40 L=1,M
406 A=X(I)
407 B=X(J)
408 C=X(K)
409 D=X(L)
410 IF (I-J) 10,40,40
411 10 IF (J-K) 20,40,40
412 20 IF (K-L) 30,40,40
413 30 Y=A*B*C*D
414 XU=XU+Y
415 40 CONTINUE
416 RETURN
417 END

```

C
C
C
C
C

-----DEF EXCEPT ITS AN ADDITION OF MULTIPLICATION OF FIVE

MK1 6550
MK1 6560
MK1 6570
MK1 6580
MK1 6590
MK1 6600
MK1 6610
MK1 6620
MK1 6630
MK1 6640
MK1 6650
MK1 6660
MK1 6670
MK1 6680
MK1 6690
MK1 6700
MK1 6710
MK1 6720
MK1 6730
MK1 6740
MK1 6750
MK1 6760
MK1 6770
MK1 6780

```

418 FUNCTION XV (X,M)
419 DIMENSION X(M)
420 XV=0.0
421 DO 50 I=1,M
422 DO 50 J=1,M
423 DO 50 K=1,M
424 DO 50 L=1,M
425 DO 50 N=1,M
426 A=X(I)
427 B=X(J)
428 C=X(K)
429 D=X(L)
430 E=X(N)
431 IF (I-J) 10,50,50
432 10 IF (J-K) 20,50,50
433 20 IF (K-L) 30,50,50
434 30 IF (L-N) 40,50,50
435 40 Y=A*B*C*D*E
436 XV=XV+Y
437 50 CONTINUE
438 WRITE (6,60) XV
439 60 FORMAT (2X,7H' XV = ,F19.10)
440 RETURN
441 END

```

C
C
C
C
C
C
C

-----AA IS COMPUTING MULTIPLICATION OF N CHARACTER .
AA = X1*X2*X3*X4*.....*XN

MK1 6790
MK1 6800
MK1 6810
MK1 6820
MK1 6830
MK1 6840
MK1 6850
MK1 6860
MK1 6870
MK1 6880
MK1 6890
MK1 6900
MK1 6910
MK1 6920
MK1 6930
MK1 6940
MK1 6950
MK1 6960
MK1 6970
MK1 6980
MK1 6990
MK1 7000
MK1 7010
MK1 7020
MK1 7030
MK1 7040
MK1 7050
MK1 7060
MK1 7070
MK1 7080

```

442 FUNCTION AA (X,M)
443 DIMENSION X(M)
444 AA=X(1)
445 DO 10 I=2,M
446 10 AA=AA*X(I)
447 RETURN
448 END

```

C
C
C
C
C
C
C
C

-----FORM N X N MATRIX BY HANSEN METHOD , G - MATRIX .
THE FIRST COLUMN , THE FIRST ROW AND THE DIAGONAL = 0,
THE REST OF THEM = 0 .
SEE EQUATION 1.4.15

SUPPORTING ROUTINE NONE

MK1 7090
MK1 7100
MK1 7110
MK1 7120
MK1 7130
MK1 7140
MK1 7150
MK1 7160
MK1 7170
MK1 7180
MK1 7190
MK1 7200
MK1 7210
MK1 7220
MK1 7230
MK1 7240
MK1 7250

```

449 SUBROUTINE GMTEX (C,B,X,XL,WU,N,I,G)
450 DIMENSION E(6), X(6), G(N,N)
451 GD(AX,H)=EXP(-AX*H)
452 GH(RBL,WJ,H,AX)=((EXP(WC*H)-EXP(PBL*H))/(WU-RBL))*AX
453 GV(XL,WU,H,AX,BX)=((EXP(WC*H)-GD(AX,H))*BX)/((WU+AX)*XL)
454 RB(R,BX)=(R-BX)/XL
455 BB=0.0
456 NN=N-1
457 DO 10 I=1,NN

```

MK1 7260
MK1 7270
MK1 7280
MK1 7290
MK1 7300
MK1 7310
MK1 7320
MK1 7330
MK1 7340

```

458 10 BB=BB+B(I)
459 DO 150 J=1,N
460 DO 140 JJ=1,N
461 JI=JJ
462 RBL=RB(R,BB)
463 IF (J-JJ) 20,60,20
464 20 IF (J-1) 30,50,30
465 30 JI=JJ+1
466 IF (JJ-N) 50,40,50
467 40 IF (J-N) 120,80,120
468 50 AX=X(JI-1)
469 GO TO 90
470 60 IF (J-1) 80,70,80
471 70 G(I,1)=EXP(RB(P,BB)*H)
472 GO TO 140
473 80 AX=X(J-1)
474 G(J,JJ)=GD(AX,H)
475 GO TO 140
476 90 IF (J-1) 110,100,110
477 100 G(I,JJ)=GH(RBL,WD,H,AX)
478 GO TO 140
479 110 IF (JJ-1) 120,130,120
480 120 G(J,JJ)=0.0
481 GO TO 140
482 130 AX=X(J-1)
483 BX=B(J-1)
484 G(J,1)=GV(XL,WD,H,AX,BX)
485 140 CONTINUE
486 I=J
C WRITE(6,144)(G(I,1),I=1,N)
C 144 FORMAT('***',2(E10.3,2X), '***')
487 150 CONTINUE
488 RETURN
489 END

```

C
C
C
C
C
C
C

-----MULTIPLY THE G - MATRIX WITH THE INITIAL VECTOR COL
TO GET THE NEXT ONE , ITERATION.
SEE EQUATION 1.4.3.

```

490 SUBROUTINE GXN (G,AN,N,GX)
491 DIMENSION G(N,N), XN(N), GX(N)
492 DO 20 I=1,N
493 GG=0.0
494 DO 10 J=1,N
495 10 GG=GG+G(I,J)*XN(J)
496 20 GX(I)=GG
497 RETURN
498 END

```

C
C
C
C
C
C
C

SUBROUTINE POLPT

PURPOSE

COMPUTES THE REAL AND COMPLEX ROOTS OF A REAL POLYNOMIAL

MK1 7350
MK1 7360
MK1 7370
MK1 7380
MK1 7390
MK1 7400
MK1 7410
MK1 7420
MK1 7430
MK1 7440
MK1 7450
MK1 7460
MK1 7470
MK1 7480
MK1 7490
MK1 7500
MK1 7510
MK1 7520
MK1 7530
MK1 7540
MK1 7550
MK1 7560
MK1 7570
MK1 7580
MK1 7590
MK1 7600
MK1 7610
MK1 7620
MK1 7630
MK1 7640
MK1 7650
MK1 7660
MK1 7670
MK1 7680
MK1 7690
MK1 7700
MK1 7710
MK1 7720
MK1 7730
MK1 7740
MK1 7750

MK1 7760
MK1 7770
MK1 7780
MK1 7790
MK1 7800
MK1 7810
MK1 7820
MK1 7830
MK1 7840
MK1 7850
MK1 7860
MK1 7870
MK1 7880
MK1 7890
MK1 7900
MK1 7910
MK1 7920

USAGE
CALL POLRT(XCF,CDF,M,ROOTR,ROOTI,IFR,M1)

DESCRIPTION OF PARAMETERS

XCF - VECTOR OF M+1 COEFFICIENTS OF THE POLYNOMIAL
ORDERED FROM SMALLEST TO LARGEST POWER
CDF - WORKING VECTOR OF LENGTH M+1
M - ORDER OF POLYNOMIAL
ROOTR - RESULTANT VECTOR OF LENGTH M CONTAINING REAL ROOTS
OF THE POLYNOMIAL
ROOTI - RESULTANT VECTOR OF LENGTH M CONTAINING THE
CORRESPONDING IMAGINARY ROOTS OF THE POLYNOMIAL
IFR - ERROR CODE WHERE
IFR=0 NO ERROR
IFR=1 M LESS THAN ONE
IFR=2 M GREATER THAN 36
IFR=3 UNABLE TO DETERMINE ROOT WITH 500 ITERATIONS
ON 5 STARTING VALUES
IFR=4 HIGH ORDER COEFFICIENT IS ZERO
M1 - NUMBER OF COEFFICIENT, M+1

REMARKS

LIMITED TO 36TH ORDER POLYNOMIAL OR LESS.
FLATTING POINT OVERFLOW MAY OCCUR FOR HIGH ORDER
POLYNOMIALS BUT WILL NOT AFFECT THE ACCURACY OF THE
RESULTS.

SUBROUTINES AND FUNCTION SUBPROGRAMS REQUIRED
NONE

METHOD

NEWTON-RAPHSON ITERATIVE TECHNIQUE. THE FINAL ITERATIONS
ON EACH ROOT ARE PERFORMED USING THE ORIGINAL POLYNOMIAL
RATHER THAN THE REDUCED POLYNOMIAL TO AVOID ACCUMULATED
ERRORS IN THE REDUCED POLYNOMIAL.

MK1 7930
MK1 7940
MK1 7950
MK1 7960
MK1 7970
MK1 7980
MK1 7990
MK1 8000
MK1 8010
MK1 8020
MK1 8030
MK1 8040
MK1 8050
MK1 8060
MK1 8070
MK1 8080
MK1 8090
MK1 8100
MK1 8110
MK1 8120
MK1 8130
MK1 8140
MK1 8150
MK1 8160
MK1 8170
MK1 8180
MK1 8190
MK1 8200
MK1 8210
MK1 8220
MK1 8230
MK1 8240
MK1 8250
MK1 8260
MK1 8270
MK1 8280
MK1 8290

499
500

SUBROUTINE POLRT (XCF,CDF,M,ROOTR,ROOTI,IFR,M1)
DIMENSION XCF(M1), CDF(M1), ROOTR(M), ROOTI(M)
DOUBLE PRECISION XU,YC,X,Y,XPR,YPP,UX,UY,V,YT,XT,U,XT2,YT2,SUMSQ,
1 DX,DY,TEMP,ALPHA,CAHS

IF A DOUBLE PRECISION VERSION OF THIS ROUTINE IS DESIRED, THE
C IN COLUMN 1 SHOULD BE REMOVED FROM THE DOUBLE PRECISION
STATEMENT WHICH FOLLOWS.

DOUBLE PRECISION XCF,CDF,ROOTR,ROOTI

THE C MUST ALSO BE REMOVED FROM DOUBLE PRECISION STATEMENTS
APPEARING IN OTHER ROUTINES USED IN CONJUNCTION WITH THIS
ROUTINE.

THE DOUBLE PRECISION VERSION MAY BE MODIFIED BY CHANGING THE
CONSTANT IN STATEMENT 78 TO 1.0D-12 AND IN STATEMENT 122 TO
1.0D-10. THIS WILL PROVIDE HIGHER PRECISION RESULTS AT THE
COST OF EXECUTION TIME

MK1 8300
MK1 8310
MK1 8320
MK1 8330
MK1 8340
MK1 8350
MK1 8360
MK1 8370
MK1 8380
MK1 8390
MK1 8400
MK1 8410
MK1 8420
MK1 8430
MK1 8440
MK1 8450
MK1 8460
MK1 8470
MK1 8480
MK1 8490
MK1 8500

501 IFIT=0
502 N=M
503 IER=0
504 IF (XCUF(N+1)) 10,40,10
505 IF (N) 20,20,60

10

C

C

C

C

506

20

507

30

C

C

C

508

40

509

C

C

C

C

510

50

511

512

513

514

515

516

517

518

519

60

C

C

C

C

520

90

521

C

C

C

C

522

523

100

C

C

C

524

525

C

C

C

C

526

527

528

529

530

531

532

110

C

C

C

C

C

533

534

535

536

120

130

IFIT=0
N=M
IER=0
IF (XCUF(N+1)) 10,40,10
IF (N) 20,20,60

SET ERROR CODE TO 1

IER=1
RETURN

SET ERROR CODE TO 4

IER=4
GO TO 30

SET ERROR CODE TO 2

IER=2
GO TO 30
IF (N-36) 70,70,50
NX=N
NXX=N+1
N2=1
KJ1=N+1
DO 40, L=1, KJ1
MT=KJ1-L+1
CUF(MT)=XCUF(L)

SET INITIAL VALUES

X0=.00500101
Y0=.01000101

ZERO INITIAL VALUE COUNTER

IN=0
X=X1

INCREMENT INITIAL VALUES AND COUNTER

X0=-10.0*Y0
Y0=-10.0*X

SET X AND Y TO CURRENT VALUE

X=X1
Y=Y1
IN=IN+1
GO TO 120
IFIT=1
XPR=X
YPR=Y

EVALUATE POLYNOMIAL AND DERIVATIVES

ICT=3
UX=0.0
UY=0.0
V=0.0

MK1 8510
MK1 8520
MK1 8530
MK1 8540
MK1 8550
MK1 8560
MK1 8570
MK1 8580
MK1 8590
MK1 8600
MK1 8610
MK1 8620
MK1 8630
MK1 8640
MK1 8650
MK1 8660
MK1 8670
MK1 8680
MK1 8690
MK1 8700
MK1 8710
MK1 8720
MK1 8730
MK1 8740
MK1 8750
MK1 8760
MK1 8770
MK1 8780
MK1 8790
MK1 8800
MK1 8810
MK1 8820
MK1 8830
MK1 8840
MK1 8850
MK1 8860
MK1 8870
MK1 8880
MK1 8890
MK1 8900
MK1 8910
MK1 8920
MK1 8930
MK1 8940
MK1 8950
MK1 8960
MK1 8970
MK1 8980
MK1 8990
MK1 9000
MK1 9010
MK1 9020
MK1 9030
MK1 9040
MK1 9050
MK1 9060
MK1 9070
MK1 9080
MK1 9090
MK1 9100

537		YT=J.C	MK1 9110
538		XT=1.0	MK1 9120
539		U=COF(N+1)	MK1 9130
540		IF (U) 140,270,140	MK1 9140
541	140	DO 150 I=1,N	MK1 9150
542		L=N-I+1	MK1 9160
543		TEMP=COF(L)	MK1 9170
544		XT2=X*XT-Y*YT	MK1 9180
545		YT2=X*YT+Y*XT	MK1 9190
546		U=U+TEMP*XT2	MK1 9200
547		V=V+TEMP*YT2	MK1 9210
548		F1=!	MK1 9220
549		UX=UX+F1*XT*TEMP	MK1 9230
550		UY=UY-F1*YT*TEMP	MK1 9240
551		XT=XT2	MK1 9250
552	150	YT=YT2	MK1 9260
553		SUMSQ=UX*UX+UY*UY	MK1 9270
554		IF (SUMSQ) 160,230,160	MK1 9280
555	160	UX=(V*UY-U*UX)/SUMSQ	MK1 9290
556		X=X+UX	MK1 9300
557		UY=-(U*UY+V*UX)/SUMSQ	MK1 9310
558		Y=Y+UY	MK1 9320
559		IF (ABS(OY)+ABS(OX)-1.0E-5) 210,170,170	MK1 9330
	C		MK1 9340
	C	STEP ITERATION COUNTER C	MK1 9350
	C		MK1 9360
560	170	ICT=ICT+1	MK1 9370
561		IF (ICT-500) 130,180,180	MK1 9380
562	180	IF (IFIT) 210,190,210	MK1 9390
563	190	IF (IN-5) 100,200,200	MK1 9400
	C		MK1 9410
	C	SET ERROR CODE TO 3	MK1 9420
	C		MK1 9430
564	200	IEP=3	MK1 9440
565		GO TO 30	MK1 9450
566	210	DO 220 L=1,NXX	MK1 9460
567		MT=XJ1-L+1	MK1 9470
568		TEMP=XCOF(MT)	MK1 9480
569		XCOF(MT)=COF(L)	MK1 9490
570	220	COF(L)=TEMP	MK1 9500
571		ITEMP=N	MK1 9510
572		N=NX	MK1 9520
573		NX=ITEMP	MK1 9530
574		IF (IFIT) 250,110,250	MK1 9540
575	230	IF (IFIT) 240,100,240	MK1 9550
576	240	X=XPR	MK1 9560
577		Y=YPR	MK1 9570
578	250	IFIT=0	MK1 9580
579		IF (ABS(Y)-1.0E-4*ABS(X)) 280,260,260	MK1 9590
580	260	ALPHA=X+X	MK1 9600
581		SUMSQ=X*X+Y*Y	MK1 9610
582		N=N-2	MK1 9620
583		GO TO 290	MK1 9630
584	270	X=0.0	MK1 9640
585		NX=NX-1	MK1 9650
586		NXX=NXX-1	MK1 9660
587	280	Y=0.0	MK1 9670
588		SUMSQ=0.0	MK1 9680
589		ALPHA=X	MK1 9690
590		N=N-1	MK1 9700

```

591 290 CDF(2)=CDF(2)+ALPHA*CDF(1)
592 IF (N.EQ.0) GO TO 310
593 DO 300 L=2,N
594 CDF(L+1)=CDF(L)+ALPHA*CDF(L)-SUMSQ*CDF(L-1)
595 310 ROOT1(N2)=Y
596 ROOTR(N2)=X
597 N2=N2+1
598 IF (SUMSQ) 320,330,320
599 320 Y=-Y
600 SUMSQ=0.0
601 GO TO 310
602 330 IF (N) 30,30,9)
603 END

```

MK1 9710
 MK1 9720
 MK1 9730
 MK1 9740
 MK1 9750
 MK1 9760
 MK1 9770
 MK1 9780
 MK1 9790
 MK1 9800
 MK1 9810
 MK1 9820
 MK1 9830
 MK1 9840
 MK1 9850
 MK1 9860
 MK1 9870
 MK1 9880
 MK1 9890
 MK1 9900
 MK1 9910
 MK1 9920
 MK1 9930
 MK1 9940

C
 C
 C
 C
 C
 C
 C
 C
 C
 C

-----SUBROUTINE TPLDT IS SINGLE PRECISION PLOTTING 5 DIFFERENT
 VARIABLES, ASSUMING CONSTANT TIME INCREMENT BETWEEN TWO
 EVENTS.

SUPPORTING ROUTINE NONE

```

604 SUBROUTINE TPLDT (M1,M2,M8,M9,M0,JX)
605 IMPLICIT REAL*4(A-H,Y-Z)
606 DIMENSION M8(JX), M9(JX), M0(JX), M1(JX), M2(JX)
607 DIMENSION LINE(41), INUM(9)
608 INTEGER PL,M1,ST,BL,SL,S9,S0,S1,S2
609 READ (5,90) PL,M1,ST,BL,SL,S9,S0,S1,S2,S6
610 MXY=0.0
611 MIN0=0.0
612 MIN2=0.0
613 MIN8=0.0
614 MIN9=0.0
615 PHI0=0.0
616 PHI2=0.0
617 PHI8=0.0
618 PHI9=0.0
619 DO 10 I=1,JX
620 IF (MIN0.GT.M0(I)) MIN0=M0(I)
621 IF (MIN2.GT.M2(I)) MIN2=M2(I)
622 IF (MIN8.GT.M8(I)) MIN8=M8(I)
623 IF (MIN9.GT.M9(I)) MIN9=M9(I)
624 IF (ABS(M0(I)).GT.PHI0) PHI0=ABS(M0(I))
625 IF (ABS(M2(I)).GT.PHI2) PHI2=ABS(M2(I))
626 IF (ABS(M8(I)).GT.PHI8) PHI8=ABS(M8(I))
627 IF (ABS(M9(I)).GT.PHI9) PHI9=ABS(M9(I))
628 10 CONTINUE
629 JJ=JX
630 JJO=JJ+6+1
631 JJI=JJ+1
632 WRITE (6,90)
633 WRITE (6,100)
634 PHI0=PHI0+ABS(MIN0)
635 PHI2=PHI2+ABS(MIN2)
636 PHI8=PHI8+ABS(MIN8)
637 PHI9=PHI9+ABS(MIN9)

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MK1 9950
 MK1 9960
 MK1 9970
 MK1 9980
 MK1 9990
 MK110000
 MK110010
 MK110020
 MK110030
 MK110040
 MK110050
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 MK110210
 MK110220
 MK110230
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 MK110260
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 MK110280

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638 > DO 20 I=1,JJ
639 IF (MIN0.LT.0.0) M0(I)=M0(I)+ABS(MIN0)
640 IF (MIN2.LT.0.0) M2(I)=M2(I)+ABS(MIN2)
641 IF (MIN8.LT.0.0) M8(I)=M8(I)+ABS(MIN8)
642 IF (MIN9.LT.0.0) M9(I)=M9(I)+ABS(MIN9)
643 M0(I)=M0(I)/PHI0
644 M2(I)=M2(I)/PHI2
645 M8(I)=M8(I)/PHI8
646 20 M9(I)=M9(I)/PHI9
647 DO 30 I=1,9
648 30 INUM(I)=I
649 WRITE (6,110) (INUM(I),I=1,9)
650 DO 70 I=1,JJ1
651 IF (1.EQ.1) GO TO 50
652 MXY=M1(I-1)
653 IP8=48(I-1)*60+1.0
654 IP9=49(I-1)*60+1.0
655 IP0=M0(I-1)*60+1.0
656 IP2=M2(I-1)*60+1.0
657 DO 40 I1=1,56,5
658 LINE(I1)=8L
659 DO 40 I2=1,5
660 I3=I1+I2
661 IF (I1.EQ.IP0) LINE(I1)=S0
662 IF (I3.EQ.IP0) LINE(I3)=S0
663 IF (I1.EQ.IP2) LINE(I1)=S2
664 IF (I3.EQ.IP2) LINE(I3)=S2
665 IF (I1.EQ.IP3) LINE(I1)=S1
666 IF (I3.EQ.IP3) LINE(I3)=S1
667 IF (I1.EQ.IP9) LINE(I1)=S9
668 IF (I3.EQ.IP9) LINE(I3)=S9
669 40 CONTINUE
670 LINE(61)=PL
671 I1L=I-1
672 IF (I1.EQ.IP0) LINE(I1)=S0
673 IF (I1.EQ.IP2) LINE(I1)=S2
674 IF (I1.EQ.IP8) LINE(I1)=S1
675 IF (I1.EQ.IP9) LINE(I1)=S9
676 IF (IP0.EQ.61) LINE(61)=S0
677 IF (IP2.EQ.61) LINE(61)=S2
678 IF (IP8.EQ.61) LINE(61)=S1
679 IF (IP9.EQ.61) LINE(61)=S9
680 IF (IP9.NE.1.OR.IP8.NE.1.OR.IP0.NE.1.OR.IP2.NE.1) LINE(I1)=PL
681 WRITE(6,120)MXY,(LINE(KK),KK=1,61)
682 IF (1.EQ.JJ1) GO TO 70
683 50 CONTINUE
684 DO 60 I1=1,56,5
685 DO 60 I2=1,5
686 I3=I1+I2
687 LINE(I3)=8L
688 60 CONTINUE
689 70 CONTINUE
690 WRITE(6,130) (INUM(I),I=1,9)
691 WRITE(6,150) PL,PI,SI,SL,S9,S0,S1,S2,S6
692 WRITE(6,160)
693 WRITE (6,140)
694 80 FORMAT(11A1)
695 70 FORMAT('1')
696 100 FORMAT(35X,'RELATIVE DENSITY')
697 110 FORMAT (17X,9(2X,'0','1',' ')/,14X,10('+-+---',' '))

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MK110290
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698	120	FORMAT(1X,'TIME',1X,F7.3,1X,61A1)	MK110890
699	130	FORMAT(14X,10('+-----'),'+',/,16X,9(3X,'0.',1))	MK110900
700	140	FORMAT ('1')	MK110910
701	150	FORMAT(15X,'INPUT CHARACTER ',11A1,/) ,/,15X,' P - POWER',/,15X,' L -	MK110920
702	160	FORMAT(' < LOG. OF POWER',/,15X,' Q - PRECURSOR DENSITY',/,15X,' R - REACTIV	MK110930
		<ITY',/,16X,'NEGATIVE VALUE PLOTTED WITH THE AXIS IN THE CENTER',	MK110940
		<////)	MK110950
703		STOP	MK110960
704		END	MK110970
			MK110980

//DATA

***** KINETICS MODULE 1 *****

INPUT PARAMETER

TYPE OF FUEL U-239

NUMBER OF DELAYED NEUTRON = 3

NEUTRON GENERATION TIME = 0.00010 (SEC)

TYPE OF REACTIVITY = 2

RO = 0.250 82 = 5.00000

TOTAL TIME USED = 1.000 (SEC) , TIME INTERVAL = 0.010 (SEC)

OUTPUT OPTION = 2

INITIAL RELATIVE FLUX 1.000

***** END OF INPUT DATA *****

SIX GROUP LAMBDA ARE 0.0128 0.0301 0.1240 0.3250 1.1200 2.6900

SIX GROUP BETA ARE 0.00009 0.00080 0.00057 0.00088 0.00023 0.00012

THE AVERAGE BETA AND LAMBDA

B(1) = 0.00000 LAMBDA(1) = 0.02636

B(2) = 0.00145 LAMBDA(2) = 0.19854

B(3) = 0.00035 LAMBDA(3) = 1.39571

COEFFICIENTS OF INHOUR FORMULA

A(1) = -0.0000E 00
 A(2) = 0.3051E-03
 A(3) = 0.3606E-02
 A(4) = 0.2862E-02
 A(5) = 0.1200E-03

THE LARGEST EIGEN-VALUE = 0.0000E 00

NO ()	TIME (SEC)	R (%)	POWER (-)	3 GROUP DELAYED NEUTRON (RELATIVE)		
1	0.000	0.00	0.1000E 01	0.3411E 03	0.7303E 02	0.2515E 01

THE G-MATRIX

0.7634E 00	0.2310E-03	0.1740E-03	0.1223E-01			
0.8991E-01	0.9997E 00	0.0000E 00	0.0000E 00			
0.1448E 00	0.0000E 00	0.9990E 00	0.0000E 00			
0.3486E-01	0.0000E 00	0.0000E 00	0.9661E 00			
2	0.000	0.00	0.1000E 01	0.3411E 03	0.7303E 02	0.2515E 01
3	0.010	0.01	0.1003E 01	0.3411E 03	0.7303E 02	0.2515E 01
4	0.020	0.02	0.1008E 01	0.3411E 03	0.7303E 02	0.2515E 01
5	0.030	0.04	0.1015E 01	0.3411E 03	0.7303E 02	0.2515E 01
6	0.040	0.05	0.1024E 01	0.3411E 03	0.7303E 02	0.2516E 01
7	0.050	0.06	0.1033E 01	0.3411E 03	0.7303E 02	0.2517E 01
8	0.060	0.07	0.1043E 01	0.3411E 03	0.7304E 02	0.2518E 01
9	0.070	0.09	0.1055E 01	0.3411E 03	0.7304E 02	0.2519E 01
10	0.080	0.10	0.1066E 01	0.3411E 03	0.7305E 02	0.2521E 01
11	0.090	0.11	0.1078E 01	0.3412E 03	0.7306E 02	0.2523E 01
12	0.100	0.12	0.1091E 01	0.3412E 03	0.7307E 02	0.2526E 01
13	0.110	0.13	0.1103E 01	0.3412E 03	0.7309E 02	0.2529E 01
14	0.120	0.14	0.1116E 01	0.3412E 03	0.7310E 02	0.2532E 01
15	0.130	0.15	0.1129E 01	0.3412E 03	0.7312E 02	0.2536E 01
16	0.140	0.16	0.1143E 01	0.3412E 03	0.7314E 02	0.2540E 01
17	0.150	0.17	0.1156E 01	0.3412E 03	0.7316E 02	0.2545E 01
18	0.160	0.18	0.1169E 01	0.3412E 03	0.7318E 02	0.2550E 01
19	0.170	0.19	0.1182E 01	0.3412E 03	0.7320E 02	0.2555E 01
20	0.180	0.20	0.1195E 01	0.3413E 03	0.7323E 02	0.2561E 01
21	0.190	0.20	0.1207E 01	0.3413E 03	0.7326E 02	0.2567E 01
22	0.200	0.21	0.1220E 01	0.3413E 03	0.7329E 02	0.2574E 01
23	0.210	0.22	0.1232E 01	0.3413E 03	0.7332E 02	0.2581E 01
24	0.220	0.22	0.1244E 01	0.3413E 03	0.7335E 02	0.2588E 01
25	0.230	0.23	0.1255E 01	0.3414E 03	0.7338E 02	0.2595E 01
26	0.240	0.23	0.1266E 01	0.3414E 03	0.7342E 02	0.2603E 01
27	0.250	0.24	0.1276E 01	0.3414E 03	0.7346E 02	0.2611E 01

28	0.260	0.24	0.1280E	01	0.3414E	03	0.7350E	02	0.2619E	01
29	0.270	0.24	0.1295E	01	0.3414E	03	0.7354E	02	0.2628E	01
30	0.280	0.25	0.1303E	01	0.3415E	03	0.7358E	02	0.2637E	01
31	0.290	0.25	0.1311E	01	0.3415E	03	0.7362E	02	0.2646E	01
32	0.300	0.25	0.1318E	01	0.3415E	03	0.7367E	02	0.2655E	01
33	0.310	0.25	0.1324E	01	0.3416E	03	0.7371E	02	0.2664E	01
34	0.320	0.25	0.1329E	01	0.3416E	03	0.7376E	02	0.2673E	01
35	0.330	0.25	0.1333E	01	0.3416E	03	0.7380E	02	0.2682E	01
36	0.340	0.25	0.1336E	01	0.3416E	03	0.7385E	02	0.2691E	01
37	0.350	0.25	0.1338E	01	0.3417E	03	0.7390E	02	0.2701E	01
38	0.360	0.24	0.1339E	01	0.3417E	03	0.7394E	02	0.2710E	01
39	0.370	0.24	0.1339E	01	0.3417E	03	0.7399E	02	0.2719E	01
40	0.380	0.24	0.1338E	01	0.3418E	03	0.7404E	02	0.2728E	01
41	0.390	0.23	0.1336E	01	0.3418E	03	0.7409E	02	0.2737E	01
42	0.400	0.23	0.1333E	01	0.3418E	03	0.7413E	02	0.2746E	01
43	0.410	0.22	0.1329E	01	0.3419E	03	0.7418E	02	0.2754E	01
44	0.420	0.22	0.1324E	01	0.3419E	03	0.7422E	02	0.2762E	01
45	0.430	0.21	0.1318E	01	0.3419E	03	0.7427E	02	0.2770E	01
46	0.440	0.20	0.1311E	01	0.3419E	03	0.7431E	02	0.2778E	01
47	0.450	0.19	0.1303E	01	0.3420E	03	0.7435E	02	0.2785E	01
48	0.460	0.19	0.1295E	01	0.3420E	03	0.7439E	02	0.2792E	01
49	0.470	0.18	0.1285E	01	0.3420E	03	0.7443E	02	0.2798E	01
50	0.480	0.17	0.1275E	01	0.3420E	03	0.7447E	02	0.2804E	01
51	0.490	0.16	0.1264E	01	0.3421E	03	0.7451E	02	0.2810E	01
52	0.500	0.15	0.1253E	01	0.3421E	03	0.7455E	02	0.2815E	01
53	0.510	0.14	0.1240E	01	0.3421E	03	0.7458E	02	0.2819E	01
54	0.520	0.13	0.1228E	01	0.3421E	03	0.7461E	02	0.2824E	01
55	0.530	0.12	0.1215E	01	0.3422E	03	0.7464E	02	0.2827E	01
56	0.540	0.11	0.1201E	01	0.3422E	03	0.7467E	02	0.2830E	01
57	0.550	0.10	0.1188E	01	0.3422E	03	0.7469E	02	0.2833E	01
58	0.560	0.09	0.1174E	01	0.3422E	03	0.7472E	02	0.2835E	01
59	0.570	0.07	0.1160E	01	0.3422E	03	0.7474E	02	0.2837E	01
60	0.580	0.06	0.1145E	01	0.3422E	03	0.7476E	02	0.2838E	01
61	0.590	0.05	0.1131E	01	0.3422E	03	0.7478E	02	0.2838E	01
62	0.600	0.04	0.1117E	01	0.3423E	03	0.7479E	02	0.2838E	01
63	0.610	0.02	0.1102E	01	0.3423E	03	0.7481E	02	0.2838E	01
64	0.620	0.01	0.1088E	01	0.3423E	03	0.7482E	02	0.2837E	01
65	0.630	-0.00	0.1074E	01	0.3423E	03	0.7483E	02	0.2836E	01
66	0.640	-0.01	0.1060E	01	0.3423E	03	0.7483E	02	0.2834E	01
67	0.650	-0.03	0.1046E	01	0.3423E	03	0.7484E	02	0.2832E	01
68	0.660	-0.04	0.1033E	01	0.3423E	03	0.7484E	02	0.2829E	01
69	0.670	-0.05	0.1020E	01	0.3423E	03	0.7484E	02	0.2826E	01
70	0.680	-0.06	0.1007E	01	0.3423E	03	0.7484E	02	0.2822E	01
71	0.690	-0.08	0.9947E	00	0.3423E	03	0.7484E	02	0.2818E	01
72	0.700	-0.09	0.9826E	00	0.3423E	03	0.7483E	02	0.2814E	01
73	0.710	-0.10	0.9708E	00	0.3423E	03	0.7483E	02	0.2809E	01
74	0.720	-0.11	0.9595E	00	0.3423E	03	0.7482E	02	0.2804E	01
75	0.730	-0.12	0.9485E	00	0.3423E	03	0.7481E	02	0.2798E	01
76	0.740	-0.13	0.9380E	00	0.3423E	03	0.7480E	02	0.2793E	01
77	0.750	-0.14	0.9279E	00	0.3423E	03	0.7479E	02	0.2787E	01
78	0.760	-0.15	0.9182E	00	0.3423E	03	0.7477E	02	0.2780E	01
79	0.770	-0.16	0.9089E	00	0.3423E	03	0.7476E	02	0.2774E	01
80	0.780	-0.17	0.9001E	00	0.3423E	03	0.7474E	02	0.2767E	01
81	0.790	-0.19	0.8918E	00	0.3422E	03	0.7472E	02	0.2760E	01
82	0.800	-0.19	0.8838E	00	0.3422E	03	0.7471E	02	0.2753E	01
83	0.810	-0.20	0.8764E	00	0.3422E	03	0.7468E	02	0.2745E	01
84	0.820	-0.20	0.8693E	00	0.3422E	03	0.7466E	02	0.2738E	01
85	0.830	-0.21	0.8627E	00	0.3422E	03	0.7464E	02	0.2730E	01
86	0.840	-0.22	0.8566E	00	0.3422E	03	0.7462E	02	0.2723E	01
87	0.850	-0.22	0.8509E	00	0.3422E	03	0.7459E	02	0.2715E	01

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88	0.860	-0.23	0.8456E 00	0.3422E 03	0.7457E 02	0.2707E 01
89	0.870	-0.23	0.8408E 00	0.3421E 03	0.7454E 02	0.2699E 01
90	0.880	-0.24	0.8364E 00	0.3421E 03	0.7452E 02	0.2691E 01
91	0.890	-0.24	0.8324E 00	0.3421E 03	0.7449E 02	0.2682E 01
92	0.900	-0.24	0.8288E 00	0.3421E 03	0.7446E 02	0.2674E 01
93	0.910	-0.25	0.8257E 00	0.3421E 03	0.7444E 02	0.2666E 01
94	0.920	-0.25	0.8230E 00	0.3421E 03	0.7441E 02	0.2658E 01
95	0.930	-0.25	0.8207E 00	0.3421E 03	0.7438E 02	0.2650E 01
96	0.940	-0.25	0.8189E 00	0.3420E 03	0.7435E 02	0.2642E 01
97	0.950	-0.25	0.8174E 00	0.3420E 03	0.7432E 02	0.2634E 01
98	0.960	-0.25	0.8164E 00	0.3420E 03	0.7429E 02	0.2626E 01
99	0.970	-0.25	0.8158E 00	0.3420E 03	0.7426E 02	0.2618E 01
100	0.980	-0.25	0.8155E 00	0.3420E 03	0.7424E 02	0.2610E 01

***** END OF CALCULATION *****

RELATIVE DENSITY

0.1 0.2 0.3 0.4 0.5 0.6 0.7 0.8 0.9

TIME	0.000			L	R		P		QI
TIME	0.010			L	R		P		QI
TIME	0.020			L	R		P		QI
TIME	0.030			L	R		P		QI
TIME	0.040			L	R		P		QI
TIME	0.050			L	R		P		QI
TIME	0.060			L	R		P		QI
TIME	0.070			L	R		P		QI
TIME	0.080			L	R		P		QI
TIME	0.090			L	R		P		QI
TIME	0.100			L	R		P		QI
TIME	0.110			L	R		P		QI
TIME	0.120			L	R		P		QI
TIME	0.130			L	R		P		QI
TIME	0.140			L	R		P		QI
TIME	0.150			L	R		P		QI
TIME	0.160			L	R		P		QI
TIME	0.170			L	R		P		QI
TIME	0.180			L	R		P		QI
TIME	0.190			L	R		P		QI
TIME	0.200			L	R		P		QI
TIME	0.210			L	R		P		QI
TIME	0.220			L	R		P		QI
TIME	0.230			L	R		P		QI
TIME	0.240			L	R		P		QI
TIME	0.250			L	R		P		QI
TIME	0.260			L	R		P		QI
TIME	0.270			L	R		P		QI
TIME	0.280			L	R		P		QI
TIME	0.290			L	R		P		QI
TIME	0.300			L	R		P		QI
TIME	0.310			L	R		P		QI
TIME	0.320			L	R		P		QI
TIME	0.330			L	R		P		QI
TIME	0.340			L	R		P		QI
TIME	0.350			L	R		P		QI
TIME	0.360			L	R		P		QI
TIME	0.370			L	R		P		QI
TIME	0.380			L	R		P		QI
TIME	0.390			L	R		P		QI
TIME	0.400			L	R		P		QI
TIME	0.410			L	R		P		QI
TIME	0.420			L	R		P		QI
TIME	0.430			L	R		P		QI
TIME	0.440			L	R		P		QI
TIME	0.450			L	R		P		QI
TIME	0.460			L	R		P		QI
TIME	0.470			L	R		P		QI
TIME	0.480			L	R		P		QI
TIME	0.490			L	R		P		QI
TIME	0.500			L	R		P		QI
TIME	0.510			L	R		P		QI
TIME	0.520			L	R		P		QI
TIME	0.530			L	R		P		QI
TIME	0.540			L	R		P		QI
TIME	0.550			L	R		P		QI

TIME	0.560
TIME	0.570
TIME	0.580
TIME	0.590
TIME	0.600
TIME	0.610
TIME	0.620
TIME	0.630
TIME	0.640
TIME	0.650
TIME	0.660
TIME	0.670
TIME	0.680
TIME	0.690
TIME	0.700
TIME	0.710
TIME	0.720
TIME	0.730
TIME	0.740
TIME	0.750
TIME	0.760
TIME	0.770
TIME	0.780
TIME	0.790
TIME	0.800
TIME	0.810
TIME	0.820
TIME	0.830
TIME	0.840
TIME	0.850
TIME	0.860
TIME	0.870
TIME	0.880
TIME	0.890
TIME	0.900
TIME	0.910
TIME	0.920
TIME	0.930
TIME	0.940
TIME	0.950
TIME	0.960
TIME	0.970
TIME	0.980
TIME	0.990

0.1 0.2 0.3 0.4 0.5 0.6 0.7 0.8 0.9.
INPUT CHARACTER: -+ LRQPS

P - POWER
L - LOG. OF POWER
Q - PRECURSOR DENSITY
R - REACTIVITY

NEGATIVE VALUE PLOTTED WITH THE AXIS IN THE CENTER

REACTOR DYNAMICS MODULE, RD-2
REACTOR KINETICS WITH FEEDBACK

by

Ronald J. Onega

The University gratefully acknowledges the support of the
Division of Higher Education of the National Science Foundation
for support of this work performed under Grant GZ-2888 and the
support of Duke Power Company, North Carolina Power and Light Company,
and Virginia Electric and Power Company.

Project Director: Milton C. Edlund

ACKNOWLEDGMENT

The author would like to extend his appreciation and thanks to Mr. Tjeri Surjanto who did the programming for this module.

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REACTOR DYNAMICS MODULE, RD-2

REACTOR KINETICS WITH FEEDBACK

2.1 Object of Module

The object of this module is to:

- (1) Examine the temperature feedback mechanism of a PWR and
- (2) Solve the one delayed neutron model with temperature feedback for a step insertion and a ramp insertion of reactivity.

The time dependence of a reactor, taking the feedback mechanisms into account, is relatively difficult. We will consider a PWR core with a two path feedback. The reactivity is diminished as the temperature of the fuel increases due to the Doppler broadening of the resonances. This feedback is instantaneous since the temperature increase follows the power generated immediately. The second feedback path is that of the moderator temperature coefficient. As the moderator temperature increases, the number density decreases and the neutron mean free path increases so that leakage increases and reactivity decreases.

We will be concerned about the stability of the reactor to a limited degree. The dynamic response depends upon the magnitude of the temperature coefficients as well as that of the signs. For a given reactor design, i.e., a given life-time power level, the reactor may or may not be stable for a given set of reactivity coefficients.

The core region is the only one of interest in this module. The rest of the primary loop as well as the secondary loop is treated in an overall dynamics module for a PWR.

Also, all of our analysis will be fundamental mode analysis. The physical phenomena are taking place so slowly that the higher harmonics of the flux distribution are all dying out so rapidly that we only need to consider the lowest or fundamental mode.

The thermal analysis really should proceed by the solution of the space-time heat conduction equation. This is a very complicated procedure and would also mean that spatial effects of the kinetics equations should be taken into account. We, therefore, will assume only a lumped parameter model and will obtain the time dependence of a reactor which is really one with the average properties of the reactor under consideration.

The program name is FUMOTEM which is an acronym for "Fundamental Mode Kinetics with Temperature Feedback."

There are four types of reactivity inputs that the program can accommodate with NRO = 1, 2, 3 or 4 respectively:

$$\begin{aligned}
 1) \rho_o(t) &= \rho_o & 0 \leq t \leq t_r \\
 &= 0 & \text{otherwise} \\
 2) \rho_o(t) &= \rho_o(1 + a t) & 0 \leq t \leq t_r \\
 &= \rho_o(1 + a t_r) & \text{otherwise} \\
 3) \rho_o(t) &= \rho_o \cos a t \\
 \text{or} \\
 4) \rho_o(t) &= \rho_o \sin a t
 \end{aligned}$$

The feedback reactivity is taken to be zero for times $t > t_r$. One can feed in a t_r greater than the calculation time however. So for $t > t_r$, the only reactivity is due to feedback effects.

2.2 The Feedback Model

We will use the one group delayed neutron model to describe the core neutronics. The kinetics equations are

$$\frac{dP(t)}{dt} = \frac{\rho(t) - \beta}{\Lambda} P(t) + \lambda Q(t) \quad (2.2.1)$$

and

$$\frac{dQ(t)}{dt} = \frac{\beta}{\Lambda} P(t) - \lambda Q(t) \quad (2.2.2)$$

where

$P(t)$ = The total reactor power (Megawatts)

$Q(t)$ = The power equivalence of the delayed neutron precursors (Megawatts).

Now we let ΔT_M be the deviation of the spatially averaged moderator temperature from its equilibrium value, i.e.,

$$\Delta T_M(t) = T_M(t) - T_{M0} \quad (2.2.3)$$

and similarly for the fuel temperature T_F we have

$$\Delta T_F(t) = T_F(t) - T_{F0} \quad (2.2.4)$$

where T_{M0} and T_{F0} are the equilibrium moderator and fuel temperatures respectively. Also, we let α_M and α_F be the moderator and fuel temperature coefficients of reactivity. Then (1)

$$\rho(t) = \rho_0(t) + \alpha_M \Delta T_M(t) + \alpha_F \Delta T_F(t) \quad (2.2.5)$$

where generally, α_M and α_F will be negative or at least their sum is negative.

The temperature coefficient of reactivity for the moderator is

$$\alpha_M \equiv \frac{\partial \rho}{\partial T_M} \approx \frac{1}{k_{eff}} \frac{\partial k_{eff}}{\partial T_M} \quad (2.2.6)$$

and

$$\alpha_F = \frac{\partial \rho}{\partial T_F} \approx \frac{1}{k_{eff}} \frac{\partial k_{eff}}{\partial T_F} \quad (2.2.7)$$

The thermal analysis of the core must now be considered and connected to the neutronics. The heat generated depends upon the fission rate or the power. We will look at the fuel temperature averaged over the core as well as an averaged coolant temperature. We will ignore the cladding of the fuel pins.

The heat balance equation for reactor fuel is

$$\left[\begin{array}{l} \text{Rate at which the} \\ \text{internal energy of} \\ \text{the fuel changes} \end{array} \right] = \left[\begin{array}{l} \text{Production rate} \\ \text{of the energy} \\ \text{in the fuel} \end{array} \right] - \left[\begin{array}{l} \text{Rate at which} \\ \text{heat is conducted} \\ \text{out of the fuel} \end{array} \right]$$

or

$$\rho_F C_F V_F \frac{dT_F(t)}{dt} = P(t) - 4\pi k_{FF} L_F (T_F - T_R) N \quad (2.2.8)$$

where

- ρ_F = density of fuel (lb/ft³)
- C_F = specific heat of fuel (Btu/lb-°F)
- V_F = volume of fuel (ft³)
- T_F = average temperature of the fuel (°F)
- $P(t)$ = total power of reactor (Btu/hr)

- k_F = thermal conductivity of fuel ($\frac{\text{Btu}}{\text{hr-ft-}^\circ\text{F}}$)
 L_F = length of fuel pin (ft)
 T_O = centerline temperature of the fuel ($^\circ\text{F}$)
 T_R = temperature of the fuel pellets at the outer edge (pellet - water interface) ($^\circ\text{F}$)
 N = total number of fuel pins in the reactor.

The expression for Equation (2.2.8) was obtained from El-Wakil, "Nuclear Heat Transport", page 123, equation 5-48.

The temperatures are averaged over the fuel pins. If we number each of the pins in the core, then the centerline temperature is

$$T_O = \frac{1}{N} (T_{O1} + T_{O2} + \dots + T_{ON})$$

where T_{O1} is the centerline temperature for the i th pin. T_F is defined similarly so that it is the spatially averaged temperature of the "average fuel pin."

Equation (2.2.8) contains the centerline temperature and the temperature at the edge of the pellet T_R which must be eliminated. We assume that a parabolic temperature distribution holds even in the transient situation (really the transient heat conduction equation holds here) so that for one fuel pin we have

$$T(r) = T_O - \frac{P(t)r^2}{4k_F V_F} \quad (2.2.9)$$

where $T(r)$ is the temperature of the fuel pin a distance r from the center of the pin. The average fuel temperature is then

$$\begin{aligned}
 T_F(t) &= \frac{2\pi}{\pi R_F} \int_0^{R_F} r dr \left[T_O - \frac{P(t)r^2}{4k_F V_F} \right] \\
 &= T_O - \frac{P(t)}{8k_F V_F} R_F^2.
 \end{aligned} \tag{2.2.10}$$

Eliminating the centerline temperatures, we have, using (2.2.10),

$$T_F(t) = T_R(t) + \frac{R_F^2 P(t)}{8k_F V_F} \tag{2.2.11}$$

Equation (2.2.8) can now be written as

$$\begin{aligned}
 \rho_F C_F V_F \frac{dT_F(t)}{dt} &= P(t) - 4\pi k_{F-F} L_F N \left[T_F(t) + \frac{P(t) R_F^2}{8k_F V_F} - T_R(t) \right] \\
 &= \left[1/2 \right] P(t) - 4\pi k_{F-F} L_F N T_F(t) + 4\pi k_{F-F} L_F N T_R(t)
 \end{aligned} \tag{2.2.12}$$

The wall temperature of the pellets $T_R(t)$ is connected to the coolant temperature $T_M(t)$ since

$$T_R(t) - T_M(t) = \frac{P(t)}{h_T A_F} \tag{2.2.13}$$

where A_F is the total area of the fuel and h_T is the heat transfer coefficient for the fuel water interface (Btu/(Ft², hr-°F)).

The energy balance for the water in the core is

$$\left[\text{Heat stored in water in core} \right] = \left[\text{Heat conducted in from fuel} \right] - \left[\text{Heat transferred out of reactor core} \right]$$

or mathematically,

$$\rho_M C_M V_M \frac{dT_M(t)}{dt} = 2\pi R_F L_F h_T N [T_R(t) - T_M(t)] + \dot{m}_M [C_{M1} T_{M1} - C_{M2} T_{M2}] \quad (2.2.14)$$

where \dot{m}_M is the mass flow rate of the water through the core and $T_M(t)$ is the average moderator temperature. Also, assume $C_{M1} = C_{M2}$.

We will assume \dot{m}_M is an input, as is T_{M1} , the moderator inlet temperature. The moderator outlet temperature is related to the inlet and average temperatures as

$$T_{M2}(t) = 2T_M(t) - T_{M1}$$

so that Equation (2.2.14) becomes (using Equation (2.2.13))

$$\begin{aligned} \rho_M C_M V_M \frac{dT_M(t)}{dt} &= 2\pi R_F L_F h_T \cdot \left[\frac{P(t)}{h_T A_F} \right] + \dot{m}_M C_M [T_{M1} - T_M(t)] \\ &= P(t) + 2 \dot{m}_M C_M T_{M1} - 2 \dot{m}_M C_M T_M(t) \end{aligned} \quad (2.2.15)$$

This can be rewritten as

$$\frac{dT_M(t)}{dt} = \frac{P(t)}{\rho_M C_M V_M} - \frac{\dot{m}_M}{\rho_M V_M} T_M(t) + \frac{\dot{m}_M}{\rho_M V_M} T_{M1} \quad (2.2.16)$$

The area of the fuel is $2NL_F \pi R_F$ and the fuel volume is $\pi R_F^2 L_F N$.

In summary, the equations we must solve are Equations (2.2.1), (2.2.2), (2.2.12) and (2.2.16), the last two of which are written

$$\rho_F C_F V_F \frac{dT_F}{dt} = (1/2)P(t) - 4\pi k_F L_F N T_F(t) + 4\pi k_F L_F N T_M(t), \quad (2.2.17)$$

and

$$\frac{dT_M(t)}{dt} = \frac{P(t)}{\rho_M C V_M} - \frac{2\dot{m}_M}{\rho_M V_M} T_M(t) + \frac{2}{\rho_M V_M} \dot{m}_M T_{M1} \quad (2.2.18)$$

If we put these equations into matrix form we have

$$\frac{d}{dt} \begin{bmatrix} P(t) \\ Q(t) \\ T_F(t) \\ T_M(t) \end{bmatrix} = \begin{bmatrix} \frac{\rho - \beta}{\Lambda} & \lambda & 0 & 0 \\ \frac{\beta}{\Lambda} & -\lambda & 0 & 0 \\ \frac{1}{2 \rho_F C V_F} & 0 & -\frac{4\pi K_F L_F N}{\rho_F C V_F} & \frac{4\pi K_F L_F N}{\rho_F C V_F} \\ \frac{1}{\rho_M C V_M} & 0 & 0 & -\frac{2\dot{m}_M}{\rho_M V_M} \end{bmatrix} \mathbf{x}$$

$$\begin{bmatrix} P(t) \\ Q(t) \\ T_F(t) \\ T_M(t) \end{bmatrix} + \begin{bmatrix} 0 \\ 0 \\ 0 \\ \frac{2}{\rho_M V_M} \dot{m}_M T_{M1}(t) \end{bmatrix} \quad (2.2.19)$$

or

$$\frac{d \underline{\phi}(t)}{dt} = \underline{A} \underline{\phi} + \underline{B} \quad (2.2.20)$$

where the \underline{A} is dependent upon the temperatures themselves. Thus, Equation (2.2.20) is non-linear.

The linearization of Equation (2.2.19) (or (2.2.20)) can be achieved assuming that we can look at changes about some operating point. Expanding about the operating point we have

$$P(t) = P^{\circ} + \Delta P \quad (2.2.21)$$

$$Q(t) = Q^{\circ} + \Delta Q \quad (2.2.22)$$

$$T_F(t) = T_F^{\circ} + \Delta T_F \quad (2.2.23)$$

$$T_M(t) = T_M^{\circ} + \Delta T_M \quad (2.2.24)$$

We assume that P° etc. are independent of time so

$$\frac{d \Delta \phi}{dt} = (\underline{A}^{\circ} + \Delta \underline{A}) (\underline{\phi}^{\circ} + \Delta \underline{\phi}) + \underline{B}$$

or neglecting the $\Delta \underline{A} \cdot \Delta \underline{\phi}$ term we have

$$\frac{d}{dt} \Delta \underline{\phi} = \underline{A}^{\circ} \underline{\phi}^{\circ} + \underline{A}^{\circ} \Delta \underline{\phi} + \Delta \underline{A} \underline{\phi}^{\circ} + \underline{B}$$

or

$$\frac{d}{dt} \Delta \underline{\phi} = \underline{A}^{\circ} \Delta \underline{\phi} + \Delta \underline{A} \underline{\phi}^{\circ} \quad (2.2.25)$$

since

$$\underline{A}^{\circ} \underline{\phi}^{\circ} + \underline{B} = \underline{0}.$$

Writing Equation (2.2.25) out explicitly, we have

$$\frac{d}{dt} \begin{bmatrix} \Delta P \\ \Delta Q \\ \Delta T_F \\ \Delta T_M \end{bmatrix} = \begin{bmatrix} \frac{\rho_o - \beta}{\Lambda} & \lambda & 0 & 0 \\ \frac{\beta}{\Lambda} & -\lambda & 0 & 0 \\ \frac{1}{2\rho_F C_F V_F} & 0 & \frac{-4\pi k_F L_F N}{\rho_F C_F V_F} & \frac{4\pi k_F L_F N}{\rho_F C_F V_F} \\ \frac{1}{\rho_M C_M V_M} & 0 & 0 & \frac{-2\pi m_M}{\rho_M V_M} \end{bmatrix} \begin{bmatrix} \Delta P \\ \Delta Q \\ \Delta T_F \\ \Delta T_M \end{bmatrix} + \begin{bmatrix} \frac{\alpha_F \Delta T_F + \alpha_M \Delta T_M}{\Lambda} & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} P^\circ \\ Q^\circ \\ T_F^\circ \\ T_M^\circ \end{bmatrix}$$

or

$$\frac{d}{dt} \begin{bmatrix} \Delta P \\ \Delta Q \\ \Delta T_F \\ \Delta T_M \end{bmatrix} = \begin{bmatrix} \frac{\rho_o - \beta}{\Lambda} & \lambda & \frac{\alpha_F P^\circ}{\Lambda} & \frac{\alpha_M P^\circ}{\Lambda} \\ \frac{\beta}{\Lambda} & -\lambda & 0 & 0 \\ \frac{1}{2\rho_F C_F V_F} & 0 & \frac{-4\pi k_F L_F N}{\rho_F C_F V_F} & \frac{4\pi k_F L_F N}{\rho_F C_F V_F} \\ \frac{1}{\rho_M C_M V_M} & 0 & 0 & \frac{-2\pi m_M}{\rho_M V_M} \end{bmatrix} \begin{bmatrix} \Delta P \\ \Delta Q \\ \Delta T_F \\ \Delta T_M \end{bmatrix} \quad (2.2.26)$$

Equation (2.2.26) is a linearized version of Equation (2.2.19). This equation is soluble by the ordinary method of eigenvalues. In matrix form, this can be written as

$$\frac{d\Delta \phi}{dt} = \underline{A}' \Delta \phi, \quad (2.2.27)$$

and \underline{A}' is given in terms of the equilibrium values.

The system of Equations (2.2.19) is highly non-linear because of the first element of the matrix \underline{A} . The method discussed in Kinetics Module 1, developed by Hansen, is still applicable even to non-linear systems. There is a slight modification that we must discuss and this will be outlined in a later section.

The overall heat transfer coefficient is also relatively difficult to obtain so we discuss this and related parameters in the next section.

Problem 2.2.1

Show that if T_o is the temperature in the center of a fuel pin and T_M is the moderator temperature, then

$$T_o = T_M + \frac{P(t)a^2}{4k_F V_F} + \frac{P(t)}{2V_F} \left[\frac{1}{k_{clad}} \ln \frac{b}{a} + \frac{1}{h_T b} \right]$$

The inner radius of the clad is "a" and the outer radius is "b".

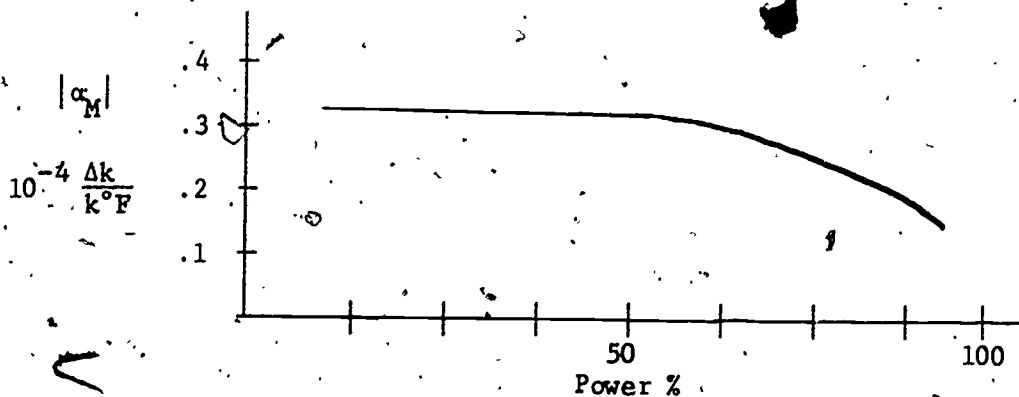
2.3 The Feedback Parameters

There are various parameters that must be obtained in order to numerically solve Equation (2.2.19). We list these parameters in order:

1. α_M - The moderator (coolant) temperature coefficient of reactivity.
2. α_F - The temperature coefficient of reactivity for the fuel.
3. C_{PF} - The total heat capacity of the fuel.
4. k_F - The thermal conductivity of the fuel.
5. h_T - The overall convective heat transfer coefficient of the fuel to moderator surface.

The flow rate F , the fuel volume V_F , etc. are all parameters that can be obtained from specifications of a particular reactor type and we shall not discuss these any further.

We first obtain a relation for α_M . The Final Safety Analysis Report of nuclear power plants generally have a curve of α_M as a function of power level for a given boron insertion and a critical rod insertion. The curves generally follow the pattern that α_M is constant from about 10% to 60% power and then decreases rather dramatically between 60 and 100% power. α_M is assumed to be negative.



We will therefore take α_M to be an input constant but an accurate analysis would necessitate a knowledge of the variation of α_M with power.

The $\alpha_F(P)$ is also rather difficult to calculate. To do this, we use Equation (2.2.7), i.e.

$$\alpha_F \approx \frac{1}{k_{eff}} \frac{\partial k_{eff}}{\partial T_F} \quad (2.3.1)$$

Now if we use the relation

$$k_{eff} = \eta p f \epsilon L_f L_{th} \quad (2.3.2)$$

then

$$\ln(k_{eff}) = \ln(\eta f \epsilon L_f L_{th}) + \ln(p)$$

and if we assume that the resonance escape probability p is the only factor which changes with the fuel temperature, then

$$\alpha_F = \frac{1}{k_{eff}} \frac{dk_{eff}}{dT_F} = \frac{1}{p} \frac{dp}{dT_F} \quad (2.3.3)$$

A standard expression for the resonance escape probability is (2)

$$p = e^{-\frac{N_F V_F}{\epsilon_M \Sigma_{SM-M}} R(T_F)} \quad (2.3.4)$$

and the resonance integral $R(T_F)$ is given by the empirical relation

$$R(T_F) = R(T_{FO}) [1 + \gamma(\sqrt{T_F} - \sqrt{T_{FO}})] \quad (2.3.5)$$

where T_{FO} is the equilibrium temperature of the fuel.

We now relate "p" to α_F in a way so that the α_F can be calculated. Let

$$a_5 \equiv \frac{N_F V_F}{\epsilon_M \Sigma_{SM} V_M} \quad (2.3.6)$$

then

$$\ln p = - a_5 R(T_F), \quad (2.3.7)$$

and

$$\alpha_F = - a_5 \frac{dR(T_F)}{dT_F} = - a_5 \frac{R(T_{FO}) \gamma}{2\sqrt{T_F}}. \quad (2.3.8)$$

IF we substitute T_{FO} into p we get

$$p(T_{FO}) = e^{-a_5 R(T_{FO})}$$

or

$$\ln \frac{1}{p(T_{FO})} = a_5 R(T_{FO}). \quad (2.3.9)$$

For an arbitrary value of temperature T_F we obtain

$$\ln \frac{1}{p(T_F)} = a_5 R(T_F). \quad (2.3.10)$$

Dividing Equation (2.3.9) by Equation (2.3.10) we get

$$\frac{R(T_{FO})}{R(T_F)} = \frac{\ln \frac{1}{p(T_{FO})}}{\ln \frac{1}{p(T_F)}} = + \frac{\ln p(T_{FO})}{\ln p(T_F)} \quad (2.3.11)$$

Inserting Equation (2.3.11) into Equation (2.3.8) we obtain

$$\alpha_F = - \frac{\gamma}{2\sqrt{T_F}} [a_5 R(T_{FO})] = - \frac{\gamma a_5}{2\sqrt{T_F}} R(T_F) \cdot \frac{\ln p(T_{FO})}{\ln p(T_F)}$$

Equation (2.3.10) yields

$$a_5 R(T_F) = - \ln p(T_F)$$

so

$$\begin{aligned} \alpha_F &= - \frac{\gamma}{2\sqrt{T_F}} \cdot [- \ln p(T_F)] \cdot \frac{\ln p(T_{FO})}{\ln p(T_F)} \\ &= - \frac{\gamma}{2\sqrt{T_F}} \frac{1}{\ln p(T_{FO})} \end{aligned} \quad (2.3.12)$$

In our module, we will assume that γ and $p(T_{FO})$ are read in.

The third item of discussion is C_{PF} , the total heat capacity of the fuel.

It is obvious that

$$C_{PF} = V_F \cdot C_{PF}^i \cdot \rho_F \quad (2.3.13)$$

where C_{PF}^i is the specific heat capacity of the fuel, i.e. the units are Btu/lbm-°F, ρ_F is the density of fuel and V_F is the volume of fuel. There is a small variation of C_{PF} with temperature but it is small for UO_2 in the region of interest. The value that will be of interest for us is 0.0590 Btu/lbm-°F.

The thermal conductivity of the fuel is a function of the temperature. El Wakil (3) lists values of k_F as a function of T_F . We use these values to obtain a polynomial regression of the k_F with the T_F .

The last quantity that we discuss is the overall convective heat transfer coefficient h_T . The heat transfer coefficient is defined by Newton's law of cooling. The relevant relation is

$$\frac{P(t)}{A_F} = h_T(T_F - T_M), \quad (2.3.14)$$

where A_F is the fuel area. There are many factors which influence h_T such as

- i) the temperature of the system
- ii) the heat flux
- iii) the physical properties of the moderating material
- iv) the geometrical shape of the cooling surface
- v) the flow rate of the coolant

In the PWR, the coolant flow is turbulent. Therefore, to obtain the heat transfer coefficient h_T we assume that,

$$h_T = \frac{C k_M Re^{0.8} Pr^{0.333}}{D_e} \quad (2.3.15)$$

where

D_e = The equivalent diameter of flow channels through the fuel rod bundles

k_M = The thermal conductivity of the moderator

C = The Colburn correction factor for fluid flow parallel to the tube bundles

Re = Reynolds number

Pr = Prandtl number of the coolant in the core

Equation (2.3.15) was a correlation recommended by Weisman. The Prandtl-number is taken from the 1967 ASME Steam Tables.

The equivalent diameter of the flow channel is shown in Fig. 2.3.2. If we let "a" be the area of the flow channel and P_w be the wetted perimeter, then

$$D_e = \frac{4A}{P_w} \quad (2.3.16)$$

For a typical PWR, D , the diameter of the fuel rods is 0.03583 ft and the pitch (the distance between centerlines of fuel rods) is designated as P_c . The Colburn correction factor is an empirical relation and is

$$C = 0.042 \frac{P_c}{D} - 0.024. \quad (2.3.17)$$

So we denote the heat transfer coefficient, which varies with power and therefore with time, as $h_T(t)$ and the equilibrium value as $h(0)$. Using Equation (2.3.15) we have

$$\frac{h_T(t)}{h_T(0)} = \frac{k_M(t) \text{Re}^{0.8} \text{Pr}^{1/3}}{k_M(0) \text{Re}^{0.8}(0) \text{Pr}^{1/3}(0)}$$

or

$$h_T(t) = \frac{h(0) k_M(t)}{k_M(0) \text{Pr}^{1/3}(0)} \frac{\rho_M(t)^{0.8}}{\mu(t)} \frac{\text{Pr}(t)^{1/3}}{\left(\frac{\rho_M(0)^{0.8}}{\mu(0)}\right)} \quad (2.3.18)$$

where we have used the fact that

$$\text{Re} = \frac{\rho_M v D_e}{\mu} \quad (2.3.19)$$

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$$\mu_M(T) = 8.5067 \times 10^{-1} - 1.7501 \times 10^{-3} T_M + 1.1142 \times 10^{-6} T_M^2 \quad (2.3.23)$$

and for the thermal conductivity of the fuel, we have

$$k_F(T) = 6.102141 - 4.636522 \times 10^{-3} T_F + 1.306299 \times 10^{-6} T_F^2 \quad (2.3.24)$$

The temperature is for the fuel in the case of k_F but the moderator temperature otherwise.

Problem 2.3.1

Given the following data:

$$R_F = 0.015042 \text{ ft}$$

$$\text{Fuel length} = 12.00 \text{ ft}$$

$$\text{Volume of reactor vessel} = 3643 \text{ ft}^3$$

$$\text{Number of fuel assemblies} = 145$$

$$\text{Fuel rods per assembly} = 208$$

$$\text{Rod pitch} = 0.04733 \text{ ft}$$

$$\rho = 43.214 \text{ lbm/ft}^3$$

$$v = 16.3 \text{ ft/sec}$$

$$\mu = 5.828 \times 10^{-5}$$

$$D_e = 0.04377 \text{ ft}$$

$$Re = 528,998$$

$$Pr = 1.01$$

$$k_M = 0.3010 \text{ Btu/hr-ft-}^\circ\text{F}$$

a) show the Colburn correction is 0.03148

b) $h = 8232 \text{ Btu/ft}^2\text{-hr-}^\circ\text{F}$.

2.4 Numerical Solution of the Dynamics Equations

The numerical integration of the differential equations can be carried out once the various parameters, such as the heat transfer coefficient, thermal conductivity, etc., have been obtained. We again write the equations so that we can proceed to use Hansen's method in their solution. FUMOTEM solves the kinetics equations using Hansen's method. The system of equations is

$$\frac{d\phi(t)}{dt} = \underline{A} \phi(t) + \underline{B} \quad (2.4.1)$$

where

$$\phi = \begin{bmatrix} P(t) \\ Q(t) \\ T_F(t) \\ T_M(t) \end{bmatrix} \quad \underline{B} = \begin{bmatrix} 0 \\ 0 \\ 0 \\ \frac{2 \cdot \dot{m}_M T_{M1}(t)}{\rho_M V_M} \end{bmatrix} \quad (2.4.2)$$

and

$$\underline{A} = \begin{bmatrix} \frac{\rho(t) - \beta}{\Lambda} & \lambda & 0 & 0 \\ \frac{\beta}{\Lambda} & -\lambda & 0 & 0 \\ \frac{1}{2\rho_F C_F V_F} & 0 & \frac{-4\pi k_F L N}{\rho_F C_F V_F} & \frac{4\pi k_F L N}{\rho_F C_F V_F} \\ \frac{1}{\rho_M C_M V_M} & 0 & 0 & \frac{-2 \cdot \dot{m}_M}{\rho_M V_M} \end{bmatrix} \quad (2.4.3)$$

subject to the constraint that

$$\rho(t) = \rho_0 + \alpha_M \Delta T_M + \alpha_F \Delta T_F. \quad (2.4.4)$$

All the parameters involved in Equations (2.4.1) - (2.4.4) are defined in section 2.2.

Hansen's method must be generalized slightly for the solution of Equation (2.4.1). We break the \underline{A} matrix into three parts:

$$\underline{A} = \underline{L} + \underline{D} + \underline{U}, \quad (2.4.5)$$

with

$$\underline{D}(t) = \begin{bmatrix} \frac{\rho(t) - \beta}{\Lambda} & 0 & 0 & 0 \\ 0 & -\lambda & 0 & 0 \\ 0 & 0 & \frac{-4\pi k_F L_F N}{\rho_F C_F V_F} & 0 \\ 0 & 0 & 0 & \frac{-2\dot{m}_M}{\rho_M V_M} \end{bmatrix}$$

$$\underline{L} = \begin{bmatrix} 0 & 0 & 0 & 0 \\ \frac{\beta}{\Lambda} & 0 & 0 & 0 \\ \frac{1}{2\rho_F C_F V_F} & 0 & 0 & 0 \\ \frac{1}{\rho_M C_M V_M} & 0 & 0 & 0 \end{bmatrix}$$

and

$$\underline{U} = \begin{bmatrix} 0 & \lambda & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & \frac{4\pi k_F L_F N}{\rho_F C_F V_F} \\ 0 & 0 & 0 & 0 \end{bmatrix}$$

Equation (2.4.1) can then be written as

$$\frac{d\phi}{dt} - \underline{D}(t) \phi(t) = (\underline{L} + \underline{U}) \phi(t) + \underline{B} \quad (2.4.6)$$

We wish to develop an iteration procedure for the solution of this system, so we begin the calculation at time t_0 and advance to time t_1 , and define

$$h = t_1 - t_0 \quad (2.4.7)$$

Before we use this relation though, let's multiply each term of Equation (2.4.6) by the integrating factor $e^{-\int_0^t \underline{D}(t') dt'}$ so

$$e^{-\int_0^t \underline{D}(t') dt'} \frac{d\phi}{dt} - e^{-\int_0^t \underline{D}(t') dt'} \underline{D}(t) \phi(t) = e^{-\int_0^t \underline{D}(t') dt'} (\underline{L} + \underline{U}) \phi(t) + e^{-\int_0^t \underline{D}(t') dt'} \underline{B}$$

or

$$\frac{d}{dt} \{ e^{-\int_0^t \underline{D}(t') dt'} \phi(t) \} = e^{-\int_0^t \underline{D}(t') dt'} (\underline{L} + \underline{U}) \phi(t) + e^{-\int_0^t \underline{D}(t') dt'} \underline{B} \quad (2.4.8)$$

The left side is an exact differential so we integrate from 0 to h and obtain

$$e^{-\int_0^t \underline{D}(t') dt'} \phi(t) \Big|_0^h = \int_0^h e^{-\int_0^t \underline{D}(t') dt'} (\underline{L} + \underline{U}) \phi(t) dt + \int_0^h e^{-\int_0^t \underline{D}(t') dt'} \underline{B}(t) dt$$

or

$$\begin{aligned} \phi(t_0 + h) = & e^{+\int_0^h \underline{D}(t') dt'} \phi(t_0) + \int_0^h e^{+\int_0^h \underline{D}(t') dt'} e^{-\int_0^t \underline{D}(t') dt'} (\underline{L} + \underline{U}) \phi(t_0 + \theta) d\theta \\ & + \int_0^h e^{+\int_0^h \underline{D}(t') dt'} - \int_0^t \underline{D}(t') dt' \underline{B}(t_0 + \theta) d\theta. \end{aligned} \quad (2.4.9)$$

Here

$$t_0 < \theta < t_1 = t_0 + h$$

and

$$d\theta = dt.$$

Notice in Equation (2.4.9) that

$$e^{\int_0^h \underline{D}(t') dt'} - \int_0^t \underline{D}(t') dt' = e^{\int_t^h \underline{D}(t') dt'} = e^{\int_{t_0+\theta}^h \underline{D}(t') dt'}$$

Therefore the equation that must be solved is

$$\phi(t_0 + h) = e^{\int_0^h \underline{D}(t') dt'} \phi(t_0) + \int_0^h e^{\int_{t_0+\theta}^h \underline{D}(t') dt'} (\underline{L} + \underline{U}) \Big|_{t_0+\theta} \phi(t_0 + \theta) dt$$

$$+ \int_0^h e^{\int_{t_0+\theta}^h \underline{D}(t') dt'} \underline{B}(t_0 + \theta) d\theta \quad (2.4.10)$$

Equation (2.4.10) is an integral equation so we must approximate its solution. To do this, we assume that

$$\underline{\phi}(t_0 + \theta) = e^{\omega_0 \theta} \underline{\phi}(t_0) \quad (2.4.11)$$

where ω_0 is the largest eigenvalue of the matrix \underline{A} evaluated at time t_0 . This means we must solve the equation

$$|\underline{A} - \omega \underline{I}| = 0$$

or

$$\det \begin{bmatrix} \frac{\rho(t_0) - \beta}{\Lambda} - \omega & \lambda & 0 & 0 \\ \frac{\beta}{\Lambda} & -\lambda - \omega & 0 & 0 \\ \frac{1}{2\rho_F C_F V} & 0 & \frac{-4\pi k_F L_F N}{\rho_F C_F V_F} - \omega & \frac{4\pi k_F L_F N}{\rho_F C_F V_F} \\ \frac{1}{\rho_M C_M V_M} & 0 & 0 & \frac{-2\dot{m}_M}{\rho_M V_M} - \omega \end{bmatrix} = 0, \quad (2.4.12)$$

to obtain the eigenvalues, ω_0 , ω_1 , ω_2 and ω_3 , where

$$\omega_0 > \omega_1 > \omega_2 > \omega_3$$

The ω 's are, of course, time dependent quantities.

We also look at $\underline{B}(t_0 + \theta)$ and assume that we can expand it in a Taylor series so that

$$\underline{B}(t_0 + \theta) = \underline{B}(t_0) + \theta \frac{d\underline{B}}{dt}(t_0) + \frac{1}{2} \theta^2 \frac{d^2 \underline{B}(t_0)}{d\theta^2} + \dots \quad (2.4.13)$$

and also

$$\underline{L}(t_0 + \theta) + \underline{U}(t_0 + \theta) = \underline{L}(t_0) + \underline{U}(t_0) + \dots \quad (2.4.14)$$

We shall keep only the 1st terms of these expansions.

Inserting Equations (2.4.14), (2.4.13) and (2.4.11) into Equation (2.4.10)

we get (after assuming that $\int_0^h \underline{D}(t') dt' = \underline{D}(t_0)h$)

$$\begin{aligned} \underline{\phi}(t_0 + h) &= e^{\underline{D}h} \underline{\phi}(t_0) + \int_0^h e^{\underline{D}(h-\theta)} (\underline{L} + \underline{U}) e^{\omega_0 \theta} \underline{\phi}(t_0) d\theta \\ &\quad + \int_0^h e^{\underline{D}(h-\theta)} \underline{B}(t_0) d\theta \\ &= e^{\underline{D}h} \underline{\phi}(t_0) + e^{\underline{D}h} (\omega_0 \underline{I} - \underline{D})^{-1} \left[e^{(\omega_0 \underline{I} - \underline{D})\theta} \right]_0^h (\underline{L} + \underline{U}) \underline{\phi}(t_0) \\ &\quad + e^{\underline{D}h} (-\underline{D})^{-1} e^{-\underline{D}\theta} \Big|_0^h \underline{B}(t_0); \end{aligned}$$

Putting the limits into this equation and simplifying we finally have

$$\begin{aligned} \underline{\phi}(t_0 + h) = & e^{\underline{D}h} \underline{\phi}(t_0) + (\omega_0 \underline{I} - \underline{D})^{-1} [e^{\omega_0 \underline{I}h} - e^{\underline{D}h}] [\underline{L}(t_0) + \underline{U}(t_0)] \underline{\phi}(t_0) \\ & + \underline{D}^{-1} (e^{\underline{D}h} - \underline{I}) \underline{B}(t_0) \end{aligned} \quad (2.4.15)$$

If we now set

$$\underline{\phi}(t_j) = \underline{\phi}_j$$

and

$$\underline{\phi}(t_j + h) = \underline{\phi}_{j+1}$$

Then Equation (2.4.15) becomes

$$\underline{\phi}_{j+1} = \underline{H} \underline{\phi}_j + \underline{R} \underline{B}_j \quad (2.4.16)$$

where

$$\underline{H} = e^{\underline{D}h} + (\omega_0 \underline{I} - \underline{D})^{-1} [e^{\omega_0 \underline{I}h} - e^{\underline{D}h}] (\underline{L} + \underline{U}) \quad (2.4.17)$$

with all quantities evaluated at time t_0 and

$$\underline{R} = \underline{D}^{-1} (e^{\underline{D}h} - \underline{I}) \quad (2.4.18)$$

We now write \underline{H} and \underline{R} explicitly for our problem. To do this, consider

Equation (2.4.17) and the explicit relation for \underline{L} and \underline{U} ; then

$$\underline{H} = \begin{bmatrix} H_1 & H_2 & 0 & 0 \\ H_3 & H_4 & 0 & 0 \\ H_5 & 0 & H_6 & H_7 \\ H_8 & 0 & 0 & H_9 \end{bmatrix} \quad (2.4.19)$$

where

$$H_1 = e^{\left[\frac{\rho(t_0) - \beta}{\Lambda} \right] h},$$

$$H_2 = \lambda \left[\frac{e^{\omega_0 h} - e^{\frac{\rho(t_0) - \beta}{\Lambda} h}}{\omega_0 - \frac{\rho(t_0) - \beta}{\Lambda}} \right],$$

$$H_3 = \frac{\beta}{\Lambda} \left[\frac{e^{\omega_0 h} - e^{-\lambda h}}{\omega_0 + \lambda} \right],$$

$$H_4 = e^{-\lambda h},$$

$$H_5 = \frac{1}{2\rho_F C_F V_F} \left[\frac{e^{\omega_0 h} - e^{\frac{4\pi k_F L_F N}{\rho_F C_F V_F} h}}{\omega_0 - \frac{4\pi k_F L_F N}{\rho_F C_F V_F}} \right],$$

$$H_6 = e^{-\frac{4\pi k_F L_F N}{\rho_F C_F V_F} h},$$

$$H_7 = \frac{4\pi k_F L_F N}{\rho_F C_F V_F} \left[\frac{e^{\omega_0 h} - e^{\frac{4\pi k_F L_F N}{\rho_F C_F V_F} h}}{\omega_0 - \frac{4\pi k_F L_F N}{\rho_F C_F V_F}} \right],$$

$$H_8 = \frac{1}{\rho_M C_M V_M} \begin{bmatrix} \omega_o h & -\frac{2\dot{m}_M h}{\rho_M V_M} \\ e^{\omega_o h} - e^{-\frac{2\dot{m}_M h}{\rho_M V_M}} \\ \omega_o & +\frac{2\dot{m}_M}{\rho_M V_M} \end{bmatrix}$$

and

$$H_9 = -\frac{2\dot{m}_M}{\rho_M V_M} \begin{bmatrix} \omega_o h & -\frac{2\dot{m}_M h}{\rho_M V_M} \\ e^{\omega_o h} - e^{-\frac{2\dot{m}_M h}{\rho_M V_M}} \\ \omega_o & +\frac{2\dot{m}_M}{\rho_M V_M} \end{bmatrix}$$

and the \underline{R} matrix is

$$\underline{R} = \begin{bmatrix} \frac{\Lambda}{\rho - \beta} \left(e^{\frac{\rho - \beta}{\Lambda}} - 1 \right) & 0 & 0 & 0 \\ 0 & \frac{1}{\lambda} (1 - e^{-\lambda h}) & 0 & 0 \\ 0 & 0 & -\frac{4\pi k_{FF} L_N h}{\rho_{FF} C_{FF} V_{FF}} \left(1 - e^{-\frac{\rho_{FF} C_{FF} V_{FF}}{4\pi k_{FF} L_N}} \right) & 0 \\ 0 & 0 & 0 & \frac{\rho_M V_M}{2\dot{m}_M} \left(1 - e^{-\frac{2\dot{m}_M}{\rho_M V_M} h} \right) \end{bmatrix} \quad (2.4.20)$$

The basic iteration procedure in FUMOTEM is as follows:

1. Calculate the pertinent reactor parameters such as α_F , h_T , \dot{m} , ρ_M , μ , C_M , k_F , and k_M for the guessed initial conditions. The parameters read in are T_{M1} , α_M , h , P_o , v (the coolant velocity), some optional settings for print out and the time length the program is to operate.

2. Construct the vector $\phi(0)$. We can choose the parameters of interest, i.e.

$$P(0) = 500 \text{ MW},$$

and

$$Q(0) = \frac{\beta}{\Lambda \lambda} P(0) \text{ MW}.$$

The fuel temperature $T_F(0)$ and the moderator temperature $T_M(0)$ are completely determined by $P(0)$ and the above read-in parameters. $T_F(0)$ and $T_M(0)$ are difficult to obtain since C_M , k_M , μ etc. depend on the temperature. This difficulty is overcome by using an iteration procedure. The method is as follows:

- a) From $P(0)$ and initially guessed values $T_F^{(0)}(0)$ and $T_M^{(0)}(0)$, calculate \dot{m}_M from the relation

$$\dot{m}_M(T_M) = \rho_M(T_M) A v. \quad (2.4.21)$$

Equation (2.3.20) is used to obtain $\rho_M(T_M)$ from the guessed moderator temperature $T_M^{(0)}(0)$. Equation (2.2.15) is now used to obtain an improved guessed power $P^{(1)}(0)$,

$$P^{(1)}(0) = 2\dot{m}_M C_M (T_M^{(0)}(0) - T_{M1}). \quad (2.4.22)$$

- b) The fractional difference between the actual power and $P^{(1)}(0)$ is

$$E \equiv \frac{|P^{(1)}(0) - P(0)|}{P(0)} \quad (2.4.23)$$

If $E \leq \Delta$ ($\Delta = 0.01$ and is an input) then the guessed temperature $T_M^{(0)}(0)$ is the correct moderator temperature. If

$$E > \Delta$$

(2.4.24)

then change $T_M^{(0)}(0)$ to $T_M^{(1)}(0)$ and repeat the above until convergence is achieved.

- c) The $T_M(0)$ is now used in Equation (2.2.17) which, for equilibrium, becomes, after solving for $T_F(0)$,

$$T_F(0) = T_M(0) + \frac{1}{8\pi k_{F,N}} P(0). \quad (2.4.25)$$

This step also requires an iterative procedure since k_F is a function of the fuel temperature as seen from Equation (2.4.25).

3. Determine the largest eigenvalue of the equation

$$|\underline{A} - \omega \underline{I}| = 0.$$

This will be the solution of a 4 x 4 determinant which is rather easy on the computer. From this determine the largest root ω_0 . The Newton-Raphson method is used to calculate all four roots of this polynomial and then the largest root is picked by comparison of the roots. The subroutine POLRT is used for the determination of ω_0 for each time step. Module RD-1 describes the Newton-Raphson method.

4. Construct the $\underline{H}(t_0)$ matrix using Equation (2.4.19) and $\underline{R}(t_0)$ using (2.4.20)
5. Determine the vector $\underline{\phi}_1$ from Equation (2.4.16), i.e.

$$\underline{\phi}_1 = \underline{H}(t_0)\underline{\phi}_0 + \underline{R}(t_0)\underline{B}_0.$$

6. Repeat the above procedure using step 3 and continue as long as required to achieve the solution over the time domain of interest. The time steps are chosen by the criterion $h = \frac{1}{\omega_0}$ and h is constrained to be in the interval $0.005 \leq h \leq 0.05$ sec.

2.5 Input-Output Data for Code FUMOTEM

The input data required for the program are presented below:

Data Card Number	Format Statement Number	Format	Unit	Variable Name	Description
1	20	I2	--	NOPLT	Plotting option 1 -- plot 0 -- no plot
2	30	F10.5	-- ⁻¹	BETA	Delayed neutron fraction
		F10.5	sec ⁻¹	X	Delayed neutron decay constant
		F10.7	sec ⁻¹	XL	Neutron generation time
		F10.7	\$	RO	reactivity, ρ_0
		I10	--	NRO	option for the type of reactivity insertion
		F10.5	sec	RTIME	Time duration reactivity is inserted for ramp input (NRO=2)
		F10.5	sec ⁻¹	A	reactivity insertion rate
3	40	F10.6	°F ⁻¹	AM	Moderator temperature coefficient of reactivity
		F10.6	(°F) ^{1/2}	G	Constant for resonance capture
		F10.6	--	PTFO	Resonance escape probability, P_0
4	50	F10.5	ft	RF	Fuel radius, R_F
		F10.5	ft	PC	Distance between fuel pin centers, P
		F10.5	$\frac{\text{Btu}}{\text{lb-°F}}$	CPFX	Specific heat of the fuel, C_F
		F10.2	lb/ft ³	FDENS	Fuel density, ρ_F
		F10.2	ft	FH	Fuel height, L
		I10	--	NA	Number of fuel assemblies
		I10	--	NFRPA	Number of fuel rods per assembly
5	60	F10.2	ft/sec	V	Average velocity of coolant, v
		F10.2	°F	TMI	Inlet coolant temperature to the core T_{M1}
6	70	F10.4	sec	TE	End of calculation time
		F10.4	sec	DH	Time increment
		F10.4	--	DELTA	Convergence criterion for the calculation of the initial equilibrium state
7	80	F10.2	MW	PPW	Initial equilibrium power
		F10.2	°F	TGUES	Guessed initial moderator temperature to obtain equilibrium
		F10.2	°F	TGF	Guessed initial fuel temperature for obtaining equilibrium conditions.

- ①
- ②
- ③
- ④
- ⑤
- ⑥
- ⑦

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the impression as shown below. This card is necessary for the plot routine.

RD2 936L

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The output of FUMOTEM consists of three parts. The first is simply a writing of all the input data. The second is the equilibrium state calculation and the third block of data is the fuel temperature, moderator temperature reactivity, exit temperature, power, α_F and Q as a function of time. If the user specifies, a plot of power, precursor power, T_F , T_M and reactivity as a function of time is provided.

Problem 2.5.1

Run FUMOTEM for the sample data cards shown.

Problem 2.5.2

Run FUMOTEM for the reactivity input

$$\rho(t) = 0.50 \cos 2.5t$$

with the following parameters:

$$\text{NOPLT} = 1$$

$$\alpha_M = -0.00005$$

$$\beta = 0.00645$$

$$\lambda = 0.07695 \text{ sec}^{-1}$$

$$\gamma = 0.0002$$

$$\text{RTIME} = 4.0 \text{ sec}$$

$$P_0 = 0.80$$

$$R_F = 0.01504 \text{ ft}$$

$$\text{Number of fuel assemblies} = 745$$

$$P_C = 0.04733 \text{ ft}$$

$$C_{PF} = 0.0590 \text{ Btu/lb-}^\circ\text{F}$$

$$\text{Number of fuel pins per assembly} = 208$$

$$\rho_F = 43.2 \text{ lb/ft}^3$$

$$L = 12 \text{ ft}$$

$$T_{M1} = 400^\circ\text{F}$$

$$\text{TE} = 4.0 \text{ sec}$$

$$\text{DELTA} = 0.01$$

DH = 0.01 sec
 P₀ = 1000 MW
 TGUES = 200.0°F
 TGF = 500.0°F
 v = 13.0 ft/sec

Also run the same calculation for v = 26.0 ft/sec.

FUMOTEM is written in FORTRAN in single precision except for the eigenvalue calculation. The solution of the equation

$$|\underline{A} - \omega \underline{I}| = 0$$

for the root $\omega_0, \omega_1, \omega_2, \omega_3$ is done in double precision. All four roots are determined and the largest one is picked to form the \underline{H} and \underline{R} matrices.

The memory requirement is about 40 kilobytes and the execution time varies with the time length. The reactor is to be simulated. Generally it takes about 1 1/2 seconds of computer time to simulate one second of reactor transient time.

REFERENCES

1. D. L. Hetrick, "Dynamics of Nuclear Reactors," The University of Chicago Press, 1971, page 159.
2. J. R. Lamarsh, "Nuclear Reactor Theory," Addison-Wesley Publishing Company, Reading, Massachusetts, 1966, pages 218, 459.
3. M. M. El-Wakil, "Nuclear Heat Transport," International Textbook Company, Scranton, Pa., 1971, page 104.

List of Symbols for FUMOTEM

The following symbols are listed in alphabetical order in the FUMOTEM program.

A	a	Period of reactivity insertion
AF	α_F	Fuel temperature coefficient
AM	α_M	Moderator temperature coefficient
AMTRX		Subroutine to form the <u>A</u> matrix
AMX		Elements of the <u>A</u> matrix
AX		Coefficients of the EIGEN4 polynomial, i.e., if $\sum_{i=1}^5 a_n \omega^n$, the a_n .
BE		<u>B</u> matrix element.
BETA	β	Delayed neutron fraction
CK		Function subroutine to calculate the water thermal conductivity as a function of temperature (Equation 2.3.21)
COL	C	Colburn number (Equation 2.3.17)
CP	C_{PM}	Heat capacity of moderator as a function of T_M
CPF	NC_{PF}	Total heat capacity of the fuel
CPFX	C_{PF}	Heat capacity of a single fuel rod
CPG, CPMX	C_{PM}	Total heat capacity of the water
CX, CGUES	k_M	Thermal conductivity of the water
DE	De	Equivalent diameter (Equation 2.3.16)
DELTA	Δ	Convergence criterion for equilibrium calculation

DF	$2R_F$	Fuel diameter
DH	Δt	Time increment (TINCR)
DH1		Lower limit time increment
DH2		Upper limit time increment
DM	t	Time at any instant
DOTMG	\dot{m}_M	Mass flow of the coolant
DX, DGUES	ρ_M	Density of the coolant
EIGEN4		Subroutine to change $ \underline{A} - \omega \underline{I} = 0$ to polynomial form
EIGENV	ω	Eigenvalues of the \underline{A} matrix
FA		Heat transfer area
FCA		Fuel cross section area
FDENS	ρ_F	Fuel density
FH	L	Fuel rod length
FK		Function subroutine used to calculate the fuel conductivity as a function of temperature
FLA		Flow area
FX, FGUES	k_F	Thermal conductivity of fuel
G	γ	Constant in Equation (2.3.12)
GXN		Subroutine to multiply an $n \times n$ matrix with a column matrix
H	h_T	Heat transfer coefficient
HH		Elements of the \underline{H} matrix (Equation 2.4.19)
HMTRX	\underline{H}	Subroutine to form the \underline{H} matrix
HP		Element of the column matrix $\underline{H}\phi$
HX, HGUES	h_T	Heat transfer coefficient

MTRX		Order of the nxn matrix
MX	t	Dimensioned time
N		Number of iterations
NA		Number of fuel assemblies
NCHAN		Total number of coolant channels
NFRPA		Number of fuel rods per assembly
NIN		Number of iterations calculated from insertion time and initial time increment
NN		Total number of iterations, calculated from end time and initial time increment
NOPLT		Option for plotting
NR		READ statement unit number
NRO		Option for reactivity insertion
NW		WRITE statement unit number
NZ		Dummy variable used for print out
PAREA		Square pitch area for channel
PC	P_C	Distance between fuel pins
PI	π	3.14159
PKL		$4\pi k_{LF}$
POW	P_o	Input power in MW
PP	P	Power in Btu/hr
PR	Pr	Prandtl number function subroutine
PW	P_w	Wetted perimeter
QQ, QW	Q	Precursor density (Power)
RIN		Reactivity inserted
RB		Element of column matrix \underline{R} \underline{B}

RCOS	$\rho = \rho_0 \cos a t$	Function subroutine to calculate reactivity for a cosine insertion
RE	Re	Function subroutine to calculate the Reynolds number
RF	R_F	Fuel radius
RHO	$\rho_M(T_M)$	Function subroutine to calculate density of H_2O
RL	$\rho = \rho_0 (1 + a t)$	Function subroutine to calculate reactivity for a ramp
RMTRX		Subroutine to form the \underline{R} matrix (Equation 2.4.20)
RN		Reynolds number
RO	ρ_0	Reactivity inserted at $t = 0$
RSIN	$\rho = \rho_0 \sin a t$	Function subroutine to calculate reactivity for a sine insertion
RP	Pr	Prandtl number
RR	$\rho_T = \rho_{in} + \rho_f$	Total reactivity, inserted plus feedback
RTIME	t_r	Time when inserted reactivity is removed (for ramp reactivity only)
RX		Element of \underline{R} matrix (Equation 2.4.20)
RY		Dimensioned reactivity
TCFA		Total cross sectional area of the fuel
TCFLA		Total cross sectional area of coolant
TE		End of calculation time
TFA		Total heat transfer area
TFG		Initial guess for fuel temperature
TFO	T_{FO}	Equilibrium fuel temperature corresponding to P_0
TGUES, TMG		Initial guess for moderator temperature

TMI

 T_{MI}

Inlet coolant temperature

TMO

 T_{MO} Equilibrium coolant temperature
corresponding to a given P_0

TMOUT

Exit temperature of core water

TPLOT

Subroutine to plot five variables at
the same time with respect to the
independent variable time

U

 μ Function subroutine to calculate the
viscosity of water as a function of
temperature

UGUES, UX

Viscosity of ~~water~~

V

 V

Mean velocity of coolant

VF

 V_F

Fuel volume

VM

 V_M

Coolant volume

WO

 ω_0 The largest eigenvalue to the
equation $|\underline{A} - \omega \underline{I}| = 0$

X

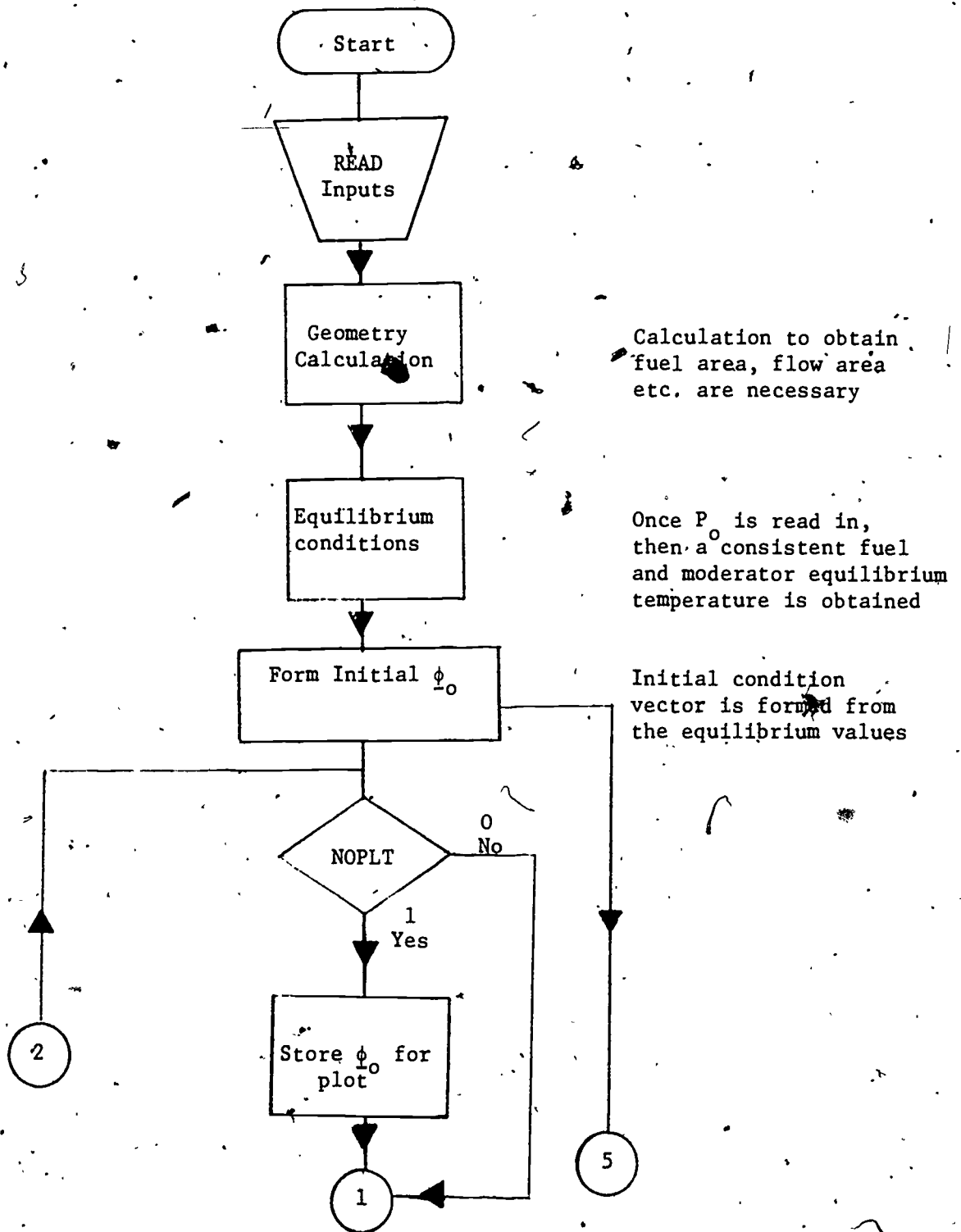
 λ Decay constant of the delayed neutron
group

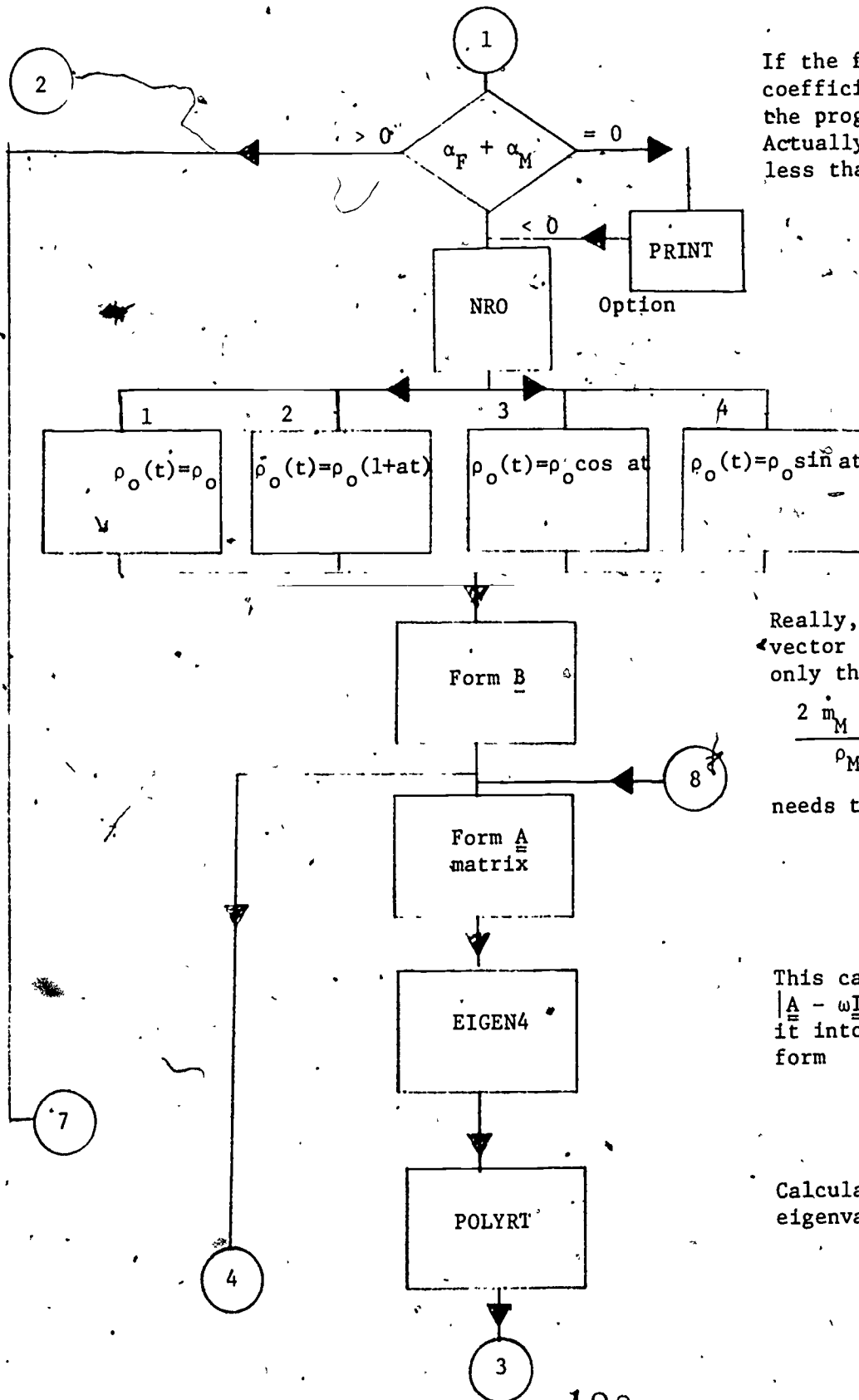
XL

 Λ

The neutron generation time

Flow Chart for FUMOTEM





If the feedback reactivity coefficients are positive, the program stops. Actually, $\alpha_F + \alpha_M$ must be less than or equal to zero

One of these four reactivity insertions is to be used

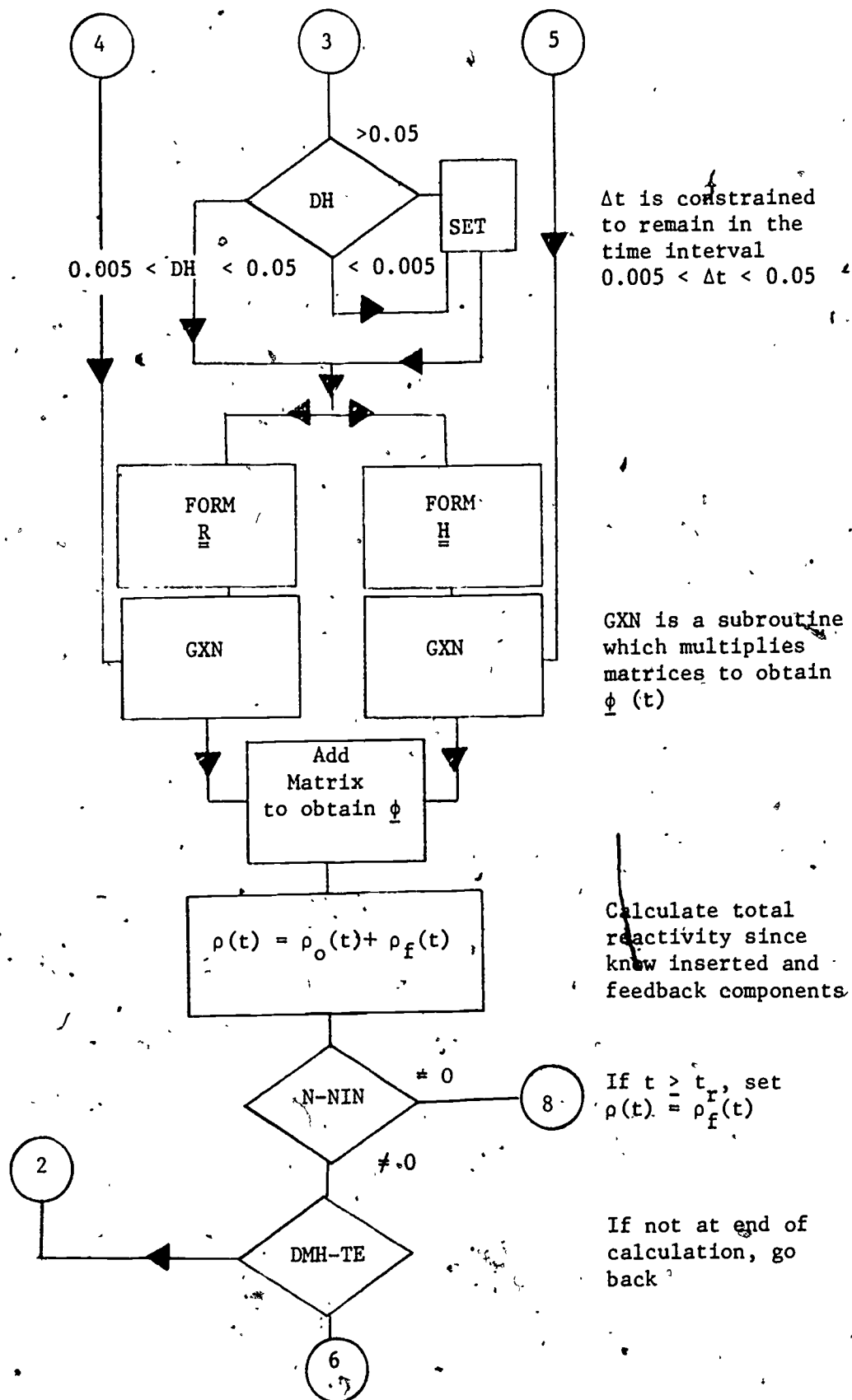
Really, since the \underline{B} vector has 3 zeroes, only the number

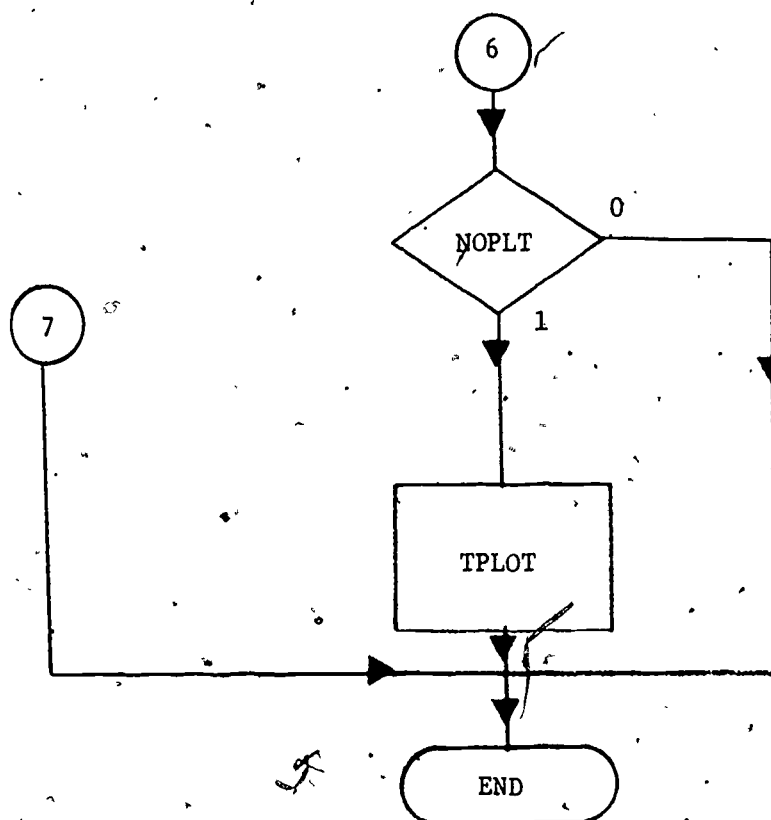
$$\frac{2 \dot{m}_M T_{M1}(t)}{\rho_M V_M}$$

needs to be solved.

This calculates the $|\underline{A} - \omega \underline{I}| = 0$ and puts it into polynomial form

Calculates the largest eigenvalue ω_o





This subroutine
plots the state
vector ϕ as a
function of time.

```
//WATFIV      ONEGA,TIME=300,PAGES=50
----- PWR FEEDBACK -----
```

CODE NAME F U M O T E M .

OBJECTIVE

- 1.EXAMINE THE TEMPERATURE FEEDBACK MECHANISM OF A PWR AND
- 2.SOLVE THE ONE DELAYED NEUTRON MODEL WITH TEMPERATURE FEEDBACK FOR A STEP INSERTION AND A RAMP INSERTION OF REACTIVITY .

PROGRAM

WRITTEN IN SINGLE PRECISION .

DESCRIPTIONS OF INPUT PARAMETERS .

FORMAT~ NUMBER	PARAMETER	FUNCTION	UNIT
20	NOPLT	OPTION FOR PLOTTING 1 - PLOT THE RESULT . 0 - NO PLOT .	
30	BETA	AVERAGE FRACTION OF DELAYED NEUTRONS , IS 0.00645 , ONLY ONE DELAYED NEUTRON GROUP IS CONSIDERED .	
	X	DELAYED NEUTRON DECAY CONSTANT .	1/SEC .
	XL	NEUTRON GENERATION TIME .	SEC
	RO	INITIAL REACTIVITY INSERTED . INPUT RO IN DOLLAR UNIT .	
	NRO	OPTION FOR REACTIVITY INSERTION AS A FUNCTION OF TIME . 1 - CONSTANT REACTIVITY . 2 - LINEAR RAMP INSERTION . 3 - COSINE VARIATION OF REACTIVITY . 4 - SINE VARIATION OF REACTIVITY WITH TIME	
	RTIME	INSERTION TIME	SEC
	A	PERIOD , CONSTANT .	1/SEC
40	AM	MODERATOR TEMPERATURE COEFFICIENT	1/F .
	G	CONSTANT GAMMA IN EQUATION 2.3.13 OF RD-2	
	PTFO	RESONANCE ESCAPE PROBABILITY .	SQRT(F)
FUEL ROD PROPERTIES .			
50	RF	RADIUS .	FT .
	PC	DISTANCE BETWEEN RODS .	FT .
	CPFX	SPECIFIC HEAT .	BTU/LB-F
	FDENS	DENSITY .	LB/FT**3
	FH	LENGTH .	FT .
	NA	NUMBER OF FUEL ASSEMBLIES .	
	NFRPA	FUEL RODS PER ASSEMBLY .	
60	V	VELOCITY OF MODERATOR/COOLANT	FT/SEC .
	TM1	INLET TEMPERATURE OF MODERATOR .	F .

C						A	60
C	70	TE	END TIME OF CALCULATION	SEC.		A	61
C		DH	TIME INCREMENT, THE RANGE WILL			A	62
C			BE BETWEEN 0.005 TO 0.05 SECOND.			A	63
C		DELTA	CONVERGING FACTOR.			A	64
C						A	65
C	80	PPW	GIVEN EQUILIBRIUM POWER AT T=0.0	MW		A	66
C		TGUES	COOLANT/MODERATOR GUESSED TEMPERATURE.	F.		A	67
C			CORRESPONDING TO PP.			A	68
C		TFG	FUEL GUESSED TEMPERATURE CORRESPONDING TO PP.	F.		A	69
C						A	70
C						A	71
C						A	72
C			MAIN PARAMETERS OTHER THAN INPUT			A	73
C						A	74
C		POWER DENSITY.		PHI(1),MPE,MPEL		A	75
C		PRECURSOR DENSITY.		PHI(2),MQU		A	76
C		FUEL TEMPERATURE.		PHI(3),MTEEF		A	77
C		MODERATOR/COOLANT TEMPERATURE.		PHI(4),MTEEM		A	78
C		REACTIVITY AT ANY INSTANT.		RR,RY		A	79
C		TIME ELAPSED.		DMH,MX		A	80
C		TOTAL NUMBER OF ITERATIONS.		NN		A	81
C		COOLANT/MODERATOR EQUILIBRIUM				A	82
C		TEMPERATURE.		TMO		A	83
C		FUEL EQUILIBRIUM TEMPERATURE.		TFO		A	84
C						A	85
C		HEAT TRANSFER COEFFICIENT.		HX,HGUES		A	86
C		TOTAL MASS FLOW		DOTMG		A	87
C		DUMMY INDICATOR FOR PRINT-OUT THE				A	88
C		EQUILIBRIUM CONDITION BEFORE AND				A	89
C		AFTER REACTIVITY INSERTION.		NZ		A	90
C		FUEL TEMPERATURE COEFFICIENT.		AAF		A	91
C		NUMBER OF ITERATION AT ANY TIME.		N		A	92
C						A	93
C						A	94
C						A	95
C						A	96
C						A	97
1		REAL MPEL(500),MPE(500),MQU(500),MTEEF(500),MTEEM(500),MX(500),RY(A	98
2		1500)				A	99
3		DIMENSION AMX(4,4),HH(4,4),RX(4,4)				A	100
4		DIMENSION ROOTI(4),EIGNV(4),COF(5),AX(5)				A	101
5		DIMENSION PHI(4),BE(4),RB(4),HP(4)				A	102
6		COMMON BETA,XL,X,FX,RF,OH,FDENS,CPF,VF,FH,VM,DX,DOTMG,W0,HX,RR,C				A	103
		1PMX,PI,NROD				A	104
		DOUBLE PRECISION AX,COF,EIGNV,ROOTI				A	105
						A	106
		-----DEFINE REYNOLDS NUMBER, PRANDTL NUMBER, HEAT TRANSFER CO				A	107
		EFFICIENTS, FUEL COEFFICIENT TEMPERATURE AND REACTIVITY.				A	108
						A	109
7		RE(RE1,RE2,RE3,RE4)=RE1*RE2*RE3/RE4				A	110
8		PR(PR1,PR2,PR3)=PR1*PR2/PR3				A	111
9		H(H1,H2,H3,H4,H5)=H1*H2*(H3**8)*(H4**3333)/H5				A	112
10		AF(AF1,AF2,AF3)=-AF1*ALOG(1.0/AF3)/SQRT(AF2)				A	113
11		RT(RT1,RT2,RT3,RT4,RT5,RT6,RT7)=RT1+(RT2+RT3)*RT4-RT5*RT2*RT6+RT7				A	114
						A	115
		-----FUEL AND WATER PROPERTY AS TEMPERATURE DEPENDENT.				A	116
						A	117
12		FK(T)=0.61021E1-0.46365E-2*T+0.13063E-5*T**2				A	118
13		CK(T)=0.11711+0.13910E-2*T-0.18102E-5*T**2				A	119

14	RHD(T)=0.57788E2+0.28018E-1*T-0.88346E-4*T**2	A 120
15	UT(T)=0.85067-0.17501E-2*T+0.11420E-5*T**2	A 121
16	CP(T)=0.47088E1-0.15753E-1*T+0.17233E-4*T**2	A 122
	C	A 123
	C	A 124
	C	A 125
	C	A 126
	C	A 127
17	NN=1600	A 128
18	NR=5	A 129
19	NW=6	A 130
20	NZ=0	A 131
21	PI=3.14159	A 132
22	N=1	A 133
23	DM=0.0	A 134
24	OH1=0.005	A 135
25	OH2=0.05	A 136
26	OMH=DM*3600.0	A 137
27	R=0.0	A 138
28	RR=0.0	A 139
29	MX(1)=0.0	A 140
	C	A 141
	C	A 142
	C	A 143
30	READ (NR,29) NOPLT	A 144
31	READ (NR,30) BETA,X,XL,RO,NRO,RTIME,A	A 145
32	READ (NR,31) AM,G,PTFO	A 146
33	READ (NR,32) RF,PC,CPFX,FDENS,FH,NA,NFRPA	A 147
34	READ (NR,33) V,TMI	A 148
	C	A 149
	C	A 150
	C	A 151
	C	A 152
	C	A 153
	C	A 154
35	READ (NR,34) TE,OH,DELTA	A 155
	C	A 156
	C	A 157
	C	A 158
	C	A 159
36	READ (NR,35) PPW,TGUES,TFG	A 160
37	PP=3.41206*PPW	A 161
	C	A 162
	C	A 163
	C	A 164
38	WRITE (NW,36)	A 165
39	WRITE (NW,37)	A 166
40	WRITE (NW,38) BETA,X,XL,RO	A 167
41	RO=RO*BETA	A 168
42	WRITE (NW,39) RTIME,A	A 169
43	WRITE (NW,40) AM,G,PTFO	A 170
44	WRITE (NW,41) RF,PC,CPFX,FDENS,FH	A 171
45	WRITE (NW,42) NA,NFRPA	A 172
46	WRITE (NW,43) V,TMI	A 173
47	WRITE (NW,44) NOPLT,NRO	A 174
48	WRITE (NW,45) PPW,TFG,TGUES	A 175
49	WRITE (NW,46) TE,OH	A 176
50	WRITE (NW,47) DELTA	A 177
	G	A 178
	C	A 179
	-----MAKE SURE THAT NN EQUAL TO DIMENSION NUMBER OF THE PLOTTED VARIABLE .	
	-----READ INPUT DATA .	
	-----THE INITIAL VALUE OF OH WILL CHANGE ACCORDING TO THE LAR - GEST EIGEN VALUE OF A-MATRIX .	
	-----READ INITIAL POWER DESIRED, AND GUESSED FUEL AND COOLANT TEMPERATURE .	
	-----PRINT OUT INPUT DATA AND THE INITIAL .	
	-----CHANGE SECOND TO UNIT HOUR .	

51	C	NOH=IFIXITE/DH+DH/2.0)	A 180
52		IF (NDH.LE.NN) GO TO 1	A 181
53		WRITE (NM,48)	A 182
54		GO TO 28	A 183
55	J	XL=XL/3600.0	A 184
56		X=X*3600.	A 185
57		V=V*3600.	A 186
58		DH1=DH1/3600.0	A 187
59		DH2=DH2/3600.0	A 188
60		DH=DH/3600.	A 189
61		A=A*3600.0	A 190
	C		A 191
	C	-----GEOMETRY CALCULATIONS FOR SQUARE PITCH.	A 192
	C		A 193
62		DF=2.0*RF	A 194
63		FA=PI*DF*FH	A 195
64		FCA=PI*DF**2/4.0	A 196
65		PAREA=PC*PC	A 197
66		FLA=PAREA-FCA	A 198
67		NROD=NA*NFRPA	A 199
68		NCHAN=NROD	A 200
69		TFA=NROD*FA	A 201
70		TCFA=NROD*FCA	A 202
71		TCFLA=NCHAN*FLA	A 203
72		VF=FCA*FH*NROD	A 204
73		VM=TCFLA*FH	A 205
74		PM=PI*DF	A 206
75		DE=4.0*FLA/PM	A 207
76		CPF=CPF*FDENS*VF	A 208
77		COL=0.042*PC/DF-0.024	A 209
	C		A 210
	C	-----ITERATE FUEL AND COOLANT TEMPERATURE UNTIL IT CONVERGES TO	A 211
	C	THE CORRESPONDING POWER AND ITS PRECURSOR.	A 212
	C		A 213
78	2	CONTINUE	A 214
79		UGUES=U(TGUES)	A 215
80		DGUES=RHO(TGUES)	A 216
81		DOTMG=DGUES*TCFLA*V	A 217
82		CGUES=CK(TGUES)	A 218
83		FGUES=FK(TFG)	A 219
84		PKL=4.0*PI*FGUES*FH	A 220
85		CPG=CP(TGUES)	A 221
86		TMG=TMI+PP/(2.0*CPG*DOTMG)	A 222
87		TFG=TMG+PP/(2.0*PKL*NROD)	A 223
88		RN=RE(DGUES,V,DE,UGUES)	A 224
89		RP=PR(CPG,DGUES,CGUES)	A 225
90		HGUES=H(COL,CGUES,RN,RP,DE)	A 226
91		PGUES=2.0*CPG*DOTMG*ABS(TGUES-TMI)	A 227
92		OP=PGUES-PP	A 228
93		DTM=UGUES-TMG	A 229
94		DELPG=ABS(OP)/PP	A 230
	C		A 231
	C	-----SET THE GUESSED POWER AND ITERATE UNTIL CONVERGE TO THE	A 232
	C	CORRESPONDING FUEL AND MODERATOR TEMPERATURE ACCURATE TO THE	A 233
	C	VALUE OF DELTA.	A 234
	C		A 235
95		IF (DELPG-DELTA) 6,6,3	A 236
96	3	IF (DTM) 4,6,5	A 237
97	4	TGUES=TGUES+ABS(DTM)/2.0	A 238
			A 239

98		GO TO 2	A 240
99	5	TGUES=TGUES-ABS(OTM)/2.0	A 241
100		GO TO 2	A 242
101	6	CONTINUE	A 243
102		QQ=BETA*PP/(X*XL)	A 244
103		QW=QQ/3.412E6	A 245
	C		A 246
	C	-----INITIALIZE MATRIX PHI .	A 247
	C		A 248
104		PHI(1)=PP	A 249
105		PHI(2)=QQ	A 250
106		PHI(3)=TFG	A 251
107		PHI(4)=TMG	A 252
108		DOTMG=DGUES*TCFLA*V	A 253
	C		A 254
	C	-----SETUP THE FUEL AND COOLANT TEMPERATURE AT EQUILIBRIUM	A 255
	C	STATE .	A 256
	C		A 257
109		TMO=TMG	A 258
110		TFO=TFG	A 259
111		RS=RR/BETA	A 260
112		WRITE (6,49)	A 261
113		WRITE (NW,50) PPW,QW,TFO,TMO,RS	A 262
114		WRITE (NW,51) RN,RP,COL,DOTMG,HGUES	A 263
115		WRITE (NW,52)	A 264
116		WRITE (NW,53)	A 265
117	7	CONTINUE	A 266
118		MTRX=4	A 267
	C		A 268
	C	-----STORE THE VALUE OF POWER , PRECURSOR DENSITY , FUEL AND MO	A 269
	C	DERATOR TEMPERATURE , AND REACTIVITY FOR PLOTTING .	A 270
	C		A 271
119		PPW=PHI(1)/3.412E6	A 272
120		QW=PHI(2)/3.412E6	A 273
121		TFG=PHI(3)	A 274
122		TMG=PHI(4)	A 275
123		IF (NOPLT.EQ.0) GO TO 8	A 276
124		MPE(N)=PPW	A 277
125		MPEL(N)=ALOG10(PPW)	A 278
126		MQU(N)=QW	A 279
127		MTEEF(N)=TFG	A 280
128		MTEEM(N)=TMG	A 281
129	8	AAF=AFIG,PHI(3),PTFO)	A 282
130		IF (AAF.LT.0.0.AND.AM.LT.0.0) GO TO 10	A 283
131		IF (AAF+AM) 10,9,28	A 284
132	9	WRITE (NW,54)	A 285
133	10	UX=U(PHI(4))	A 286
134		DX=RHO(PHI(4))	A 287
135		CX=CK(PHI(4))	A 288
136		CPMX=CP(PHI(4))	A 289
137		FX=FK(PHI(3))	A 290
138		DOTMG=DX*TCFLA*V	A 291
139		HX=H(COL,CX,RE(DX,V,DE,UX),PR(CPMX,UX,CX),DE)	A 292
140		TMOU=2.0*PHI(4)-TMI	A 293
	C		A 294
	C	-----COMPUTE THE REACTIVITY INSERTED AS A FUNCTION OF TIME	A 295
	C	AND REACTIVITY .	A 296
	C		A 297
141		GO TO (11,12,13,14), NRO	A 298
142	11	RIN=RO	A 299

143		GO TO 15	A 300
144	12	RIN=RO*(1+A*DM)	A 301
145		GO TO 15	A 302
146	13	RIN=RO*COS(A*DM)	A 303
147		GO TO 15	A 304
148	14	RIN=RO*SIN(A*DM)	A 305
149	15	CONTINUE	A 306
150		IF (DMH-RTIME) 19,19,16	A 307
151	16	IF (NRO.GT.2) GO TO 19	A 308
152		IF (NRO-1) 18,17,18.	A 309
153	17	RIN=0.0	A 310
154		GO TO 19	A 311
155	18	RTMH=RTIME/3600.	A 312
156		RIN=RO*(1+A*RTMH)	A 313
157		GO TO 19	A 314
158	19	RY(N)=RR/BETA	A 315
	C		A 316
	C	-----FORM B-MATRIX .	A 317
	C		A 318
159		BE(1)=0.0	A 319
160		BE(2)=0.0	A 320
161		BE(3)=0.0	A 321
162		BE(4)=+2.0*DDTMG*TM1/(DX*VM)	A 322
	C		A 323
	C	-----FORM A-MATRIX .	A 324
	C		A 325
163		CALL AMTRX (AMX)	A 326
	C		A 327
	C	-----CHANGE A-DETERMINANT TO POLYNOMIAL FORM .	A 328
	C		A 329
164		CALL EIGEN4 (AMX,AX)	A 330
165		M=5	A 331
166		M1=M	A 332
167		IF (AX(1)) 22,20,22	A 333
168	20	DO 21 JK=1,4	A 334
169	21	AX(JK)=AX(JK+1)	A 335
170		MTRX=3	A 336
171		M1=4	A 337
	C		A 338
	C	-----COMPUTE THE EIGENVALUE OF THE A-MATRIX .	A 339
	C		A 340
172	22	CALL POLRT (AX,COF,MTRX,EIGNV,ROOTI,IER,M1)	A 341
173		MTRX=4	A 342
174		ML=M1-1	A 343
	C		A 344
	C	-----FIND THE LARGEST EIGENVALUE .	A 345
	C		A 346
175		WO=EIGNV(1)	A 347
176		DO 23 IE=1,ML	A 348
177		WE=EIGNV(IE)	A 349
178	23	WO=AMAX1(WE,WO)	A 350
	C		A 351
	C	-----CALCULATE THE TIME INCREMENT , SUCH THAT THE TIME INCRE	A 352
	C	MENT RECIPROCAL TO THE LARGEST EIGEN-VALUE .	A 353
	C	MAKE SURE THAT THE RANGE IS BETWEEN 0.005 TO 0.05 SECOND , IF	A 354
	C	SMALLER THAN 0.005 SECOND CHANGE THE VALUE TO 0.005 AND IF LAR	A 355
	C	GER THAN 0.05 SET IT TO 0.05 SECOND .	A 356
	C		A 357
179		DH=ABS(1.0/WO)/10.0	A 358
180		DOH=DH	A 359

```

181      IF (DDH.LT.DH1) DH=DH1
182      IF (DDH.GT.DH2) DH=DH2
      C
      C      -----FORM R-MATRIX .
      C
183      CALL RMTRX (RX)
      C
      C      -----FORM H-MATRIX .
      C
184      CALL HMTRX (HH)
      C
      C      -----MULTIPLY R-MATRIX WITH COLUMN MATRIX-B AND MULTIPLY
      C      H-MATRIX WITH COLUMN VECTOR PHI .
      C
185      CALL GXN (RX,BE,MTRX,RB)
186      CALL GXN (HH,PHI,MTRX,HP)
      C
      C      -----FORM NEW PHI-MATRIX BY ADDING THE TWO COLUMN MATRIX .
      C
187      DO 24 IP=1,MTRX
188      24  PHI(IP)=RB(IP)+HP(IP)
189      WRITE (NW,55) N,DMH,RY(N),PPW,QW,TFG,TMG,TMOUT,AAF
190      IF (DMH-TE) 25,27,27
      C
      C      -----IF MODERATOR TEMPERATURE EXCEEDING 700 DEGREE F ,PRINT A
      C      WARNING AND GET CUT .
      C
191      25  IF (PHI(4).GE.1000.0) GO TO 26
192      RR=AM*(PHI(4)-TMO)+AAF*(PHI(3)-TFO)+RIN
193      R=RR
194      DM=DM+DH
195      DMH=DM*3600.0
196      IF (DMH.GE.TE) GO TO 27
197      N=N+1
198      MX(N)=DMH
199      GO TO 7
200      26  WRITE (NW,56)
201      27  NN=N
202      IF (NOPLT.EQ.0) GO TO 28
      C
      C      -----PLOT THE RESULT SIMULTANEDUSLY IN ONE GRAPH .
203      CALL TPLOT (MX,MPE,MPEL,MQU,MTEEM,MTEEF,RY,NN)
204      STOP
      C
205      29  FORMAT (I2)
206      30  FORMAT (2F10.5,2F10.7,I10,2F10.5)
207      31  FORMAT (3F10.6)
208      32  FORMAT (3F10.5,2F10.2,2I10)
209      33  FORMAT (2F10.2)
210      34  FORMAT (3F10.4)
211      35  FORMAT (3F10.2)
212      36  FORMAT (1H)
213      37  FORMAT (1X,23H*****//,1X,19HMODULE 2 , FEEDBACK,
1/,1X,23H*****//,1X,10HINPUT DATA,/,1X,14H***
2*****//)
214      38  FORMAT (5X,4HBETA,5X,6HLAMBDA,5X,17HNEUTRON GEN. TIME,5X,18HINITIA
1L REACTIVITY,/,12X,11H( SEC*-1 ),6X,7H( SEC ),18X,5H( $ ),/,3X,F8
2.6,3X,F7.5,7X,E9.3,17X,F5.2,/)
215      39  FORMAT (5X,14HINSERTION TIME,5X,15HCONSTANT PERIOD,/,9X,5H( SEC ),14
1X,7H(1/SEC),/,8X,F7.3,13X,F7.3,/)

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216	40	FORMAT (5X,14HCOOLANT COEFF.,5X,10HCONSTANT 8,5X,20HRESONANCE ESC.	A 420
		1 PROB.,/,6X,12H() ,7X,8H() ,7X,18H()	A 421
		2,/,6X,E12.5,7X,E8.2,13X,F5.3,/))	A 422
217	41	FORMAT (5X,11HFUEL RADIUS,5X,5HPITCH,5X,8HFUEL CP.,5X,12HFUEL DENS	A 423
		ITY,5X,11HFUEL HEIGHT,5X,18H) ,5X,12H)	A 424
		2/,7X,6H(FT) ,8X,6H(FT) ,2X,12H(BTU/LB-F) ,3X,12H(LB/FT**3) ,7X	A 425
		3,6H(FT) ,9X,16H) ,7X,10H) ,/,7X,F7.5,7X,F7.	A 426
		45,4X,F6.4,8X,F6.2,10X,F6.2,/))	A 427
218	42	FORMAT (5X,18HNUMBER OF ASSEMBLY,5X,26HNUMBER OF ROD PER ASSEMBLY,	A 428
		1/,5X,18H() ,5X,26H() ,/,12X,	A 429
		2,4,23X,14,/))	A 430
219	43	FORMAT (5X,16HCOOLANT VELOCITY,5X,20HINLET COOLANT TEMPT.,/,8X,10H	A 431
		1 FT/SEC) ,15X,5H(F) ,/,11X,F4.1,18X,F5.1,/))	A 432
220	44	FORMAT (9X,6HOPTION,4X,4HPLOT,9X,11,/,19X,11HTYPE INSR.,2X,11,/))	A 433
221	45	FORMAT (5X,13HINITIAL POWER,9X,4H(MW),1X,E11.4,/,5X,19HGUESSED FUE	A 434
		1L TEMPT.,4X,3H(F),5X,F6.1,/,5X,22HGUESSED COOLANT TEMPT.,1X,3H(F),	A 435
		26X,F5.1,/,2X,17HEND OF INPUT DATA,/,2X,18(1H*),/))	A 436
222	46	FORMAT (/,5X,12H END TIME ,3X,14HTIME INCREMENT,/,8X,F6.3,10X,F	A 437
		16,4,/))	A 438
223	47	FORMAT (5X,22H CONVERGENCE FACTOR = ,F9.5,/))	A 439
224	48	FORMAT (/,5X,118H*** CHECK DIMENSION , IFIX(TE/DH) HAS TO BE SMA	A 440
		1LLER OR EQUAL TO THE DIMENSION OF MPEL,MPE,MTEEF,MTEEM,MY,RY AND M	A 441
		2QU .,/))	A 442
225	49	FORMAT (/,5X,27HEQUILIBRIUM STATE INITIALLY,/,4X,28(1H*),/))	A 443
226	50	FORMAT (5X,5HPPOWER,5X,9HPRECURSOR,5X,11HFUEL TEMPT.,5X,14HCOOLANT	A 444
		1TEMPT.,5X,10HREACTIVITY,5X,11H) ,/,3X,E10.4,2X,E10.4,6X,F	A 445
		27.2,11X,F6.2,11X,F6.3,/))	A 446
227	51	FORMAT (85(1H*),/),5X,10HREYNOLDS #,3X,9HPRANDTL #,3X,9HCOLBURN	A 447
		1#,3X,9HMASS FLOW,3X,19HHEAT TRANSFER COEFF.,/,42X,9H(LB/HR) ,4X,18	A 448
		2H(BTU/HR-F-FT**2) ,/,6X,F8.1,6X,F4.1,5X,F7.3,5X,E11.4,8X,F7.1,/))	A 449
		3)	A 450
228	52	FORMAT (5X,2HND,10X,4HTIME,6X,10HREACTIVITY,7X,5HPPOWER,11X,9HPRECU	A 451
		1RSOR,/,4X,4H() ,8X,5H(SEC),9X,3H(\$),11X,4H(MW),14X,4H(MW))	A 452
229	53	FORMAT (7X,9HFUEL TEMP,5X,9HMOD. TEMP,3X,9HEXIT TEMP,8X,11HFUEL CO	A 453
		IEFF.,/,10X,3H(F),11X,3H(F),9X,3H(F),14X,5H(1/F),/))	A 454
230	54	FORMAT (/,8X,19H... NO FEEDBACK .../))	A 455
231	55	FORMAT (3X,14,8X,F7.3,7X,F5.2,7X,E11.4,7X,E11.4,/,8X,F7.2,7X,F7.2,	A 456
		15X,F7.2,9X,E11.4)	A 457
232	56	FORMAT (/,5X,39H C R I T I C A L T E M P E R A T U R E ,/))	A 458
233		END	A 459
234		SUBROUTINE RMTRX (R)	B 1
	C		B 2
	C		B 3
	C	-----FORM R-MATRIX , DIAGONAL MATRIX IN EQUATION 2.4.20	B 4
	C	SUPPORTING ROUTINE NONE	B 5
	C		B 6
	C		B 7
235		DIMENSION R(4,4),A(4,4)	B 8
236		COMMON BETA,XL,X,FX,RF,DH,FDENS,CPF,VF,FH,VM,DX,DOTMG,WJ,HX,RR,C	B 9
		1PMX,PI,NROD	B 10
237		CALL AMTRX (A)	B 11
238		DO 3 I=1,4	B 12
239		DO 3 II=1,4	B 13
240		IF (I-II) 1,2,1	B 14
241	1	R(I,II)=0.0	B 15
242		GO TO 3	B 16
243	2	R(I,II)=(EXP(A(I,II)*DH)-1.0)/A(I,II)	B 17
244	3	CONTINUE	B 18

245	RETURN	B	19
246	END	B	20
247	SUBROUTINE HMTRX (H)	C	1
	-----FORM H-MATRIX AS IN THE EQUATION 2.4.19 .	C	2
	SUPPORTING ROUTINE NONE	C	3
		C	4
		C	5
		C	6
		C	7
248	DIMENSION H(4,4),A(4,4)	C	8
249	COMMON BETA,XL,X,FX,RF,DH,FDENS,CPF,VF,FH,VM,DX,DOTMG,W0,HX,RR,C	C	9
	1PMX,PI,NROD	C	10
250	CALL AMTRX (A)	C	11
251	H(1,1)=EXP(A(1,1)*DH)	C	12
252	H(1,2)=A(1,2)*(EXP(W0*DH)-EXP(A(1,1)*DH))/(W0-A(1,1))	C	13
253	H(1,3)=0.0	C	14
254	H(1,4)=0.0	C	15
255	H(2,1)=A(2,1)*(EXP(W0*DH)-EXP(A(2,2)*DH))/(W0-A(2,2))	C	16
256	H(2,2)=EXP(A(2,2)*DH)	C	17
257	H(2,3)=0.0	C	18
258	H(2,4)=0.0	C	19
259	H(3,1)=A(3,1)*(EXP(W0*DH)-EXP(A(3,3)*DH))/(W0-A(3,3))	C	20
260	H(3,2)=0.0	C	21
261	H(3,3)=EXP(A(3,3)*DH)	C	22
262	H(3,4)=A(3,4)*(EXP(W0*DH)-EXP(A(3,3)*DH))/(W0-A(3,3))	C	23
263	H(4,1)=A(4,1)*(EXP(W0*DH)-EXP(A(4,4)*DH))/(W0-A(4,4))	C	24
264	H(4,2)=0.0	C	25
265	H(4,3)=0.0	C	26
266	H(4,4)=EXP(A(4,4)*DH)	C	27
267	RETURN	C	28
268	END		
269	SUBROUTINE POLRT (XCOF,COF,M,ROOTR,ROOTI,IER,M1)	D	1
		D	2
		D	3
	SUBROUTINE POLRT	D	4
		D	5
	PURPOSE	D	6
	COMPUTES THE REAL AND COMPLEX ROOTS OF A REAL POLYNOMIAL	D	7
		D	8
	USAGE	D	9
	CALL POLRT(XCOF,COF,M,ROOTR,ROOTI,IER,M1)	D	10
		D	11
	DESCRIPTION OF PARAMETERS	D	12
	XCOF -VECTOR OF M+1 COEFFICIENTS OF THE POLYNOMIAL	D	13
	ORDERED FROM SMALLEST TO LARGEST POWER	D	14
	COF -WORKING VECTOR OF LENGTH M+1	D	15
	M -ORDER OF POLYNOMIAL	D	16
	ROOTR-RESULTANT VECTOR OF LENGTH M CONTAINING REAL ROOTS	D	17
	OF THE POLYNOMIAL	D	18
	ROOTI-RESULTANT VECTOR OF LENGTH M CONTAINING THE	D	19
	CORRESPONDING IMAGINARY ROOTS OF THE POLYNOMIAL	D	20
	IER -ERROR CODE WHERE	D	21
	IER=0 NO ERROR	D	22
	IER=1 M LESS THAN ONE	D	23
	IER=2 M GREATER THAN 36	D	24
	IER=3 UNABLE TO DETERMINE ROOT WITH 500 ITERATIONS	D	25
	ON 5 STARTING VALUES	D	26

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C IER=4 HIGH ORDER COEFFICIENT IS ZERO. D 27
C M1 -NUMBER OF COEFFICIENT , M+1 D 28
C (ADDED ARGUMENT FROM THE ORIGINAL TO GET MORE FLEXIBLE D 29
C DIMENSION ) D 30
C D 31
C REMARKS D 32
C LIMITED TO 36TH ORDER POLYNOMIAL OR LESS. D 33
C FLOATING POINT OVERFLOW MAY OCCUR FOR HIGH ORDER D 34
C POLYNOMIALS BUT WILL NOT AFFECT THE ACCURACY OF THE RESULTS. D 35
C SUBROUTINES AND FUNCTION SUBPROGRAMS REQUIRED D 36
C NONE Q 37
C METHOD D 38
C NEWTON-RAPHSON ITERATIVE TECHNIQUE. THE FINAL ITERATIONS D 39
C ON EACH ROOT ARE PERFORMED USING THE ORIGINAL POLYNOMIAL D 40
C RATHER THAN THE REDUCED POLYNOMIAL TO AVOID ACCUMULATED D 41
C ERRORS IN THE REDUCED POLYNOMIAL. D 42
C D 43
C D 44
C D 45
C D 46
270 DIMENSION XCOF(M1),COF(M1),ROOTR(M),ROOTI(M) D 47
271 DOUBLE PRECISION XC,YO,X,Y,XPR,YPR,UX,UY,V,YT,XT,U,XT2,YT2,SUMSQ, D 48
1 DX,OY,TEMP,ALPHA,DABS D 49
C D 50
C -----IF A DOUBLE PRECISION VERSION OF THIS ROUTINE IS DESIRED, D 51
C C IN COLUMN 1 SHOULD BE REMOVED FROM THE DOUBLE PRECISION D 52
C STATEMENT WHICH FOLLOWS. D 53
C D 54
272 DOUBLE PRECISION XCCF,CDF,ROOTR,ROOTI D 55
C D 56
C -----THE C MUST ALSO BE REMOVED FROM DOUBLE PRECISION STATEMENT D 57
C APPEARING IN OTHER ROUTINES USED IN CONJUNCTION WITH THIS D 58
C ROUTINE. D 59
C THE DOUBLE PRECISION VERSION MAY BE MODIFIED BY CHANGING THE D 60
C CONSTANT IN STATEMENT 78 TO 1.00-12 AND IN STATEMENT 122 TO D 61
C 1.00-10. THIS WILL PROVIDE HIGHER PRECISION RESULTS AT THE D 62
C COST OF EXECUTION TIME D 63
C D 64
C D 65
C D 66
273 IFIT=0 D 67
274 N=M D 68
275 IER=0 D 69
276 IF (XCOF(N+1)) 1,4,1 D 70
277 IF (N) 2,2,6 D 71
C D 72
C -----SET ERROR CODE TO 1 D 73
C D 74
278 IER=1 D 75
279 RETURN D 76
C D 77
C -----SET ERROR CODE TO 4. D 78
C D 79
280 IER=4 D 80
281 GO TO 3 D 81
C D 82
C -----SET ERROR CODE TO 2 D 83
C D 84
282 IER=2 D 85
283 GO TO 3 D 86

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284	6	IF (N-36) 7,7,5,?	D	87
285	7	NX=N	D	88
286		NXX=N+1	D	89
287		N2=1	D	90
288		KJ1=N+1	D	91
289		DO 8 L=1,KJ1	D	92
290		MT=KJ1-L+1	D	93
291	8	COF(MT)=XCOF(L)	D	94
	C		D	95
	C	-----SET INITIAL VALUES	D	96
	C		D	97
292	9	XO=.00500101	D	98
293		YO=0.01000101	D	99
	C		D	100
	C	-----ZERO INITIAL VALUE COUNTER	D	101
	C		D	102
294		IN=0	D	103
295	10	X=XO	D	104
	C		D	105
	C	-----INCREMENT INITIAL VALUES AND COUNTER	D	106
	C		D	107
296		XO=-10.0*YO	D	108
297		YO=-10.0*X	D	109
	C		D	110
	C	-----SET X AND Y TO CURRENT VALUE	D	111
	C		D	112
298		X=XO	D	113
299		Y=YO	D	114
300		IN=IN+1	D	115
301		GO TO 12	D	116
302	11	IF IT=1	D	117
303		XPR=X	D	118
304		YPR=Y	D	119
	C		D	120
	C	-----EVALUATE POLYNOMIAL AND DERIVATIVES	D	121
	C		D	122
305	12	ICT=0	D	123
306	13	UX=0.0	D	124
307		UY=0.0	D	125
308		V=0.0	D	126
309		YT=0.0	D	127
310		XT=1.0	D	128
311		U=COF(N+1)	D	129
312		IF (U) 14,27,14	D	130
313	14	DO 15 I=1,N	D	131
314		L=N-I+1	D	132
315		TEMP=COF(L)	D	133
316		XT2=X*XT-Y*YT	D	134
317		YT2=X*YT+Y*XT	D	135
318		U=U+TEMP*XT2	D	136
319		V=V+TEMP*YT2	D	137
320		FI=1	D	138
321		UX=UX+FI*XT*TEMP	D	139
322		UY=UY-FI*YT*TEMP	D	140
323		XT=XT2	D	141
324	15	YT=YT2	D	142
325		SUMSQ=UX*UX+UY*UY	D	143
326		IF (SUMSQ) 16,23,16	D	144
327	16	DX=(V*UY-U*UX)/SUMSQ	D	145
328		X=X+DX	D	146

329		DY=-(U*UY+V*UX)/SUMSQ	D 147
330		Y=Y+DY	D 148
331		IF (DABS(DY)+DABS(DX)-1.0D-5) 21,17,17	D 149
	C		D 150
	C	-----STEP ITERATION COUNTER	D 151
	C		D 152
332	17	ICT=ICT+1	D 153
333		IF (ICT-500) 13,18,18	D 154
334	18	IF (IFIT) 21,19,21	D 155
335	19	IF (IN-5) 10,20,20	D 156
	C		D 157
	C	-----SET ERROR CODE TO 3	D 158
	C		D 159
336	20	IER=3	D 160
337		GO TO 3	D 161
338	21	DO 22 L=1,NXX	D 162
339		MT=KJ1-L+1	D 163
340		TEMP=XCOF(MT)	D 164
341		XCOF(MT)=COF(L)	D 165
342	22	COF(L)=TEMP	D 166
343		ITEMP=N	D 167
344		N=NX	D 168
345		NX=ITEMP	D 169
346		IF (IFIT) 25,11,25	D 170
347	23	IF (IFIT) 24,10,24	D 171
348	24	X=XPR	D 172
349		Y=YPR	D 173
350	25	IFIT=0	D 174
351		IF (DABS(Y)-1.0D-4*DABS(X)) 28,26,26	D 175
352	26	ALPHA=X*X	D 176
353		SUMSQ=X*X+Y*Y	D 177
354		N=N-2	D 178
355		GO TO 29	D 179
356	27	X=0.0	D 180
357		NX=NX-1	D 181
358		NXX=NXX-1	D 182
359	28	Y=0.0	D 183
360		SUMSQ=0.0	D 184
361		ALPHA=X	D 185
362		N=N-1	D 186
363	29	COF(2)=COF(2)+ALPHA*COF(1)	D 187
364		IF (N.EQ.0) GO TO 31	D 188
365		DO 30 L=2,N	D 189
366	30	COF(L+1)=COF(L+1)+ALPHA*COF(L)-SUMSQ*COF(L-1)	D 190
367	31	ROOT1(N2)=Y	D 191
368		ROOTR(N2)=X	D 192
369		N2=N2+1	D 193
370		IF (SUMSQ) 32,33,32	D 194
371	32	Y=-Y	D 195
372		SUMSQ=0.0	D 196
373		GO TO 31	D 197
374	33	IF (N) 3,3,9	D 198
375		END	D 199
376		SUBROUTINE TPLOT (M1,M0,M2,M5,M6,M8,M9,JX)	E 1
	C		E 2
	C		E 3
	C	-----TPLOT IS PLOTTING SEVERAL VARIABLES IN ONE GRAPH . THE X-A	E 4
	C	DUES NOT REPRESENT ANY VARIABLE , IT IS INTEGER SEQUENCES .	E 5

FOR NEGATIVE VALUES , THE ZERO LINE IS IN THE 0.5 LINE .

SUPPORTING ROUTINE NONE

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377 IMPLICIT REAL*4(A-H,M-Z)
378 DIMENSION M8(JX),M9(JX),M0(JX),M1(JX),M2(JX),M5(JX),M6(JX)
379 DIMENSION LINE(61),INUM(9)
380 INTEGER PL,M1,S5,BL,SL,S9,S0,S6,S1,S2
381 READ (5,8) PL,M1,S5,BL,SL,S9,S0,S1,S2,S6
382 MXY=0.0
383 MIN0=0.0
384 MIN5=0.0
385 MIN6=0.0
386 MIN2=0.0
387 MIN8=0.0
388 MIN9=0.0
389 PHI5=0.0
390 PHI6=0.0
391 PHI0=0.0
392 PHI2=0.0
393 PHI8=0.0
394 PHI9=0.0
395 DO 1 I=1,JX
396 IF (MIN0.GT.M0(I)) MIN0=M0(I)
397 IF (MIN5.GT.M5(I)) MIN5=M5(I)
398 IF (MIN6.GT.M6(I)) MIN6=M6(I)
399 IF (MIN2.GT.M2(I)) MIN2=M2(I)
400 IF (MIN8.GT.M8(I)) MIN8=M8(I)
401 IF (MIN9.GT.M9(I)) MIN9=M9(I)
402 IF (ABS(M0(I)).GT.PHI0) PHI0=ABS(M0(I))
403 IF (ABS(M2(I)).GT.PHI2) PHI2=ABS(M2(I))
404 IF (ABS(M8(I)).GT.PHI8) PHI8=ABS(M8(I))
405 IF (ABS(M9(I)).GT.PHI9) PHI9=ABS(M9(I))
406 IF (ABS(M5(I)).GT.PHI5) PHI5=ABS(M5(I))
407 IF (ABS(M6(I)).GT.PHI6) PHI6=ABS(M6(I))
408 CONTINUE
409 JJ=JX
410 JJO=JJ*6+1
411 JJ1=JJ+1
412 WRITE (6,9)
413 WRITE (6,10)
414 PHI0=PHI0+ABS(MIN0)
415 PHI5=PHI5+ABS(MIN5)
416 PHI6=PHI6+ABS(MIN6)
417 PHI2=PHI2+ABS(MIN2)
418 PHI8=PHI8+ABS(MIN8)
419 PHI9=PHI9+ABS(MIN9)
420 DO 2 I=1,JJ
421 IF (MIN0.LT.0.0) M0(I)=M0(I)+ABS(MIN0)
422 IF (MIN5.LT.0.0) M5(I)=M5(I)+ABS(MIN5)
423 IF (MIN6.LT.0.0) M6(I)=M6(I)+ABS(MIN6)
424 IF (MIN2.LT.0.0) M2(I)=M2(I)+ABS(MIN2)
425 IF (MIN8.LT.0.0) M8(I)=M8(I)+ABS(MIN8)
426 IF (MIN9.LT.0.0) M9(I)=M9(I)+ABS(MIN9)
427 M0(I)=M0(I)/PHI0
428 M5(I)=M5(I)/PHI5
429 M6(I)=M6(I)/PHI6
430 M2(I)=M2(I)/PHI2
431 M8(I)=M8(I)/PHI8
432 M9(I)=M9(I)/PHI9

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433	DO 3 I=1,9	E 66
434	3 INUM(I)=I	E 67
435	WRITE (6,11) (INUM(I),I=1,9)	E 68
436	DO 7 I=1,JJ1	E 69
437	IF (I.EQ.1) GO TO 5	E 70
438	MXY=MI(I-1)	E 71
439	IP8=M8(I-1)*60+1.0	E 72
440	IP5=M5(I-1)*60+1.0	E 73
441	IP9=M9(I-1)*60+1.0	E 74
442	IP6=M6(I-1)*60+1.0	E 75
443	IP0=M0(I-1)*60+1.0	E 76
444	IP2=M2(I-1)*60+1.0	E 77
445	DO 4 I1=1,56,5	E 78
446	LINE(I1)=BL	E 79
447	DO 4 I2=1,4	E 80
448	I3=I1+I2	E 81
449	IF (I1.EQ.IP0) LINE(I1)=S0	E 82
450	IF (I3.EQ.IP0) LINE(I3)=S0	E 83
451	IF (I1.EQ.IP2) LINE(I1)=S2	E 84
452	IF (I3.EQ.IP2) LINE(I3)=S2	E 85
453	IF (I1.EQ.IP5) LINE(I1)=S5	E 86
454	IF (I3.EQ.IP5) LINE(I3)=S5	E 87
455	IF (I1.EQ.IP6) LINE(I1)=S6	E 88
456	IF (I3.EQ.IP6) LINE(I3)=S6	E 89
457	IF (I1.EQ.IP8) LINE(I1)=SL	E 90
458	IF (I3.EQ.IP8) LINE(I3)=SL	E 91
459	IF (I1.EQ.IP9) LINE(I1)=S9	E 92
460	IF (I3.EQ.IP9) LINE(I3)=S9	E 93
461	4 CONTINUE	E 94
462	LINE(61)=PL	E 95
463	I1L=I-1	E 96
464	IF (I1.EQ.IP0) LINE(I1)=S0	E 97
465	IF (I1.EQ.IP2) LINE(I1)=S2	E 98
466	IF (I1.EQ.IP5) LINE(I1)=S5	E 99
467	IF (I1.EQ.IP6) LINE(I1)=S6	E 100
468	IF (I1.EQ.IP8) LINE(I1)=SL	E 101
469	IF (I1.EQ.IP9) LINE(I1)=S9	E 102
470	IF (IP0.EQ.61) LINE(61)=S0	E 103
471	IF (IP5.EQ.61) LINE(61)=S5	E 104
472	IF (IP2.EQ.61) LINE(61)=S2	E 105
473	IF (IP6.EQ.61) LINE(61)=S6	E 106
474	IF (IP8.EQ.61) LINE(61)=SL	E 107
475	IF (IP9.EQ.61) LINE(61)=S9	E 108
476	IF (IP0.NE.1.OR.IP2.NE.1.OR.IP5.NE.1.OR.IP6.NE.1.OR.IP8.NE.1.OR.IP9.NE.1) LINE(I1)=PL	E 109
477	WRITE (6,12) MXY,(LINE(KK),KK=1,61)	E 110
478	IF (I.EQ.JJ1) GO TO 7	E 111
479	5 CONTINUE	E 112
480	DO 6 I1=1,56,5	E 113
481	DO 6 I2=1,4	E 114
482	I3=I1+I2	E 115
483	LINE(I3)=BL	E 116
484	6 CONTINUE	E 117
485	7 CONTINUE	E 118
486	WRITE (6,13) (INUM(I),I=1,9)	E 119
487	WRITE (6,15) PL,MI,S5,BL,SL,S9,S0,S1,S2,S6	E 120
488	WRITE (6,16)	E 121
489	WRITE (6,14)	E 122
490	STOP	E 123
		E 124
		E 125

491	8	FORMAT (11A1)	E 126
492	9	FORMAT (1H1)	E 127
493	10	FORMAT (35X,16HRELATIVE DENSITY)	E 128
494	11	FORMAT (17X,9(2X,2H0.,11,1H)/,14X,10(6H+-----),1H+)	E 129
495	12	FORMAT (1X,4HTIME,1X,F7.3,1X,61A1)	E 130
496	13	FORMAT (14X,10(6H+-----),1H+/,16X,9(3X,2H0.,11))	E 131
497	14	FORMAT (1H1)	E 132
498	15	FORMAT (15X,17HINPUT CHARACTER ,11A1,/)	E 133
499	16	FORMAT (26H ,/,15X,10H P - POWER,/,15X,18	E 134
		1H L - LOG. OF. POWER,/,15X,22H Q - PRECURSOR DENSITY,/,15X,21H F -	E 135
		2FUEL TEMPERATURE,/,15X,26H M - MODERATOR TEMPERATURE,/,15X,15H R -	E 136
		3 REACTIVITY,/,15X,42HIF VARIABLE PLOTTED HAS NEGATIVE VALUE THE,/,	E 137
		415X,41HAXIS FOR THIS VARIABLE IS SHIFTED TO THE ,/,15X,16HCENTER O	E 138
		5F Y-AXIS,///)	E 139
500		END	E 140
501		SUBROUTINE GXN (G,XN,N,GX)	F 1
	C		F 2
	C		F 3
	C	-----MULTIPLY THE G - MATRIX WITH THE INITIAL VECTOR CO	F 4
	C	LUMN TO GET THE NEXT ITERATION .	F 5
	C		F 6
	C	SUPPORTING ROUTINE, NONE	F 7
	C		F 8
	C		F 9
502		DIMENSION G(N,N),XN(N),GX(N)	F 10
503		DO 2 I=1,N	F 11
504		GG=0.0	F 12
505		DO 1 J=1,N	F 13
506	1	GG=GG+G(I,J)*XN(J)	F 14
507	2	GX(I)=GG	F 15
508		RETURN	F 16
509		END	F 17
510		SUBROUTINE AMTRX (A)	G 1
	C		G 2
	C		G 3
	C	-----FORM A-MATRIX IN EQUATION 2.4.12 .	G 4
	C	SUPPORTING ROUTINE, NONE	G 5
	C		G 6
	C		G 7
	C		G 8
511		DIMENSION A(4,4)	G 9
512		COMMON BETA,XL,X,FX,RF,DH,FDENS,CPF,VF,FH,VM,DX,DOTMG,W0,HX,RR,C	G 10
		1PMX,PI,NROD	G 11
513		A(1,1)=(RR-BETA)/XL	G 12
514		A(1,2)=X	G 13
515		A(1,3)=0.0	G 14
516		A(1,4)=0.0	G 15
517		A(2,1)=BETA/XL	G 16
518		A(2,2)=-X	G 17
519		A(2,3)=0.0	G 18
520		A(2,4)=0.0	G 19
521		A(3,1)=1.0/(2.0*CPF)	G 20
522		A(3,2)=0.0	G 21
523		A(3,3)=-4.*PI*FX*FH*NROD/CPF	G 22
524		A(3,4)=-A(3,3)	G 23
525		A(4,1)=1.0/(DX*CPMX*VM)	G 24
526		A(4,2)=0.0	

527	A(4,3)=0.0	G	25
528	A(4,4)=-2.0*DOTMG/(DX*VM)	G	26
529	RETURN	G	27
530	END	G	28
531	SUBROUTINE EIGEN4 (A,B)	H	1
C		H	2
C		H	3
C	-----EIGEN4 IS SUBSTITUTING DETERMINANT A TO POLYNOMIAL FORM .	H	4
C	SUPPORTING ROUTINE NONE	H	5
C		H	6
C		H	7
532	DIMENSION A(4,4),B(5)	H	8
533	DOUBLE PRECISION,B	H	9
534	B(1)=A(1,1)*A(2,2)*A(3,3)*A(4,4)-A(1,2)*A(2,1)*A(3,3)*A(4,4)	H	10
535	B(2)=-(A(1,1)*A(2,2)*A(3,3)+A(1,1)*A(2,2)*A(4,4)+A(1,1)*A(3,3)*A(4,4)+A(2,2)*A(3,3)*A(4,4)-A(1,2)*A(2,1)*A(3,3)-A(1,2)*A(2,1)*A(4,4)+A(2,2)*A(3,3)*A(4,4)-A(1,2)*A(2,1)*A(3,3)-A(1,2)*A(2,1)*A(4,4))	H	11
		H	12
		H	13
536	B(3)=A(3,3)*A(4,4)+A(1,1)*A(2,2)+A(1,1)*A(3,3)+A(1,1)*A(4,4)+A(2,2)*A(3,3)+A(2,2)*A(4,4)-A(1,2)*A(2,1)	H	14
		H	15
537	B(4)=-(A(1,1)+A(2,2)+A(3,3)+A(4,4))	H	16
538	B(5)=1.0	H	17
539	RETURN	H	18
540	END	H	19

//DATA

MODULE 2 , FEEDBACK

INPUT DATA

BETA	LAMBDA	NEUTRON GEN. TIME	INITIAL REACTIVITY	
()	(SEC**-1)	(SEC)	(%)	
0.006450	0.07695	0.100E-03	0.30	
INSERTION TIME		CONSTANT PERIOD		
(SEC)		(1/SEC)		
1.000		1.500		
COOLANT COEFF.	CONSTANT B	RESONANCE ESC. PROB.		
()	()	()		
-0.50000E-04	*****	0.800		
FUEL RADIUS	PITCH	FUEL CP.	FUEL DENSITY	FUEL HEIGHT
(FT)	(FT)	(BTU/LB-F)	(LB/FT**3)	(FT)
0.01504	0.04733	0.0590	43.20	12.00
NUMBER OF ASSEMBLY		NUMBER OF ROD PER ASSEMBLY		
()		()		
145		208		
COOLANT VELOCITY	INLET COOLANT TEMPT.			
(FT/SEC)	(F)			
13.0	400.0			
OPTION	PLOT	1		
	TYPE INSR.	1		
INITIAL POWER	(MW)	0.1000E 04		
GUESSED FUEL TEMPT.	(F)	500.0		
GUESSED COOLANT TEMPT.	(F)	200.0		

END OF INPUT DATA

END TIME TIME INCREMENT
8.000 0.0100

CONVERGENCE FACTOR = 0.01000

EQUILIBRIUM STATE INITIALLY

POWER	PRECURSOR	FUEL TEMPT.	COOLANT TEMPT.	REACTIVITY
0.1000E 04	0.8382E 06	504.31	412.73	0.000

REYNOLDS #	PRANDTL #	COLBURN #	MASS FLOW (LB/HR)	HEAT TRANSFER COEFF. (BTU/HR-F-FT**2)
509503.6	1.0	0.042	0.1172E 09	9044.7

NO	FUEL TEMP (F)	TIME (SEC)	MOD. TEMP (F)	REACTIVITY (%)	EXIT TEMP (F)	POWER (MW)	FUEL COEFF. (1/E)	PRECURSOR (MW)
1	504.31	0.000	412.73	0.00	425.47	0.1000E 04	0.4968E-06	0.8382E 06
2	496.42	0.046	413.04	0.30	426.08	0.9345E 03	0.5008E-06	0.8381E 06
3	497.10	0.096	412.92	0.30	425.84	0.1373E 04	0.5004E-06	0.8379E 06
4	510.84	0.146	413.39	0.30	426.79	0.1420E 04	0.4937E-06	0.8391E 06
5	521.65	0.196	413.88	0.29	427.76	0.1422E 04	0.4885E-06	0.8404E 06
6	529.63	0.246	414.32	0.29	428.64	0.1419E 04	0.4848E-06	0.8418E 06
7	535.33	0.296	414.71	0.29	429.42	0.1416E 04	0.4822E-06	0.8431E 06
8	539.41	0.346	415.06	0.28	430.12	0.1413E 04	0.4804E-06	0.8444E 06
9	542.34	0.396	415.37	0.28	430.74	0.1410E 04	0.4791E-06	0.8457E 06
10	544.44	0.446	415.65	0.28	431.30	0.1408E 04	0.4782E-06	0.8470E 06
11	545.97	0.496	415.90	0.28	431.79	0.1406E 04	0.4775E-06	0.8483E 06
12	547.09	0.546	416.12	0.28	432.23	0.1404E 04	0.4770E-06	0.8496E 06
13	547.91	0.596	416.31	0.28	432.62	0.1403E 04	0.4766E-06	0.8508E 06
14	548.53	0.646	416.49	0.27	432.97	0.1403E 04	0.4764E-06	0.8521E 06
15	549.01	0.696	416.64	0.27	433.29	0.1402E 04	0.4762E-06	0.8533E 06
16		0.746		0.27		0.1402E 04		0.8546E 06

17	549.38	416.78	433.57	0.4760E-06	
	0.796	0.27	0.1402E 04	0.8558E 06	
18	549.69	416.91	433.82	0.4759E-06	
	0.846	0.27	0.1402E 04	0.8570E 06	
19	549.94	417.02	434.05	0.4758E-06	
	0.896	0.27	0.1402E 04	0.8583E 06	
20	550.17	417.13	434.26	0.4757E-06	
	0.946	0.27	0.1403E 04	0.8595E 06	
21	550.37	417.22	434.44	0.4756E-06	
	0.996	0.27	0.1403E 04	0.8607E 06	
22	550.56	417.31	434.61	0.4755E-06	
	1.046	0.27	0.1404E 04	0.8619E 06	
23	550.74	417.38	434.76	0.4754E-06	
	1.096	-0.03	0.1405E 04	0.8631E 06	
24	550.91	417.45	434.90	0.4753E-06	
	1.146	-0.03	0.1011E 04	0.8643E 06	
25	550.95	417.52	435.03	0.4753E-06	
	1.196	-0.03	0.9980E 03	0.8643E 06	
26	539.15	417.05	434.10	0.4805E-06	
	1.246	-0.03	0.1000E 04	0.8642E 06	
27	530.19	416.62	433.23	0.4845E-06	
	1.296	-0.03	0.1003E 04	0.8641E 06	
28	523.75	416.23	432.46	0.4875E-06	
	1.346	-0.02	0.1005E 04	0.8640E 06	
29	519.15	415.88	431.77	0.4897E-06	
	1.396	-0.02	0.1007E 04	0.8639E 06	
30	515.87	415.58	431.15	0.4912E-06	
	1.446	-0.02	0.1009E 04	0.8638E 06	
31	513.54	415.30	430.60	0.4923E-06	
	1.496	-0.02	0.1011E 04	0.8637E 06	
32	511.87	415.06	430.11	0.4931E-06	
	1.546	-0.02	0.1013E 04	0.8637E 06	
33	510.67	414.84	429.68	0.4937E-06	
	1.596	-0.01	0.1014E 04	0.8636E 06	
34	509.82	414.64	429.29	0.4941E-06	
	1.646	-0.01	0.1016E 04	0.8636E 06	
35	509.20	414.47	428.95	0.4944E-06	
	1.696	-0.01	0.1017E 04	0.8635E 06	
36	508.76	414.32	428.64	0.4946E-06	
	1.746	-0.01	0.1018E 04	0.8635E 06	
37	508.43	414.18	428.37	0.4948E-06	
	1.796	-0.01	0.1019E 04	0.8634E 06	
38	508.20	414.06	428.12	0.4949E-06	
	1.846	-0.01	0.1020E 04	0.8634E 06	
39	508.02	413.95	427.91	0.4950E-06	
	1.896	-0.01	0.1021E 04	0.8634E 06	
40	507.89	413.86	427.72	0.4951E-06	
	1.946	-0.01	0.1021E 04	0.8633E 06	
41	507.79	413.77	427.55	0.4951E-06	
	1.996	-0.01	0.1022E 04	0.8633E 06	
42	507.71	413.70	427.39	0.4952E-06	
	2.046	-0.01	0.1023E 04	0.8633E 06	
43	507.66	413.63	427.26	0.4952E-06	
	2.096	-0.01	0.1023E 04	0.8632E 06	
44	507.61	413.57	427.14	0.4952E-06	
	2.146	-0.01	0.1023E 04	0.8632E 06	
45	507.57	413.52	427.03	0.4952E-06	
	2.196	-0.01	0.1024E 04	0.8632E 06	
46	507.54	413.47	426.94	0.4952E-06	
	2.246	-0.01	0.1024E 04	0.8632E 06	

	507.52	413.43	426.85	0.4953E-06
47	2.296	-0.00	0.1025E 04	0.8632E 06
	507.50	413.39	426.78	0.4953E-06
48	2.346	-0.00	0.1025E 04	0.8631E 06
	507.48	413.35	426.71	0.4953E-06
49	2.396	-0.00	0.1025E 04	0.8631E 06
	507.47	413.33	426.65	0.4953E-06
50	2.446	-0.00	0.1025E 04	0.8631E 06
	507.46	413.30	426.60	0.4953E-06
51	2.496	-0.00	0.1025E 04	0.8631E 06
	507.45	413.28	426.55	0.4953E-06
52	2.546	-0.00	0.1026E 04	0.8631E 06
	507.44	413.25	426.51	0.4953E-06
53	2.596	-0.00	0.1026E 04	0.8631E 06
	507.43	413.24	426.47	0.4953E-06
54	2.646	-0.00	0.1026E 04	0.8631E 06
	507.42	413.22	426.44	0.4953E-06
55	2.696	-0.00	0.1026E 04	0.8631E 06
	507.41	413.20	426.41	0.4953E-06
56	2.746	-0.00	0.1026E 04	0.8630E 06
	507.41	413.19	426.38	0.4953E-06
57	2.796	-0.00	0.1026E 04	0.8630E 06
	507.40	413.18	426.36	0.4953E-06
58	2.846	-0.00	0.1026E 04	0.8630E 06
	507.40	413.17	426.34	0.4953E-06
59	2.896	-0.00	0.1026E 04	0.8630E 06
	507.39	413.16	426.32	0.4953E-06
60	2.946	-0.00	0.1026E 04	0.8630E 06
	507.39	413.15	426.30	0.4953E-06
61	2.996	-0.00	0.1026E 04	0.8630E 06
	507.39	413.14	426.29	0.4953E-06
62	3.046	-0.00	0.1027E 04	0.8630E 06
	507.38	413.14	426.28	0.4953E-06
63	3.096	-0.00	0.1027E 04	0.8630E 06
	507.38	413.13	426.26	0.4953E-06
64	3.146	-0.00	0.1027E 04	0.8630E 06
	507.38	413.13	426.25	0.4953E-06
65	3.196	-0.00	0.1027E 04	0.8629E 06
	507.38	413.12	426.24	0.4953E-06
66	3.246	-0.00	0.1027E 04	0.8629E 06
	507.37	413.12	426.24	0.4953E-06
67	3.296	-0.00	0.1027E 04	0.8629E 06
	507.37	413.11	426.23	0.4953E-06
68	3.346	-0.00	0.1027E 04	0.8629E 06
	507.37	413.11	426.22	0.4953E-06
69	3.396	-0.00	0.1027E 04	0.8629E 06
	507.37	413.11	426.22	0.4953E-06
70	3.446	-0.00	0.1027E 04	0.8629E 06
	507.36	413.11	426.21	0.4953E-06
71	3.496	-0.00	0.1027E 04	0.8629E 06
	507.36	413.10	426.21	0.4953E-06
72	3.546	-0.00	0.1027E 04	0.8629E 06
	507.36	413.10	426.20	0.4953E-06
73	3.596	-0.00	0.1027E 04	0.8629E 06
	507.36	413.10	426.20	0.4953E-06
74	3.646	-0.00	0.1027E 04	0.8629E 06
	507.36	413.10	426.20	0.4953E-06
75	3.696	-0.00	0.1027E 04	0.8629E 06
	507.36	413.10	426.19	0.4953E-06
76	3.746	-0.00	0.1027E 04	0.8628E 06

77	507.35	3.796	413.10	-0.00	426.19	0.1027E 04	0.4953E-06	0.8628E 06
78	507.35	3.846	413.09	-0.00	426.19	0.1027E 04	0.4953E-06	0.8628E 06
79	507.35	3.896	413.09	-0.00	426.19	0.1027E 04	0.4953E-06	0.8628E 06
80	507.35	3.946	413.09	-0.00	426.18	0.1027E 04	0.4953E-06	0.8628E 06
81	507.35	3.996	413.09	-0.00	426.18	0.1027E 04	0.4953E-06	0.8628E 06
82	507.35	4.046	413.09	-0.00	426.18	0.1027E 04	0.4953E-06	0.8628E 06
83	507.35	4.096	413.09	-0.00	426.18	0.1027E 04	0.4953E-06	0.8628E 06
84	507.34	4.146	413.09	-0.00	426.18	0.1027E 04	0.4953E-06	0.8628E 06
85	507.34	4.196	413.09	-0.00	426.17	0.1027E 04	0.4953E-06	0.8628E 06
86	507.34	4.246	413.09	-0.00	426.17	0.1027E 04	0.4953E-06	0.8628E 06
87	507.34	4.296	413.09	-0.00	426.17	0.1027E 04	0.4953E-06	0.8627E 06
88	507.34	4.346	413.09	-0.00	426.17	0.1027E 04	0.4953E-06	0.8627E 06
89	507.34	4.396	413.09	-0.00	426.17	0.1027E 04	0.4953E-06	0.8627E 06
90	507.34	4.446	413.09	-0.00	426.17	0.1027E 04	0.4953E-06	0.8627E 06
91	507.34	4.496	413.09	-0.00	426.17	0.1027E 04	0.4953E-06	0.8627E 06
92	507.33	4.546	413.09	-0.00	426.17	0.1027E 04	0.4953E-06	0.8627E 06
93	507.33	4.596	413.08	-0.00	426.17	0.1027E 04	0.4953E-06	0.8627E 06
94	507.33	4.646	413.08	-0.00	426.17	0.1027E 04	0.4953E-06	0.8627E 06
95	507.33	4.696	413.08	-0.00	426.17	0.1027E 04	0.4953E-06	0.8627E 06
96	507.33	4.746	413.08	-0.00	426.17	0.1027E 04	0.4953E-06	0.8627E 06
97	507.33	4.796	413.08	-0.00	426.17	0.1027E 04	0.4953E-06	0.8627E 06
98	507.33	4.846	413.08	-0.00	426.17	0.1027E 04	0.4953E-06	0.8627E 06
99	507.33	4.896	413.08	-0.00	426.17	0.1027E 04	0.4953E-06	0.8626E 06
100	507.33	4.946	413.08	-0.00	426.17	0.1027E 04	0.4953E-06	0.8626E 06
101	507.32	4.996	413.08	-0.00	426.17	0.1027E 04	0.4953E-06	0.8626E 06
102	507.32	5.046	413.08	-0.00	426.17	0.1027E 04	0.4953E-06	0.8626E 06
103	507.32	5.096	413.08	-0.00	426.17	0.1027E 04	0.4953E-06	0.8626E 06
104	507.32	5.146	413.08	-0.00	426.17	0.1027E 04	0.4953E-06	0.8626E 06
105	507.32	5.196	413.08	-0.00	426.17	0.1027E 04	0.4953E-06	0.8626E 06
106	507.32	5.246	413.08	-0.00	426.17	0.1027E 04	0.4953E-06	0.8626E 06

	507.32	413.08	426.17	0.4954E-06	
107	5.296	-0.00	0.1027E 04	0.8626E 06	
	507.32	413.08	426.17	0.4954E-06	
108	5.346	-0.00	0.1027E 04	0.8626E 06	
	507.32	413.08	426.17	0.4954E-06	
109	5.396	-0.00	0.1027E 04	0.8626E 06	
	507.32	413.08	426.17	0.4954E-06	
110	5.446	-0.00	0.1027E 04	0.8625E 06	
	507.32	413.08	426.17	0.4954E-06	
111	5.496	-0.00	0.1026E 04	0.8625E 06	
	507.31	413.08	426.17	0.4954E-06	
112	5.546	-0.00	0.1026E 04	0.8625E 06	
	507.31	413.08	426.17	0.4954E-06	
113	5.596	-0.00	0.1026E 04	0.8625E 06	
	507.31	413.08	426.16	0.4954E-06	
114	5.646	-0.00	0.1026E 04	0.8625E 06	
	507.31	413.08	426.16	0.4954E-06	
115	5.696	-0.00	0.1026E 04	0.8625E 06	
	507.31	413.08	426.16	0.4954E-06	
116	5.746	-0.00	0.1026E 04	0.8625E 06	
	507.31	413.08	426.16	0.4954E-06	
117	5.796	-0.00	0.1026E 04	0.8625E 06	
	507.31	413.08	426.16	0.4954E-06	
118	5.846	-0.00	0.1026E 04	0.8625E 06	
	507.31	413.08	426.16	0.4954E-06	
119	5.896	-0.00	0.1026E 04	0.8625E 06	
	507.31	413.08	426.16	0.4954E-06	
120	5.946	-0.00	0.1026E 04	0.8625E 06	
	507.30	413.08	426.16	0.4954E-06	
121	5.996	-0.00	0.1026E 04	0.8625E 06	
	507.30	413.08	426.16	0.4954E-06	
122	6.046	-0.00	0.1026E 04	0.8624E 06	
	507.30	413.08	426.16	0.4954E-06	
123	6.096	-0.00	0.1026E 04	0.8624E 06	
	507.30	413.08	426.16	0.4954E-06	
124	6.146	-0.00	0.1026E 04	0.8624E 06	
	507.30	413.08	426.16	0.4954E-06	
125	6.196	-0.00	0.1026E 04	0.8624E 06	
	507.30	413.08	426.16	0.4954E-06	
126	6.246	-0.00	0.1026E 04	0.8624E 06	
	507.30	413.08	426.16	0.4954E-06	
127	6.296	-0.00	0.1026E 04	0.8624E 06	
	507.30	413.08	426.16	0.4954E-06	
128	6.346	-0.00	0.1026E 04	0.8624E 06	
	507.30	413.08	426.16	0.4954E-06	
129	6.396	-0.00	0.1026E 04	0.8624E 06	
	507.30	413.08	426.16	0.4954E-06	
130	6.446	-0.00	0.1026E 04	0.8624E 06	
	507.29	413.08	426.16	0.4954E-06	
131	6.496	-0.00	0.1026E 04	0.8624E 06	
	507.29	413.08	426.16	0.4954E-06	
132	6.546	-0.00	0.1026E 04	0.8624E 06	
	507.29	413.08	426.16	0.4954E-06	
133	6.596	-0.00	0.1026E 04	0.8624E 06	
	507.29	413.08	426.16	0.4954E-06	
134	6.646	-0.00	0.1026E 04	0.8623E 06	
	507.29	413.08	426.16	0.4954E-06	
135	6.696	-0.00	0.1026E 04	0.8623E 06	
	507.29	413.08	426.16	0.4954E-06	
136	6.746	-0.00	0.1026E 04	0.8623E 06	

137	507.29	6.796	413.08	-0.00	426.16	0.1026E 04	0.4954E-06	0.8623E 06
138	507.29	6.846	413.08	-0.00	426.16	0.1026E 04	0.4954E-06	0.8623E 06
139	507.29	6.896	413.08	-0.00	426.16	0.1026E 04	0.4954E-06	0.8623E 06
140	507.28	6.946	413.08	-0.00	426.16	0.1026E 04	0.4954E-06	0.8623E 06
141	507.28	6.996	413.08	-0.00	426.16	0.1026E 04	0.4954E-06	0.8623E 06
142	507.28	7.046	413.08	-0.00	426.16	0.1026E 04	0.4954E-06	0.8623E 06
143	507.28	7.096	413.08	-0.00	426.16	0.1026E 04	0.4954E-06	0.8623E 06
144	507.28	7.146	413.08	-0.00	426.16	0.1026E 04	0.4954E-06	0.8623E 06
145	507.28	7.196	413.08	-0.00	426.16	0.1026E 04	0.4954E-06	0.8623E 06
146	507.28	7.246	413.08	-0.00	426.16	0.1026E 04	0.4954E-06	0.8622E 06
147	507.28	7.296	413.08	-0.00	426.16	0.1026E 04	0.4954E-06	0.8622E 06
148	507.28	7.346	413.08	-0.00	426.16	0.1026E 04	0.4954E-06	0.8622E 06
149	507.27	7.396	413.08	-0.00	426.16	0.1026E 04	0.4954E-06	0.8622E 06
150	507.27	7.446	413.08	-0.00	426.16	0.1026E 04	0.4954E-06	0.8622E 06
151	507.27	7.496	413.08	-0.00	426.16	0.1026E 04	0.4954E-06	0.8622E 06
152	507.27	7.546	413.08	-0.00	426.15	0.1026E 04	0.4954E-06	0.8622E 06
153	507.27	7.596	413.08	-0.00	426.15	0.1026E 04	0.4954E-06	0.8622E 06
154	507.27	7.646	413.08	-0.00	426.15	0.1026E 04	0.4954E-06	0.8622E 06
155	507.27	7.696	413.08	-0.00	426.15	0.1026E 04	0.4954E-06	0.8622E 06
156	507.27	7.746	413.08	-0.00	426.15	0.1026E 04	0.4954E-06	0.8622E 06
157	507.27	7.796	413.08	-0.00	426.15	0.1026E 04	0.4954E-06	0.8622E 06
158	507.27	7.846	413.08	-0.00	426.15	0.1026E 04	0.4954E-06	0.8621E 06
159	507.26	7.896	413.08	-0.00	426.15	0.1026E 04	0.4954E-06	0.8621E 06
160	507.26	7.946	413.08	-0.00	426.15	0.1026E 04	0.4954E-06	0.8621E 06
161	507.26	7.996	413.07	-0.00	426.15	0.1026E 04	0.4954E-06	0.8621E 06

		RELATIVE DENSITY								
		0.1	0.2	0.3	0.4	0.5	0.6	0.7	0.8	0.9
TIME	0:000	R						P		F LQM
TIME	0:046							P		F L QRI
TIME	0:096									F PQMR
TIME	0:146									F QRI
TIME	0:196									F QRL
TIME	0:246									FRMI
TIME	0:296									RM
TIME	0:346									RFMI
TIME	0:396									RQFI
TIME	0:446									R QFI
TIME	0:496									R QFI
TIME	0:546									R QFI
TIME	0:596									R F
TIME	0:646									R F
TIME	0:696									R F
TIME	0:746									R F
TIME	0:796									R F
TIME	0:846									R F
TIME	0:896									R R F
TIME	0:946									R R F
TIME	0:996									R R F
TIME	1:046									R F
TIME	1:096									R F
TIME	1:146							P		L F
TIME	1:196							P		LFMI
TIME	1:246							P		F MI
TIME	1:296	R						P		F MI
TIME	1:346	R						P		FL MI
TIME	1:396	R						P		FL MI
TIME	1:446	R						P		F L MI
TIME	1:496	R						P		F L MI
TIME	1:546	R						P		F L MI
TIME	1:596	R						P		F L MI
TIME	1:646	R						P		F L MI
TIME	1:696	R						P		F L MI
TIME	1:746	R						P		F L MI
TIME	1:796	R						P		F L MI
TIME	1:846	R						P		F L MI
TIME	1:896	R						P		F L MI
TIME	1:946	R						P		F L MI
TIME	1:996	R						P		F L MI
TIME	2:046	R						P		F L MI
TIME	2:096	R						P		F L MI
TIME	2:146	R						P		F L MI
TIME	2:196	R						P		F L MI
TIME	2:246	R						P		F L MI
TIME	2:296	R						P		F L MI
TIME	2:346	R						P		F L MI
TIME	2:396	R						P		F L MI
TIME	2:446	R						P		F L MI
TIME	2:496	R						P		F L MI
TIME	2:546	R						P		F L MI
TIME	2:596	R						P		F L MI
TIME	2:646	R						P		F L MI
TIME	2:696	R						P		F L MI
TIME	2:746	R						P		F L MI

TIME	2.796	R	P	F	L	M
TIME	2.846	R	P	F	L	M
TIME	2.896	R	P	F	L	M
TIME	2.946	R	P	F	L	M
TIME	2.996	R	P	F	L	M
TIME	3.046	R	P	F	L	M
TIME	3.096	R	P	F	L	M
TIME	3.146	R	P	F	L	M
TIME	3.196	R	P	F	L	M
TIME	3.246	R	P	F	L	M
TIME	3.296	R	P	F	L	M
TIME	3.346	R	P	F	L	M
TIME	3.396	R	P	F	L	M
TIME	3.446	R	P	F	L	M
TIME	3.496	R	P	F	L	M
TIME	3.546	R	P	F	L	M
TIME	3.596	R	P	F	L	M
TIME	3.646	R	P	F	L	M
TIME	3.696	R	P	F	L	M
TIME	3.746	R	P	F	L	M
TIME	3.796	R	P	F	L	M
TIME	3.846	R	P	F	L	M
TIME	3.896	R	P	F	L	M
TIME	3.946	R	P	F	L	M
TIME	3.996	R	P	F	L	M
TIME	4.046	R	P	F	L	M
TIME	4.096	R	P	F	L	M
TIME	4.146	R	P	F	L	M
TIME	4.196	R	P	F	L	M
TIME	4.246	R	P	F	L	M
TIME	4.296	R	P	F	L	M
TIME	4.346	R	P	F	L	M
TIME	4.396	R	P	F	L	M
TIME	4.446	R	P	F	L	M
TIME	4.496	R	P	F	L	M
TIME	4.546	R	P	F	L	M
TIME	4.596	R	P	F	L	M
TIME	4.646	R	P	F	L	M
TIME	4.696	R	P	F	L	M
TIME	4.746	R	P	F	L	M
TIME	4.796	R	P	F	L	M
TIME	4.846	R	P	F	L	M
TIME	4.896	R	P	F	L	M
TIME	4.946	R	P	F	L	M
TIME	4.996	R	P	F	L	M
TIME	5.046	R	P	F	L	M
TIME	5.096	R	P	F	L	M
TIME	5.146	R	P	F	L	M
TIME	5.196	R	P	F	L	M
TIME	5.246	R	P	F	L	M
TIME	5.296	R	P	F	L	M
TIME	5.346	R	P	F	L	M
TIME	5.396	R	P	F	L	M
TIME	5.446	R	P	F	L	M
TIME	5.496	R	P	F	L	M
TIME	5.546	R	P	F	L	M
TIME	5.596	R	P	F	L	M
TIME	5.646	R	P	F	L	M
TIME	5.696	R	P	F	L	M
TIME	5.746	R	P	F	L	M

TIME	5.796	R	P	F	L	M
TIME	5.846	R	P	F	L	M
TIME	5.896	R	P	F	L	M
TIME	5.946	R	P	F	L	M
TIME	5.996	R	P	F	L	M
TIME	6.046	R	P	F	L	M
TIME	6.096	R	P	F	L	M
TIME	6.146	R	P	F	L	M
TIME	6.196	R	P	F	L	M
TIME	6.246	R	P	F	L	M
TIME	6.296	R	P	F	L	M
TIME	6.346	R	P	F	L	M
TIME	6.396	R	P	F	L	M
TIME	6.446	R	P	F	L	M
TIME	6.496	R	P	F	L	M
TIME	6.546	R	P	F	L	M
TIME	6.596	R	P	F	L	M
TIME	6.646	R	P	F	L	M
TIME	6.696	R	P	F	L	M
TIME	6.746	R	P	F	L	M
TIME	6.796	R	P	F	L	M
TIME	6.846	R	P	F	L	M
TIME	6.896	R	P	F	L	M
TIME	6.946	R	P	F	L	M
TIME	6.996	R	P	F	L	M
TIME	7.046	R	P	F	L	M
TIME	7.096	R	P	F	L	M
TIME	7.146	R	P	F	L	M
TIME	7.196	R	P	F	L	M
TIME	7.246	R	P	F	L	M
TIME	7.296	R	P	F	L	M
TIME	7.346	R	P	F	L	M
TIME	7.396	R	P	F	L	M
TIME	7.446	R	P	F	L	M
TIME	7.496	R	P	F	L	M
TIME	7.546	R	P	F	L	M
TIME	7.596	R	P	F	L	M
TIME	7.646	R	P	F	L	M
TIME	7.696	R	P	F	L	M
TIME	7.746	R	P	F	L	M
TIME	7.796	R	P	F	L	M
TIME	7.846	R	P	F	L	M
TIME	7.896	R	P	F	L	M
TIME	7.946	R	P	F	L	M
TIME	7.996	R	P	F	L	M

0.1 0.2 0.3 0.4 0.5 0.6 0.7 0.8 0.9
 INPUT CHARACTER I*Q FRPILM

P - POWER
 L - LOG. OF POWER
 Q - PRECURSOR DENSITY
 F - FUEL TEMPERATURE
 M - MODERATOR TEMPERATURE
 R - REACTIVITY

IF VARIABLE PLOTTED HAS NEGATIVE VALUE THE
 AXIS FOR THIS VARIABLE IS SHIFTED TO THE
 CENTER OF Y-AXIS